

1 **On the petrology and microstructures of small-scale ductile shear zones in granitoid rocks: an**  
2 **overview**

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10 **Abstract.**

11 Since the pioneering works of John Ramsay in the 1970's and 1980's, the analysis of exceptional  
12 exposures of small-scale shear zones (i.e.  $10^{-3}$  –  $10^{-1}$  m thick) in granitoid rocks provided invaluable  
13 insights into the processes controlling strain localisation in the middle and lower continental crust.  
14 Indeed, recent advancement in field, microstructural and petrological analyses of such small-scale  
15 shear zone have shed new light on the metamorphic, tectonic and fluid conditions promoting shear  
16 zone nucleation and development in granitoid rocks. In this paper we provide an overview of these  
17 new insights, comparing and integrating the results obtained from field, and microstructural and  
18 petrological analyses of small-scale shear zones in granitoid plutons and meta-granitoids from the  
19 Alps. A review of the deformation temperature shows that the granitoid shear zones development  
20 occurs between 350 to 600°C, with most of them localising in a restricted temperature window  
21 between 450 and 500°C. At these conditions, the magmatic assemblage is metastable and subjected  
22 to a series of metamorphic reactions. Furthermore, the development of shear zone does not occur  
23 under-closed system conditions. Introducing or expelling fluids and mass (i.e. metasomatism) during  
24 deformation has mineralogical consequences that control the rheology and the way shear zone  
25 evolves. Among the main mineralogical and microstructural changes, the breakdown of magmatic

26 feldspar(s) into fine-grained aggregates steers both the rheology and fabric evolution of shear zones  
27 in granitoid rocks, triggering further mechano-chemical feedback mechanisms. Future research  
28 should consider the occurrence of feedback processes between deformation, metamorphic and  
29 metasomatic processes to understand and quantify the evolution with time and strain of shear zone  
30 geometry and rheology, as well as of the development of larger-scale shear zone networks.

31 **Keywords**

32 Small-scale shear zones; Strain localisation, Granitoid rocks, Microstructure; Fluid-rock interaction.

33

## 35 1. Introduction

36 The seminal paper by Ramsay and Graham (1970) has formalised the concept of localised shear zone  
37 as a tabular-shaped zone of high strain. The shear zone they described is typically characterized by  
38 the occurrence of a strain-induced foliation with a sigmoidal-shaped geometry reflecting the  
39 symmetric decrease of strain from the centre to the boundary of the shear zone (Fig. 1a-b). The  
40 inferences of Ramsay and Graham about ductility and strain softening processes controlling shear  
41 zone development (“[...] *the structure in them appear to have been produced by ductile flow of the*  
42 *rock within the zone. [...] Why then should the deformation become localized in these narrow zones?*  
43 *Perhaps it is because the initially deforming material became “strain softened” so that it acquired a*  
44 *progressively decreasing yield strength or higher ductility as the progress of deformation led to*  
45 *changes in its internal fabric.*”) have inspired researchers for many years afterwards. Ramsay (1980)  
46 linked quantitatively the geometry of shear zones (quantified by  $\varphi$ , which represents the angle  
47 between the normal to the shear plane and the inclination of the strain marker) to the amount of shear  
48 strain  $\gamma$ , through the equation (Fig. 1c):

$$49 \quad \gamma = \tan \varphi.$$

50 This equation effectively served as a pivotal link between data from field analyses, microstructural  
51 observations, experimental rock deformation and numerical and analytical modelling. Therefore, this  
52 quantitative approach laid the foundations for the multidisciplinary analysis of strain localisation  
53 processes at the meso- to micro-scale, leading to the development of the modern concept of (ductile)  
54 shear zone (Fossen and Cavalcante, 2017).

55 Ramsay developed important concepts on shear zone formation from field analysis of small-scale  
56 shear zones of cm-dm thickness within the meta-granitoid Maggia nappe (Central Alps, Fig. 1a) and  
57 the gabbroic rocks of Castell Odair (North Uist) (Ramsay and Graham, 1970). Indeed, meta-granitoid

58 units and granitoid plutons represent ideal natural laboratories for investigations of strain localization  
59 within macroscopically homogeneous and isotropic rocks, providing spectacular field exposures of  
60 shear zones and (incipient) shear zone networks (e.g., Ramsay and Allison, 1979; Ingles et al., 1999;  
61 Pennacchioni, 2005; Pennacchioni and Mancktelow, 2007; Pennacchioni and Zucchi, 2013; Ceccato  
62 and Pennacchioni, 2018).

63 In this paper, we report the recent advancements on the understanding of the microstructural  
64 processes controlling the nucleation and development of small-scale ductile shear zones in granitoid  
65 rocks. Such small-scale shear zones preferentially develop at retrograde metamorphic conditions at  
66 the transition between amphibolite to greenschist facies. However, small-scale shear zones similar to  
67 those described in the Maggia nappe have been also reported from different rock types and  
68 metamorphic conditions (e.g. granulite and eclogite facies in granitoid and gabbroic rocks: Van  
69 Roermund et al., 1978; Austrheim, 1987; Pennacchioni, 1996; Getsinger et al., 2013; Smith et al.,  
70 2015; Macente et al., 2017; Lee et al., 2022). Field analysis has highlighted the role of precursor  
71 brittle structures and compositional heterogeneities in controlling ductile strain localisation (Segall  
72 and Simpson, 1986; Bürgmann and Pollard, 1994; Christiansen and Pollard, 1997; Guermani and  
73 Pennacchioni, 1998; Pennacchioni, 2005; Fousseis et al., 2006; Pennacchioni and Mancktelow, 2007,  
74 2018; Schrank et al., 2008b, Pennacchioni and Zucchi, 2013). Recent analytical and modelling  
75 advancements have provided new (semi-)quantitative insights on the role of microstructural and  
76 metamorphic processes in controlling the evolution of both the internal fabric and the yield strength  
77 of the deforming rock (e.g., Smith et al., 2015; Gardner et al., 2017; Meyer et al., 2016). In particular,  
78 the systematic investigation of deformation microstructures with electron backscatter diffraction  
79 (EBSD) and quantitative image analysis has allowed to constrain the microstructural and  
80 crystallographic signatures of specific deformation mechanisms and metamorphic reactions  
81 controlling strain accommodation (Menegon et al., 2008a,b; Pennacchioni et al., 2010; Kilian et al.,  
82 2011; Oliot et al., 2014; Spruzeniece and Piazzolo, 2015; Ceccato et al., 2017, 2018; Gilgannon et al.,  
83 2017). Similarly, micro-analytical techniques (electron probe micro analyser – EPMA,

84 cathodoluminescence – CL, secondary ion mass spectrometer – SIMS, Fourier-transform infrared  
85 spectroscopy – FTIR; micro-computerized tomography –  $\mu$ CT) coupled to petrological modelling  
86 (phase diagram section calculation, geothermobarometry) allowed the role of reaction processes and  
87 fluids (mainly H<sub>2</sub>O) in controlling the rheology and microstructural evolution to be identified  
88 (Kronenberg and Wolf, 1990; Oliot et al., 2010; Goncalves et al., 2012; Bestmann and Pennacchioni,  
89 2015; Finch et al., 2016; Macente et al., 2017; Ceccato et al., 2020, 2022; Kronenberg et al., 2020).  
90 In addition, numerical and analytical models allowed to test inferences retrieved from field and  
91 microstructural analyses and to quantify the effect of deformation-metamorphism-metasomatism  
92 feedback processes on the geometry and rheology of shear zones (Mancktelow, 2002; Shrank et al.,  
93 2008; Huet et al., 2014; Ceccato et al., 2018; Kaatz et al., 2021). These new insights questioned  
94 existing models of purely strain-driven shear zone development (Ramsay and Graham, 1970;  
95 Ramsay, 1980; Lamouroux et al., 1994; Means, 1995). Indeed, recent studies have highlighted the  
96 role of both inherited and dynamically evolving (strain-dependent) gradients of fluid-rock  
97 interactions in controlling shear zone nucleation and development (Oliot et al., 2014; Pennacchioni  
98 and Mancktelow, 2018; Kaatz et al., 2021).

99 In this paper, starting from the concepts developed by John Ramsay and co-authors in the 1970', we  
100 provide a brief overview and an integrated critical discussion of the improvements that innovative  
101 approaches have brought to our understanding of the mesostructural, microstructural and petrological  
102 processes controlling the development of small-scale shear zones in granitoid rocks. This work aims  
103 to complement and expand the review work of Pennacchioni and Mancktelow (2018) on small-scale  
104 shear zones developed in granitoid rocks. For what regards the development of multiscale shear zones  
105 in general, the reader is addressed to more comprehensive and extensive reviews (e.g., Fossen and  
106 Cavalcante, 2017). Finally, we highlight several topics that still require further research efforts and  
107 provide some hints for further research development.

## 109        **2. Field analyses of the meso-scale geometry of shear zones in granitoid rocks**

110    The geometrical model of the ductile shear zone proposed by Ramsay and Graham (1970) allows to  
111    quantify the amount of shear strain accommodated across a shear zone from the deformation  
112    geometry (Fig. 1c). The model was based on the field analysis of small-scale (cm-dm thick) ductile  
113    shear zones developed within isotropic, homogeneous rocks (Fig 1a). In addition, it formalized the  
114    occurrence of “strain markers” in shear zones, such as the inclined sigmoidal foliation, the gradual  
115    deflection of crosscutting geological structures (e.g., dykes) and passive distortion of initially  
116    subrounded, suitably sized objects (Ramsay and Allison, 1979; Ramsay, 1980). In particular, the  
117    quantitative geometrical description and recognition of strain markers allowed the acknowledgement  
118    and quantification of the dominant role of a simple shear component of deformation, over pure shear,  
119    during the development of ductile shear zones (Ramsay and Graham, 1970; Ramsay, 1980). Since  
120    then, the analysis of the field geometry of small-scale shear zones and meso-scale networks led to  
121    significant improvement in our understanding of, as well as opened new debates on: (1) the process  
122    of shear zone nucleation; (2) the strain history recorded with increasing displacement; (3) the  
123    propagation and networking of shear zones.

### 124        *2.1. Nucleation and geometry of incipient shear zones.*

125    Contrasting models on shear zone nucleation have been proposed. Based on theoretical, experimental,  
126    microstructural and numerical studies, shear zones nucleate “spontaneously” as the result of either  
127    thermal softening at large (km) scale (Fleitout and Froidevaux, 1980; Regenauer-Lieb and Yuen,  
128    2003; Duretz et al., 2014; Kiss et al., 2019), or the interconnection of zones of concentrated high  
129    strain caused by micro-scale, point-like flaws in an otherwise isotropic, homogeneous medium (Ingles  
130    et al., 1999; Mancktelow, 2002; Holyoke and Tullis, 2006; Meyer et al., 2016). In contrast, field  
131    observations on granitoid rocks provide increasing evidence for the necessary occurrence of a tabular,  
132    meso-scale precursor structure (either a brittle fracture or a compositional/structural heterogeneity)  
133    to allow a shear zone to nucleate (Mancktelow and Pennacchioni, 2013; Wehrens et al., 2017;  
134    Pennacchioni and Mancktelow, 2018, and reference therein). It is worth noting that brittle fractures

135 can be sometimes associated with pseudotachylytes (frictional melt produced by seismic slip along  
136 faults in silicate rocks: Sibson, 1975), which are then exploited as precursors to small-scale ductile  
137 shear zones in granitoids (e.g., Passchier, 1982; Takagi et al., 2000; Sullivan et al., 2013).

138 Furthermore, field observations suggest the likely occurrence of a threshold dimension of the  
139 precursor structures at the meter to tens of meters scale, below which precursor structures are not  
140 exploited as site for ductile strain localization in meta-granitoids (Mancktelow and Pennacchioni,  
141 2013) as well as in foliated metamorphic rocks (Fusseis and Handy, 2008). The geometry of the  
142 precursor has a two-fold effect on the finite strain geometry of the shear zone (Mancktelow and  
143 Pennacchioni, 2005; Pennacchioni, 2005): (i) it determines the spatial extent (length) of the finite  
144 shear zone, and (ii) it controls the geometry (thickness, continuous vs. discontinuous) of the finite  
145 strain gradient. These observations have had important implications for the understanding of the  
146 processes controlling the growth of shear zones in both thickness and length, as well as networking,  
147 during strain accommodation.

## 148 *2.2. Finite strain gradients and shear zone history*

149 The symmetric, continuous decrease of shear strain from the core to the boundary represented by the  
150 sigmoidal foliation was interpreted to reflect the evolution through time and space of the rheology  
151 during strain localisation at constant temperature and pressure conditions (Ramsay and Graham,  
152 1970; Schrank et al., 2008a,b). This interpretation has led previous authors to the formulation of some  
153 idealised models of shear zone, either narrowing or thickening with increasingly accommodated  
154 displacement (Type I-II-III-IV) (Hull, 1988; Means, 1995; Schrank et al., 2008b; Vitale and Mazzoli,  
155 2008; see Fossen and Cavalcante, 2017 for a detailed discussion of such models). Such models  
156 generally include three different classes of shear zones (e.g., Means, 1995): (i) thickening shear zones  
157 in which the width increases with increasing displacement, usually related to strain-hardening  
158 rheology within the shear zone; (ii) narrowing shear zones in which the width decreases with  
159 increasing displacement, usually related to the occurrence of strain-softening rheology within the

160 shear zone; (iii) constant-width shear zones in which the width remains constant with increasing  
161 displacement. However, so far, no systematic and general relationship between accommodated  
162 displacement and thickness of ductile shear zone has been found (e.g., Fossen and Cavalcante, 2017).  
163 The datasets of thickness-displacement data are scattered and mainly limited to the small-scale ( $\leq 10$   
164 m thick). These scattered data generally describe a not-better-defined trend of thickening with  
165 increasing accommodated displacement (Fig. 2; Fossen and Cavalcante, 2017). However, the concept  
166 itself of shear zones thickening contrasts with the concept of strain localisation *within* the shear zone,  
167 which would rather promote shear zone narrowing (Ramsay and Graham, 1970; Fossen and  
168 Cavalcante, 2017; Pennacchioni and Mancktelow, 2018).

169 In some cases, the finite thickness and the discontinuous/continuous character of the finite strain  
170 gradient might reflect the geometry of the compositional and rheological heterogeneity on which the  
171 shear zone nucleated (Mancktelow and Pennacchioni, 2005; Schrank et al., 2008b; Spruzeniec and  
172 Piazzolo, 2015; Pennacchioni and Mancktelow, 2018; Kaatz et al., 2021). This interpretation is  
173 supported by thickness vs. displacement data retrieved from detailed field mapping, showing no  
174 systematic trends in shear zone thickness variation with (small,  $<1-10$  m) displacement (Pennacchioni  
175 and Mancktelow, 2018). At larger (m to tens of meter) widths, and at larger shear strains, strain  
176 gradients around shear zones are commonly associated with an effective gradient in fluid-rock  
177 interaction and mass transfer (Marquer et al., 1985; Fourcade et al., 1989; Tobisch et al., 1991; Früh-  
178 Green, 1994; Barnes et al., 2004; Keller et al., 2004; Oliot et al., 2014; Finch et al., 2016; Kaatz et  
179 al., 2021). These multiscale observations suggest that the analysis of the development and the role of  
180 such fluid gradients may be key to improve our understanding of the evolution of shear zone thickness  
181 during shearing. Indeed, recent theoretical and numerical models considering fluid expulsion related  
182 to chemo-mechanical pressure gradients induced by rheological and chemical contrasts during shear  
183 zone development reproduce and nicely explain the pattern of shear zone evolution and networking  
184 inferred from field observations of multiscale shear zone networks in crystalline rocks in general  
185 (e.g., Oliot et al., 2014; Finch et al., 2016; Kaatz et al., 2021). Therefore, even though widely accepted



186 and adopted to interpret shear zones geometry and rheology, the standard classification of thickening-  
187 narrowing-constant thickness shear zones might oversimplify the complex (and possibly cyclical)  
188 thickening-narrowing and rheological history recorded in shear zones during displacement  
189 accommodation (e.g., Steffen et al., 2001; Schrank et al., 2008b).

### 190 *2.3. Insights into shear zone growth and networking*

191 Granitoid rocks commonly display well-exposed outcrops of multi-scale shear zone networks,  
192 presenting shear zones ranging from  $10^{-4}$  m to  $10^5$  m in length (Fossen and Cavalcante, 2017). Despite  
193 the overall agreement on the theoretical process of formation of interconnected shear zone network  
194 through the coalescence of single, isolated shear zones, the geologic process controlling linkage and  
195 actual coalescence remains matter of debate (Carreras et al., 2010; Fossen and Cavalcante, 2017). A  
196 network of pre-existing compositional and structural heterogeneities within an otherwise  
197 homogeneous granitoid rock could provide the seed on which a shear zone network develops,  
198 controlling the incipient stages of shear zone formation and networking (Pennacchioni and  
199 Mancktelow, 2007). However, the mechanisms controlling the formation of large scale, complex  
200 networks formed by shear zones of different width and length remains matter of debate.

201 Especially when shear zones nucleate from brittle precursors, contrasting models have been proposed  
202 to explain the transition from isolated, small-scale shears to meso-scale interconnected networks  
203 (Fusseis et al., 2006; Pennacchioni and Mancktelow, 2007; Fusseis and Handy, 2008; Schrank et al.,  
204 2008). Detailed field mapping of incipient meso-scale networks in meta-granitoids has shown that, at  
205 least for small values of accommodated shear strain, the shear zone length is limited to the length of  
206 the precursor structure. In the latter case, the shear zone did not appear to extend in length with  
207 increasingly accommodated strain, as it has been inferred from the preservation of contractional jogs  
208 at the tip of sheared precursor structures (Pennacchioni, 2005; Pennacchioni and Mancktelow, 2007;  
209 Pennacchioni and Zucchi, 2013). On the other hand, based on microstructural and meso-scale  
210 observations in foliated meta-pelites of Cap de Creus, Fusseis et al. (2006) proposed that small-scale

211 shear zones can grow laterally following the synchronous propagation of the migrating brittle  
212 precursor tip parallel to the shear plane in the shear direction. Reaction-induced softening of the host  
213 rock around the shear zones, triggered by fluid diffusion during shearing, may provide another  
214 feasible way to extend and widen single shears which will invariably connect into a complex network  
215 with increasing strain (e.g., Oliot et al., 2014; Kaatz et al., 2021). Therefore, several processes might  
216 be responsible for shear zone networking, and both metamorphic, ductile and brittle-ductile processes  
217 need to be considered to understand the processes of shear zone growth in time and space (e.g.,  
218 Schrank et al., 2008b, Nevitt et al., 2014; Nevitt and Pollard, 2017; Nardini et al., 2018; Döhmann et  
219 al., 2019).

### 221 3. Petrological constraints on shear zone development in granitoid rocks

222 Shear zones are the loci of numerous interactions between deformation, metamorphism and fluid  
223 flow. For instance, the development of shear zones in granitoid rocks is concomitant with  
224 mineralogical changes that, in most cases, favour strain softening and localisation (McCaig, 1984;  
225 Gueydan et al., 2003; Holyoke and Tullis, 2006). In turn, deformation enhances metamorphic  
226 reactions either by overcoming kinetic barriers or by favouring the development of permeability and  
227 fluid infiltration (Fusseis et al., 2009; Oliot et al., 2014). Understanding the feedbacks between all  
228 these interdependent processes is critical to better model strain localisation.

229 In general, metamorphic reactions may be driven by: (i) change in pressure ( $P$ )-temperature ( $T$ )  
230 conditions, (ii) change in fluid content and composition ( $fl$ ), (iii) change in bulk chemistry ( $X$ ) due to  
231 mass transfer (i.e. metasomatism). During deformation of a granitoid rock, it is very likely that the  
232 magmatic assemblage is metastable at the  $P$ - $T$ - $fl$ - $X$  conditions of the deformation (Goncalves et al.,  
233 2012). In this case, metamorphic and metasomatic reactions are also driven by (iv) disequilibrium  
234 and the equilibration of the metastable magmatic assemblage at the  $P$ - $T$ - $fl$ - $X$  conditions of the  
235 deformation (Goncalves et al., 2012). In the following section, we review phase relations of granitoid

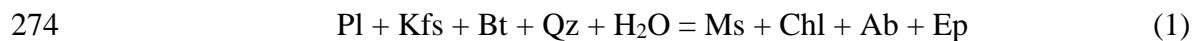
236 rocks as a function of  $P$ - $T$ - $f$ - $X$  conditions and discuss their influence on strain localisation and shear  
237 zone development through the analysis of  $P$ - $T$ ,  $T$ - $M_{H_2O}$  and  $T$ - $X$  phase diagram sections. Phase  
238 diagrams were computed with *Perple\_X* 6.9.1 (Connolly, 2005) for the bulk composition of an  
239 average meta-granodiorite from the Neves Lake area (e.g., Leydier et al., 2019), using the  
240 thermodynamic database of (Holland and Powell, 2011; TC-DS62) and the a-X solid solution models  
241 of White et al. (2014). The chemical system includes  $K_2O$ - $FeO$ - $MgO$ - $Al_2O_3$ - $SiO_2$ - $H_2O$  (KFMASH).  
242  $MnO$ ,  $TiO_2$  and  $Fe_2O_3$  are not considered here. Mineral abbreviations are from Whitney and Evans  
243 (2010).

### 244 *3.1. $P$ - $T$ conditions of shear zones in meta-granitoids and granitoid plutons*

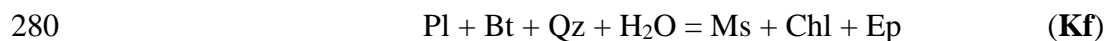
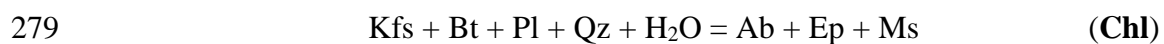
245 Since the description of small-scale shear zones in the meta-granitoid unit of the Maggia Nappe of  
246 Ramsay and Graham (1970), many other occurrences of similar small-scale shear zones were reported  
247 from the crystalline basements and granitoid plutons of the Alps and other localities worldwide  
248 (Pennacchioni and Mancktelow, 2018, and reference therein). Despite the different tectonic and  
249 metamorphic history of (Alpine) meta-granitoid and magmatic plutons, all the reported cases of  
250 small-scale shear zones developed between 350 °C and 600 °C, during cooling of the pluton or during  
251 the retrograde  $P$ - $T$  path of the meta-granitoid unit (Fig. 3a). In detail, shear zone nucleation and  
252 development in meta-granitoid units occurred mainly at  $P > 0.4$  GPa and  $T = 420$ - $600$  °C (light blue  
253 areas in Fig. 3a). In granitoid plutons, emplaced at middle and upper crustal levels, shear zone  
254 nucleation occurred at  $P < 0.5$  GPa and  $T = 350$ - $550$  °C (pale orange areas in Fig. 3a), similarly to  
255 what was already reported by Gapais (1989). Overall, these data show that small-scale shear zones in  
256 meta-granitoids preferentially develop in a restricted  $T$  window at the transition between amphibolite  
257 and greenschist facies. The main differences between shear zones developed in meta-granitoid units  
258 and those from cooling granitoid plutons probably reside in (i) the timescale of deformation (much  
259 shorter for cooling plutons) and (ii) fluid (mainly  $H_2O$ ) availability, which has enormous effects for  
260 the rheology, metamorphic processes and longevity of shear zones (see sections below).

261 3.2. *P-T phase relations of granitoid systems under H<sub>2</sub>O-saturated conditions*

262 We explore, through phase relation modelling, the main mineralogical changes expected at the  
263 conditions of shear zone development as defined above. Figure 3b-c shows the *P-T* phase diagram  
264 sections computed for the KFMASH system, which is suitable to model the main phase relations in  
265 granitoid systems. Here, we firstly consider phase relations calculated at H<sub>2</sub>O-saturated conditions  
266 for the sake of simplicity, and then we discuss the effect of H<sub>2</sub>O-undersaturation on phase relations  
267 in granitoid rocks in the next Sections (see below). Figure 3b shows that from ca. 700 °C (i.e. the  
268 temperature of the hydrated solidus) to about 550 °C and a pressure from 0.1 to 0.8 GPa, there is a  
269 large multivariant field that consists of Pl + Bt + Qz ± Amp ± Kfs corresponding to the expected  
270 magmatic assemblage for a granodiorite. During cooling of a pluton or a retrograde path, there are no  
271 major mineralogical changes from 700 to 550 °C. In contrast, between 450 and 500 °C, the phase  
272 diagram shows several phase boundaries that are mostly related to a main univariant reaction (1) in  
273 the KFMASH system:



275 that separates amphibolite and greenschist facies. For the average granodiorite bulk composition,  
276 there are four divariant reactions associated with reaction (1). Among them, we discuss two divariant  
277 reactions located at pressure above 0.5 GPa that account for the main mineralogical changes observed  
278 in granitoid shear zones:



281 These two divariant reactions occur in a narrow temperature window, between 450 °C and 500 °C  
282 (Fig. 3c), and they can be linked with classic petrological observations made in shear zones: (a) the  
283 saussuritisation of magmatic plagioclase into a fine grained assemblage consisting of albite, epidote  
284 and muscovite (e.g., (Fitz Gerald and Stünitz, 1993; Stünitz and Fitz Gerald, 1993; Oliot et al., 2010;

285 Goncalves et al., 2012; Menegon et al., 2008a) (b) the breakdown of K-feldspar to produce muscovite  
286 (Hippertt, 1998; Gueydan et al., 2003; Tschegg and Grasmann, 2009); and (c) the formation of  
287 muscovite and chlorite at the expense of biotite (Goncalves et al., 2012; Airaghi et al., 2020).  
288 Additionally, myrmekite formation at the expense of K-feldspar is another commonly reported  
289 reaction which contributes to the development of shear zones in granitoid rocks in this temperature  
290 window (Tsurumi et al., 2003; Menegon et al., 2006; Ceccato et al., 2018; Cisneros-Lazaro et al.,  
291 2019).

292 To conclude, the breakdown of the magmatic load-bearing framework (LBF; Handy, 1990) of Pl +  
293 Qz  $\pm$  Kfs (almost ~85% of the rock volume) into a fine-grained assemblage consisting of Ab + Ms +  
294 Ep  $\pm$  Chl is a first order softening mechanism that may favour strain localisation. Thus, in granitoid  
295 rock, the 450-500 °C temperature range corresponds to the most favourable conditions for the  
296 initiation of reactions promoting strain localisation. The identified temperatures provide the  
297 maximum temperatures at which such reactions may occur. Any occurrence of strain localization and  
298 shear zone formation below these temperatures might be effectively controlled by the same  
299 parageneses.

### 300 3.3. Fluid (H<sub>2</sub>O) saturation and mineralogical effects

301 In the previous section, the phase relations have been investigated assuming H<sub>2</sub>O-saturated conditions  
302 during deformation. However, H<sub>2</sub>O-saturation might not even be the case or persist during the whole  
303 history of shear zone development, as it has been suggested by several recent studies (Finch et al.,  
304 2016; Ceccato et al., 2020; Menegon and Fagereng, 2021). Thus, we investigate here the role of H<sub>2</sub>O  
305 content, and more particularly H<sub>2</sub>O undersaturation on phase relations. Figure 3d is a  $T$ - $M_{H_2O}$  phase  
306 diagram computed at a constant pressure of 0.6 GPa. The black dashed curve traces the boundary  
307 between H<sub>2</sub>O-saturated (to the right) and H<sub>2</sub>O-undersaturated (to the left) phase assemblages. The  $T$ -  
308  $M_{H_2O}$  diagram shows that the meta-granodiorite composition is rather sensitive to very little changes  
309 in  $M_{H_2O}$  at undersaturated conditions. Indeed,  $M_{H_2O}$  varies overall from 1.2 to 6 mol (0.33 to 1.63

310 wt%). Conversely, at H<sub>2</sub>O-saturated conditions, the mineral assemblage is mainly dependent on the  
311 temperature. During deformation and cooling from 550 to 400 °C, the amount of H<sub>2</sub>O in mol per  
312 kilogram of rock must increase by more than 160% to keep the assemblage H<sub>2</sub>O-saturated. The main  
313 mineralogical effects of H<sub>2</sub>O-undersaturation are (1) the stabilization of garnet, (2) the increase of  
314 plagioclase stability to lower temperature (485 to 450°C) and (3) the breakdown of chlorite. Under  
315 closed-system conditions (no gain and loss of H<sub>2</sub>O), cooling from 550 °C and deformation would  
316 result in the formation of garnet-bearing, at intermediate degrees of under-saturation, or even kyanite  
317 bearing shear zones for very low (geologically unrealistic) H<sub>2</sub>O contents and low temperature ( $M_{H_2O}$   
318  $< 0.12$  mol,  $T < 425$  °C). The stabilization of garnet and kyanite may be due also to bulk composition  
319 effects (abundant MnO and Fe<sub>2</sub>O<sub>3</sub> for garnet), or extensive mass-transfer at fluid-saturated conditions  
320 (e.g., kyanite-bearing, granitoid-derived micaschists reported by Sassier et al., 2006). Nonetheless,  
321 the occurrence of such index minerals could suggest shear zone formation at H<sub>2</sub>O under-saturated  
322 conditions. At intermediate stages of H<sub>2</sub>O-undersaturation, i.e. in which a fluid is available but not  
323 enough to result in H<sub>2</sub>O as a free-fluid phase, metamorphic parageneses containing progressively  
324 more proportions (in terms of both modal quantity and number of different phases) of hydrated  
325 minerals will be stabilised. Then, after H<sub>2</sub>O saturation, the availability of H<sub>2</sub>O as a free-fluid phase  
326 has enormous effects on the kinetics of metamorphic reactions (having H<sub>2</sub>O as limiting reactant) and  
327 fluid-assisted deformation processes (e.g., dissolution-precipitation and mass transfer, e.g., Yardley  
328 and Valley, 1997; Wintsch and Yi, 2002; Ceccato et al., 2020; Tursi, 2022). Therefore, H<sub>2</sub>O-  
329 saturation conditions directly affects the bulk rheology, and need to be investigated in detail for a  
330 thorough understanding of shear zone formation and evolution.

### 331 *3.4. Fluid conditions during shear zone development*

332 H<sub>2</sub>O-saturation conditions, and thus the occurrence of a free-fluid phase, may be affected by several  
333 parameters of the geological setting in which granitoid shear zones develop. Indeed, initial fluid  
334 availability, the supply rate from any (either internal or external) fluid source, the retrograde cooling  
335 rate of the (meta-)granitoid and the occurrence of retrograde hydration metamorphic reactions may

336 determine the fluid-saturated vs. -undersaturated conditions of the system regardless of the “open vs.  
337 closed” conditions of the system with respect to fluid flow (e.g., Yardley, 1981).

338 A variable fluid behaviour during shear zone development can be inferred from field, microstructural  
339 and petrological evidence. When small scale shear zone in granitoid rocks exploit precursor  
340 structures, common precursors are: (i) mineralized joints; (ii) veins; (iii) alteration haloes due to fluid-  
341 rock interaction along the selvages of brittle fractures and veins (e.g., Mancktelow and Pennacchioni,  
342 2005; Pennacchioni, 2005). This set of precursor types involves the occurrence of (mineralising)  
343 fluids during at least the formation of the precursors, and thus, local fluid-saturated conditions. Many  
344 case studies report that, during precursor formation, the fluid is likely to be of local origin and the  
345 granitoid rock behaves as a closed system (e.g., Cesare et al., 2001; Pennacchioni, 2005;  
346 Mittempergher et al., 2014; Leydier et al., 2019; Ceccato et al., 2020). After precursor nucleation,  
347 phase diagram section calculations have shown that incipient shearing on small-scale shear zones  
348 may occur at both fluid-saturated (e.g., in meta-granitoid rocks of the Neves Lake area, Leydier et al.,  
349 2019) and fluid-undersaturated conditions (e.g., shear zones from the Adamello, Rieserferner and  
350 Sierra Nevada granitoid plutons, Ceccato et al., 2020). In addition, isotopic and major element  
351 chemical analyses suggest that, at least within granitoid plutons, the incipient stages of shearing as  
352 well might occur at closed system conditions (Pennacchioni, 2005; Mittempergher et al., 2014). At  
353 larger-scale, (m to tens of m in width) shear zones in meta-granitoid rocks deform under fluid-  
354 saturated, open-system conditions, commonly accompanied by significant mass transfer (e.g.,  
355 Marquer, 1985; Fourcade et al., 1989; McCaig et a., 1990; Früh-Green, 1994; Barnes et al., 2004;  
356 Rossi et al., 2005; Goncalves et al. 2012; Oliot et al., 2014; Rolland and Rossi, 2016; Wehrens et al.,  
357 2017). At such large scale, given the paucity of connate fluids in (meta-)granitoid rocks, an external  
358 fluid source is needed, as it is commonly reported for other large-scale, retrograde orogenic shear  
359 zones (e.g., Kuckaus shear zone, Diener et al., 2016). As already proposed by Fourcade et al. (1989),  
360 the development of a network of shear zones “well-interconnected” among them, and connected to  
361 an external fluid source, likely promote fluid-flow, open system conditions and significant mass-

362 transfer (see below; e.g., Ollot et al., 2014). However, this topic deserves further investigations and  
363 research.

### 364 *3.5. Mass transfer and metasomatism in shear zones*

365 Shear zones developed under fluid-saturated conditions are also likely affected by mass transfer either  
366 via diffusion or infiltration metasomatism (Korzhinskii, 1970). Metasomatic reactions result in  
367 mineral assemblages that differ significantly from those formed where the system is either open or  
368 closed to only H<sub>2</sub>O transfer. Mass transfer (i.e., metasomatism) has been reported and quantified in  
369 many granitoid shear zones (e.g., Goncalves et al., 2012 and reference therein). Although each shear  
370 zone should be treated as a unique case (e.g., Tursi et al., 2018), granitoid shear zones developed  
371 under greenschist to amphibolite facies conditions show very consistent mass transfer with, in  
372 general, a gain in MgO and H<sub>2</sub>O coupled with a loss of Na<sub>2</sub>O and CaO (Goncalves et al., 2012). It is  
373 beyond the scope of this contribution to explain this coupling, but we explore via phase relation  
374 modelling the effect of this specific mass transfer on the mineralogical composition of granitoid shear  
375 zones. Based on the literature review, we propose a hypothetical mass transfer scenario, with only  
376 four mobile components: gain of H<sub>2</sub>O and MgO coupled with the loss of CaO and Na<sub>2</sub>O. Mass transfer  
377 evolves linearly from the protolith composition to a metasomatic shear zone characterized by +200%  
378 for MgO and -100% and -75% for Na<sub>2</sub>O and CaO respectively. The amount of H<sub>2</sub>O is computed such  
379 that the system is kept H<sub>2</sub>O saturated. The phase diagram is computed at 0.6 GPa.

380 The change in bulk composition induced by mass transfer has no effects on the absolute *P-T* position  
381 of the **(Chl)** and **(Kf)** reactions (Fig. 4a). At 0.6 GPa, these reactions are located at 480-485 °C (Fig.  
382 4a). Under closed system conditions ( $X=0$ , Fig. 4b), no major mineralogical changes occur before  
383 reaching the **(Chl)** and **(Kf)** reactions. At 485-480 °C, 65 vol% of the rock (Pl and Kfs) is  
384 recrystallized into fine-grained aggregates of Ab + Ep (~55%) ± phyllosilicates (Bt + Ms, ~20%).  
385 Figures 4a-4c show the effect of mass transfer at a temperature of 475 °C just below the onset of shear  
386 zone development. Progressive mass transfer induces a continuous breakdown of Ab + Ep to produce



387 large amounts of Ms + Chl + Qz, (Fig. 4c). In a strongly metasomatized system ( $X=1$ ), the resulting  
388 rock in the shear zone would be a phyllonite consisting of 50% of Qz and 50 % of Ms + Chl (Fig.  
389 4c).

390 Therefore, the main result of mass transfer and metasomatic reactions in shear zones is the progressive  
391 replacement of the protolith feldspars building up the LBF (Handy, 1990). This results in softening  
392 of the deforming shear zone (Gueydan et al., 2003; Oliot et al., 2014). The formation of fine-grained  
393 aggregates of Ab + Ep at intermediate degrees of metasomatism first, and then the stabilization of  
394 large amounts of phyllosilicates (Bt + Ms + Chl) at advanced stages, have enormous effects on the  
395 rheology and fabric development of shear zones, leading to progressive softening and foliation  
396 development.

#### 397 **4. Shear zone microstructure and rheology**

398 The “strain softening mechanisms” proposed by Ramsay and Graham (1970) involve the modification  
399 of the spatial arrangement, geometrical characters, chemical composition and phase proportions of  
400 the mineral constituents at the micro- to meso-scale (e.g., Poirier, 1980; White et al., 1980).

##### 401 *4.1. Microstructural processes and rheological effects*

402 At the temperature conditions at which small-scale shear zones are inferred to nucleate and develop  
403 (350-600 °C), strain accommodation and the development of the shear zone fabric involve grain-size  
404 reduction (GSR) of mineral phases (Ramsay and Allison, 1979). GSR in quartz is mainly  
405 accommodated by dynamic recrystallization (Stipp et al., 2002; Bestmann and Pennacchioni, 2015);  
406 whereas feldspars undergo predominantly brittle GSR (Ree et al., 2005; Viegas et al., 2016) in  
407 addition to the retrograde metamorphic reactions described above (Fig. 5a-b). The contribution of  
408 dynamic recrystallisation to GSR in feldspars is negligible in the temperature range considered here  
409 (Gapais, 1989; Fitz Gerald and Stünitz, 1993; Menegon et al., 2008b).

410 Associated with GSR processes, especially when local mass-transfer is involved, phase mixing (PM)  
411 may result from neo-crystallisation and phase nucleation related to metamorphic reactions or to the

412 emergence of synkinematic porosity (Fusseis et al., 2009; Kilian et al., 2011; Menegon et al., 2015;  
413 Ceccato et al., 2018; Gilgannon et al., 2020, 2021) (Fig. 5a-b). The main result of GSR and PM  
414 processes consists in the formation of fine-grained, mixed polyphase aggregates characterised by a  
415 rheology (i.e., deformation mechanisms) “weaker” than that of the parent mineral phases (Kilian et  
416 al., 2011; Platt and Behr, 2011a,b; Sullivan et al., 2013; Wehrens et al., 2016; Mansard et al., 2018;  
417 Stenvall et al., 2019). Indeed, fine-grained, polyphase aggregates preferentially deform by grain-size  
418 sensitive (GSS) creep processes, including grain-boundary sliding (GBS) and diffusion creep, which  
419 are inherently characterised by a “weak” linear-viscous rheology (Fig. 5c). However, the softening  
420 potential of GSR by dynamic recrystallisation of monomineralic aggregates is limited by the  
421 piezometric relationships and the balance between grain size reduction and crystal growth processes  
422 (De Bresser et al., 2001; Cross et al., 2015; Tokle and Hirth, 2021) (Fig. 5d). Effective softening is  
423 mainly related to the combined activation of GSR and PM processes resulting in polyphase aggregates  
424 (Kilian et al., 2011; Spruzeniece and Piazzolo, 2015; Ceccato et al., 2018; Stenvall et al., 2019). Indeed,  
425 the polyphase mixed nature and fine grain size of these aggregates have two major consequences on  
426 the kinetics of rheology-controlling deformation mechanisms (Herwegh et al., 2011): (i). the fine  
427 grain size enhances the efficiency of diffusion-assisted creep processes (Coble creep, dissolution-  
428 precipitation creep) through the reduction of the diffusion length scales; (ii) PM hinders the efficiency  
429 of crystal growth by pinning, maintaining the overall grain size of the polyphase aggregates in the  
430 range where diffusion-assisted GSS creep is dominant (Fig. 5d). Therefore, GSR and PM are the key  
431 processes steering the evolution of shear zone strain-softened rheology, as well as of the shear zone  
432 micro- and meso-scale fabric. Whilst the strength of mono-phase aggregates can be estimated by  
433 paleopiezometers for dislocation creep (e.g., for quartz, Cross et al., 2017), the strength of polyphase  
434 mixed aggregates is still poorly constrained and requires the adoption of analytical mixing models  
435 under significant assumptions (Huet et al., 2014; Platt, 2015). Nevertheless, rheological mixing  
436 models can be very effective in the quantification of softening related to PM processes (e.g., Fig. 5c;  
437 Huet et al., 2014; Gerbi et al., 2016; Ceccato et al., 2018).

438      *4.2.Fabric and microstructural development*

439 Progressive GSR with increasing strain and metamorphic reactions lead to the development of fine-  
440 grained, recrystallized aggregates. These aggregates progressively merge, with increasing strain, to  
441 form the typical strain-dependent foliation observed along the ductile shear zone walls, leading to the  
442 transition from protomylonite, to mylonite and ultramylonite (Ramsay and Allison, 1979; Ramsay,  
443 1980; White et al., 1980) (Fig. 1c). In particular, the development of the characteristic sigmoidal  
444 foliation of protomylonites is primarily controlled by the formation of a phyllosilicate-dominated  
445 interconnected weak layer (IWL) microstructure during the incipient stages of deformation (e.g.,  
446 Handy, 1990; Holyoke and Tullis, 2006). With increasing strain, this phyllosilicate dominated IWL  
447 is however overtaken and disrupted by the development of an IWL microstructure formed by  
448 coalesced fine-grained, polyphase aggregates resulting from feldspars breakdown and neo-  
449 crystallisation (Hippertt and Hongn, 1998; Stenvall et al., 2019). As strain accommodation,  
450 coalescence and phase mixing progress, fine-grained homogeneous polyphase ultramylonites form,  
451 representing the final product of strain accommodation and rheological softening (Ramsay and  
452 Allison, 1979; Kilian et al., 2011; Cross and Skemer, 2017; Stenvall et al., 2019) (Fig. 6a). Thus, the  
453 strength and textural evolution of ductile shear zones strictly depend on the fate of feldspar phases,  
454 forming the LBF in the protolith (Handy, 1990) (Fig. 6a). As described above, the dismantling of the  
455 feldspars-LBF is deeply affected by the H<sub>2</sub>O-saturation conditions of the system and the occurrence  
456 of mass-transfer processes, which in turn affect the final ultramylonitic mineral assemblage.

457      *4.3.Deformation mechanisms and fluids*

458 Aqueous fluids are agents of chemical (hydrolytic) softening in nominally anhydrous minerals (e.g.,  
459 Kronenberg and Wolf, 1990; Bras et al., 2021; Ceccato et al., 2022). Similarly, the occurrence of a  
460 fluid phase may enhance the kinetics of retrograde metamorphic reactions aiding strain localisation,  
461 then affecting the shear zone bulk rheology as well (e.g., Spruzeniece and Piazzolo, 2015; Stenvall et  
462 al., 2020). Consequently, the evolution of both the bulk fluid content of shear zone and the  
463 intracrystalline fluid content of mineral phases have major effects on the geometrical and rheological

464 evolution of shear zones (Spruzeniece and Piazzolo, 2015; Finch et al., 2016; Stenvall et al., 2020;  
465 Kaatz et al., 2021).

466 In turn, the microstructural and rheological evolution of shear zone with increasing strain may affect  
467 both the bulk fluid content of the shear zone and of its mineral constituents (Bestmann and  
468 Pennacchioni, 2015; Finch et al., 2016). Indeed, fluid content may evolve with increasing strain as an  
469 effect of feedback processes between deformation (GSR-PM), and metamorphic-metasomatic  
470 processes (Hobbs et al., 2010; Oliot et al., 2014; Finch et al., 2016; Kaatz et al., 2021). Second-order  
471 processes related to GSR may lead to both the release of intracrystalline fluids (e.g. fluid inclusions  
472 and intracrystalline H<sub>2</sub>O) into the shear zone, as well as to the drying up of the shear zones during  
473 recrystallisation as a consequence of increased grain boundary area (Pennacchioni and Cesare, 1997;  
474 Mancktelow and Pennacchioni, 2004; Bestmann and Pennacchioni, 2015). Furthermore, the  
475 activation of GSS creep and PM processes has a two-fold effect on fluid content (Fig. 6b): (i) they  
476 promote fluid flow toward the most strained portions of the shear zone through the development of  
477 local grain-scale dilatancy (i.e., creep cavitation; Füsseis et al., 2009; Menegon et al., 2015;  
478 Spruzeniece and Piazzolo, 2015; Gilgannon et al., 2021); (ii) at the same time, the rheological contrast  
479 between the weak shear zone and the undeformed host rock may lead to pressure gradients driving  
480 fluids away from the shear zone toward the host rock (Mancktelow, 2006; Oliot et al., 2014; Finch et  
481 al., 2016; Menegon and Fagereng, 2021). The relevance of all these processes may change depending  
482 on the open vs. closed conditions of the system with respect to fluid flow, and upon the stabilization  
483 of a stable retrograde metamorphic paragenesis (Yardley and Valley, 1997; Diener et al., 2016;  
484 Stenvall et al., 2020). The rheology and geometry evolution of shear zones are therefore the results  
485 of a delicate equilibrium between fluid availability and deformation-metamorphic-metasomatic  
486 feedback processes at the microscale.

487 **5. Final remarks and future directions**

488 The analysis of meta-granitoid rocks provided very helpful inferences for the understanding of  
489 (incipient) strain localization processes in isotropic geological bodies at mid crustal conditions.  
490 Granitoid units, however, represent the exception rather than the rule for what regards the typical  
491 characters of metamorphic units in the continental crust and orogenic belts. The occurrence of  
492 pervasive heterogeneities (e.g., foliations) in metamorphic orogenic units, and the composition of  
493 such units themselves, might deeply affect the development of shear zones at general mid-crustal  
494 conditions (Rennie et al., 2013; Diener et al., 2016).

495 Small-scale shear zones in granitoid rocks predominantly nucleate exploiting precursor brittle  
496 structures and compositional heterogeneities. They nucleate and develop within a restricted  
497 temperature window (350-600°C), during the retrograde or cooling path from amphibolite to  
498 greenschist facies metamorphic conditions. At these conditions, the meta-stable granitoid mineral  
499 assemblage is affected by a series of retrograde divariant reactions that assist strain accommodation  
500 within the shear zones. These reactions are spatially restricted, i.e. localised, within shear zones (and  
501 not in the host rock in general), suggesting that metamorphic retrograde equilibration of the  
502 metastable assemblage is promoted by either fluid-rock interaction or deformation processes taking  
503 place within the shear zones.

504 These reactions mainly involve breakdown of feldspar(s) with the resulting development of either  
505 fine-grained polyphase or phyllosilicate-rich aggregates. Thus, feldspar breakdown at retrograde  
506 conditions controls the evolution of both the rheology (through the activation of GSS creep  
507 mechanisms) and the fabric (through the development of phyllosilicate- or polyphase-aggregate-  
508 bearing IWL) of shear zones. Indeed, shear zone rheology is controlled by the activation of GSS creep  
509 and PM processes, which result in the strain softening rheology inferred by Ramsay and Graham  
510 (1970). Such microstructural processes have a major role in controlling the diffusion of fluids and the

511 extent of metasomatic and retrograde metamorphic processes leading to host rock softening and shear  
512 zone widening.

513 Notwithstanding the improvements provided by detailed field, analytical and numerical studies,  
514 several points remain unclear and deserve further analyses. We think that three main topics should be  
515 at the core of future research on shear zone development and strain localisation in both granitoid rocks  
516 and the continental crust in general, including: (1) quantification of strain evolution and rheology; (2)  
517 understanding transient vs. steady state evolution; (3) integrated, multiscale shear zone networking  
518 models.

519 The first aspect consists in the quantification of rates and magnitudes (e.g., progression of  
520 recrystallisation with strain) of deformation processes in natural shear zones, in order to be then  
521 compared to the results of experimental rock deformation and numerical modelling. Field,  
522 microstructural and petrological integrated analyses should aim at the quantification of the effects of  
523 the feedbacks between deformation, metamorphism and metasomatism on the rates and intensity of  
524 strain localisation and on the strength of ductile shear zones. The integration of rheological mixing  
525 models with quantitative microstructural observations along strain/fluid gradients of small-scale shear  
526 zones will allow to quantify the effects of specific deformation, metamorphic and metasomatic  
527 processes on the bulk rheology of shear zones as a function of strain and of different degrees of fluid-  
528 rock interaction (Shigematsu, 1999; Gerbi et al., 2010; Cross and Skemer, 2019; Whyte et al., 2021).  
529 Both strain and fluid-gradients develop through time, and thus their analysis may provide  
530 fundamental insights into the evolution in time of microstructural processes, feedback mechanisms  
531 and rheology.

532 In addition, further investigations are necessary to understand the strict relationship between strain  
533 accommodation and the gradients of fluid-rock interaction intensity observed at various scale. Is there  
534 an effective transition from closed to open system with increasing strain and networking (Fourcade  
535 et al., 1989, Oliot et al. 2014)? What is the actual process controlling this transition?

536 These insights will in turn shed new light on other relevant questions: (i) what are the factors  
537 controlling the role of ductile shear zone as fluid pathways?, and (ii), what controls the evolution of  
538 the equilibrium between fluid intake and expulsion? Transient vs. continuous fluid intake and fluid  
539 flow is then reflected on the softening mechanisms controlling strain accommodation through time.  
540 This will have significant implications for the understanding of transient vs. steady-state character of  
541 deformation. Therefore, the final questions become: is strain accommodation in shear zones  
542 continuous or discontinuous (i.e., transient vs. steady state deformation)? What are the parameters  
543 controlling transient vs. steady state deformation? Further improvements on geochronological dating  
544 techniques may provide the necessary tools to address this question (Oriolo et al., 2018).

545 Considering the nucleation of shear zones on precursor structures, what are the boundary conditions  
546 that drive shear zone widening and networking? Do small-scale shear zones effectively represent the  
547 seeds of larger scale mylonitic belts? To what scale do brittle and compositional precursors control  
548 the nucleation of shear zones? The existing models for shear zone growth and networking still  
549 partially fail to integrate field, microstructural and petrological observations consistently. So far,  
550 thickness-length-displacement data retrieved from detailed mapping of networks of meso-scale shear  
551 zones do not show any systematic trend in shear zone thickness variation, neither at small-scale, nor  
552 connecting the small-scale toward the large scale (Fossen and Cavalcante, 2017; Pennacchioni and  
553 Mancktelow, 2018).

554 Recent results from field, microstructural and numerical analyses have highlighted the effects of  
555 scale-dependent feedback mechanisms in controlling the geometry and the networking of general  
556 shear zones in space and time (Regenauer-Lieb and Yuen, 2003; Schrank et al., 2008b; Hobbs et al.,  
557 2010, 2011; Mancktelow and Pennacchioni, 2013; Oliot et al., 2014; Finch et al., 2016; Meyer et al.,  
558 2016; Pennacchioni and Mancktelow, 2018; Kiss et al., 2019; Bras et al., 2021; Kaatz et al., 2021).  
559 Such feedback mechanisms largely originate from the interaction between deformation, mass  
560 transfer, chemical and heat diffusion (i.e. metamorphic and metasomatic) processes, which might take

561 place at different length- and timescales during shear zone evolution (Fig. 2; e.g., Regenauer-Lieb  
562 and Yuen, 2003; Hobbs et al., 2010; Kaatz et al., 2021). Chemical diffusion and (metamorphic)  
563 reactions seem to control the thickness of small- and meso-scale shear zones. For example, Oliot et  
564 al. (2014) proposed a model for shear zone thickening and networking driven by fluid expulsion  
565 (meso-scale mass transfer) from shear zones, triggered by chemically and mechanically induced  
566 pressure gradients, with consequent metasomatic-softening reactions in the host rock. At larger scales  
567 ( $>10^3$  m), shear heating may contribute significantly to the softening of the host rock leading to shear  
568 zone thickening (e.g., Kiss et al., 2019). The main result of such feedback mechanisms is the softening  
569 of the host rock, which in turn would lead to shear zone thickening. At first sight, these softening-  
570 driven model of shear zone thickening might seem to dispute the classical model of softening-  
571 narrowing shear zones (e.g., Means, 1995). However, rather than considering softening within the  
572 shear zone itself as the classical model does, such model expect softening in the host rock along the  
573 selvages of the shear zone and shear zone tips, allowing more rock to be included in the deforming  
574 shear zone during softening (Oliot et al., 2014).

575 During shear zone thickening and networking, the deforming systems may be subjected to different  
576 feedback processes controlling both the deformation mechanisms and the rate at which strain is  
577 accommodated (Schrank et al., 2008b; Hobbs et al., 2010; Meyer et al., 2016; Bras et al., 2021).  
578 Therefore, we must consider these feedback mechanisms in the future studies that seek to explore the  
579 processes of networking and mesoscale evolution, integrating field, microstructural and petrological  
580 observations. These studies will need to be supported by an extended database of displacement vs.  
581 thickness data retrieved from quantitative, multiscale mapping of shear zone networks on optimally  
582 exposed outcrops. Accordingly, numerical models should aim at comparing their results with such  
583 databases provided by field analyses.

584



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592

593 **Figure captions**

594 **Figure 1. (a)** Example of small-scale shear zone from the Ramsay's type locality close to Poncione  
595 dei Laghetti (Maggia nappe, Central Alps). Courtesy of Neil Mancktelow. **(b)** Example of small-scale  
596 shear zone from the Neves Lake area (Tauern Window, Eastern Alps). Courtesy of Giorgio  
597 Pennacchioni. **(c)** Schematic representation of the geometrical model of a Ramsay-type shear zone,  
598 showing the transition, developed along the strain gradient identified by the sigmoidal foliation  
599 (marked by grey dashed curves), from the undeformed granitoid host rock (HR), to protomylonite  
600 (Proto), mylonite (Myl) and ultramylonite (Umyl). The model presents the geometrical parameters  
601 ( $\alpha$ ,  $\alpha'$ ,  $\varphi$ ,  $\theta$ ) and equations adopted to calculate shear strain  $\gamma$ . The angles  $\alpha$  and  $\alpha'$  represent the angle  
602 between the shear plane and the orientation of a passive strain marker (e.g., an aplitic dyke, dark  
603 grey). Redrawn and modified after Ramsay (1980). T: thickness, D: displacement.

604 **Figure 2.** Thickness vs. Displacement diagram, redrawn from Fossen and Cavalcante (2017)  
605 including the data reported in Hull (1988), Takagi et al. (2000) and Pennacchioni and Mancktelow  
606 (2018). The reader is referred to each of these publications for further details on the data. The  
607 transparent arrows suggest the likely Thickness-Displacement trends described by some datasets from  
608 different small-scale shear zone networks. At the bottom, the length scales for different diffusion  
609 (heat, fluids, mass) processes are reported (from Hobbs et al., 2011). The inset shows the theoretical  
610 Thickness-Displacement trends with the inferred softening-hardening rheology as proposed by the  
611 models of Means (1995).

612 **Figure 3.** *P-T-fl* conditions of shear zones development in the Alpine (and neighbouring) meta-  
613 granitoid units and granitoid plutons. **(a)** *P-T* diagram showing the estimated metamorphic conditions  
614 of shear zone nucleation in meta-granitoid (blue rectangles) and granitoid plutons (orange rectangles)  
615 from selected case studies. Legend: (1) (Ramsay and Allison, 1979); (2) Mont Blanc (Guermani and  
616 Pennacchioni, 1998); (3a) Neves shear zones, regional metamorphism (Mancktelow and

617 Pennacchioni, 2005); (3b) Neves shear zones, nucleation of precursor (Leydier et al., 2019); (4)  
618 Gotthard massif – Fibbia meta-granite (Oliot et al., 2010, 2014); (5a) Gran Paradiso, regional  
619 metamorphism (Menegon and Pennacchioni, 2010); (5b-c) Gran Paradiso, single samples of localised  
620 shear zones (samples GPP2 and GPP5, respectively, from Rosenbaum et al., 2012); (6) Aar massif –  
621 Grimsel granodiorite (Goncalves et al., 2012; Wehrens et al., 2016); (7) Suretta nappe – Roffna  
622 rhyolite (Goncalves et al., 2016); (8a) Adamello and Sierra Nevada shear zones (Ceccato et al., 2020);  
623 (8b) Sierra Nevada (Nevitt et al., 2017); (9) Serifos granodiorite (Tschegg and Grasemann, 2009);  
624 (10) Capanne pluton, Elba Island (Smith et al., 2011); (11) Bergell pluton (Guastoni et al., 2014);  
625 (12a) Rieserferner mylonites on Ep-veins (Ceccato et al., 2020); (12b) Rieserferner high-temperature  
626 sheared joints and mylonites (Ceccato and Pennacchioni, 2018). Mineral abbreviations from Whitney  
627 and Evans (2010). **(b)**  $P$ - $T$  phase diagram section computed for an average meta-granodiorite  
628 composition from the Neves Lake area, Eastern Alps (see Leydier et al., 2019). The dashed rectangle  
629 represents the  $P$ - $T$  conditions area shown in detail in (c). **(c)** Phase diagram section computed as in  
630 (b), showing in detail the boundaries between greenschists and amphibolite facies for the meta-  
631 granitoid rock. This boundary is defined by a series of univariant in- and out-reactions which assist  
632 strain localisation and shear zone formation in granitoid rocks. **(d)**  $T$ - $M_{H_2O}$  phase diagram section  
633 showing the phase relationships computed at different  $H_2O$  content ( $M_{H_2O}$ ). The phase fields located  
634 on the right side of the  $H_2O$ -in phase boundary (black dashed curve) corresponds to the same  $H_2O$ -  
635 saturated assemblage as those observed in the  $P$ - $T$  section (b-c).

636 **Figure 4.**  $T$ - $X$  diagrams and mass-transfer. See text for explanation.

637 **Figure 5.** Microstructures and rheology. Subscripts indicate either the magmatic (<sub>1</sub>) or  
638 recrystallised/metamorphic (<sub>2</sub>) origin of mineral phases. **(a)** Backscatter electron (BSE) image of a  
639 Kfs porphyroclast system in a mylonite from the Rieserferner pluton (modified after Ceccato et al.,  
640 2018). The microstructure includes the development of myrmekite (Myrm) from Kfs<sub>1</sub>, which lead to  
641 both grain size reduction and phase mixing, producing fine aggregates of Pl<sub>2</sub> + Qz<sub>2</sub>. Note the Kfs<sub>2</sub>

642 tails on Kfs<sub>1</sub> porphyroclasts, suggesting the occurrence of dissolution-precipitation processes and the  
643 chemical stability of Kfs. Other metamorphic reactions concur to the productions of fine-grained  
644 aggregates of Pl<sub>2</sub> + Ep + Ms + Bt. (b) Microstructure of an orthogneiss (low-strain stage of a large  
645 scale shear zone from the Fibbia meta-granite, Oliot et al., 2014) showing the occurrence of fine-  
646 grained layers of Ab<sub>2</sub> + Pl<sub>2</sub> derived from Pl<sub>1</sub> breakdown and neo-crystallisation warping around layers  
647 of dynamically recrystallised Qz<sub>2</sub> (modified from Fig. 3f of Oliot et al., 2014). (c) Log-Log  
648 differential stress vs. strain rate diagram showing the rheological curves computed for the  
649 deformation at 450 °C of quartz (Qz, red curve), a granitoid rock in which Qz + Pl deform by  
650 dislocation creep only (GSI, blue curve), and a fine-grained aggregate of Pl<sub>2</sub> + Qz<sub>2</sub> deforming by  
651 diffusion creep (GSS) only (green curve) (redrawn from Ceccato et al., 2018). The fine-grained Pl<sub>2</sub> +  
652 Qz<sub>2</sub> aggregate dominated by GSS creep mechanisms is several orders of magnitude weaker than both  
653 pure quartz and the Qz + Pl granitoid rock. (d) Log-Log grain size vs. differential stress diagram  
654 showing the contoured curves for strain rate computed for quartz at 450 °C from the combination of  
655 flow laws for dislocation creep (GSI) and diffusion creep (GSS) (see Ceccato et al., 2018 for more  
656 details on the adopted flow laws and computation). Each contour line is marked by the Log<sub>10</sub> value  
657 of the strain rate. The horizontal arrows indicate the effects on quartz microstructure and rheology  
658 during different processes of grain-size reduction: grain-size reduction by dynamic recrystallization  
659 (grey arrow, GSR) does not lead to any significant softening, shown by the similar strain rate  
660 conditions at which the white star and the grey star fall (at constant differential stress). Grain-size  
661 reduction by phase mixing (orange arrow, PM) may lead to effective softening, shown by the  
662 increased strain rate at which the orange star resides. The red curve represents the paleopiezometric  
663 relationship defined by Cross et al. (2017).

664 **Figure 6.** Sketch of the fabric evolution and microstructural processes within a shear zone with  
665 increasing shear strain. (a) Sketch of the microstructural evolution within the shear zone. (See the  
666 main text for explanation). (b) Sketch of grain-boundary sliding in equigranular Pl + Qz aggregates,  
667 showing also the nucleation of Kfs + Bt in the developed dynamic porosity. The red arrows indicate

668 the expulsion of fluids triggered by the mechanical contrast between the GSS creep-dominated Pl +  
669 Qz aggregate and the surrounding host rock (dominated by GSI-dislocation creep). The blue arrows  
670 represent the inward fluid flux triggered by the development of dynamic porosity during grain-  
671 boundary sliding.

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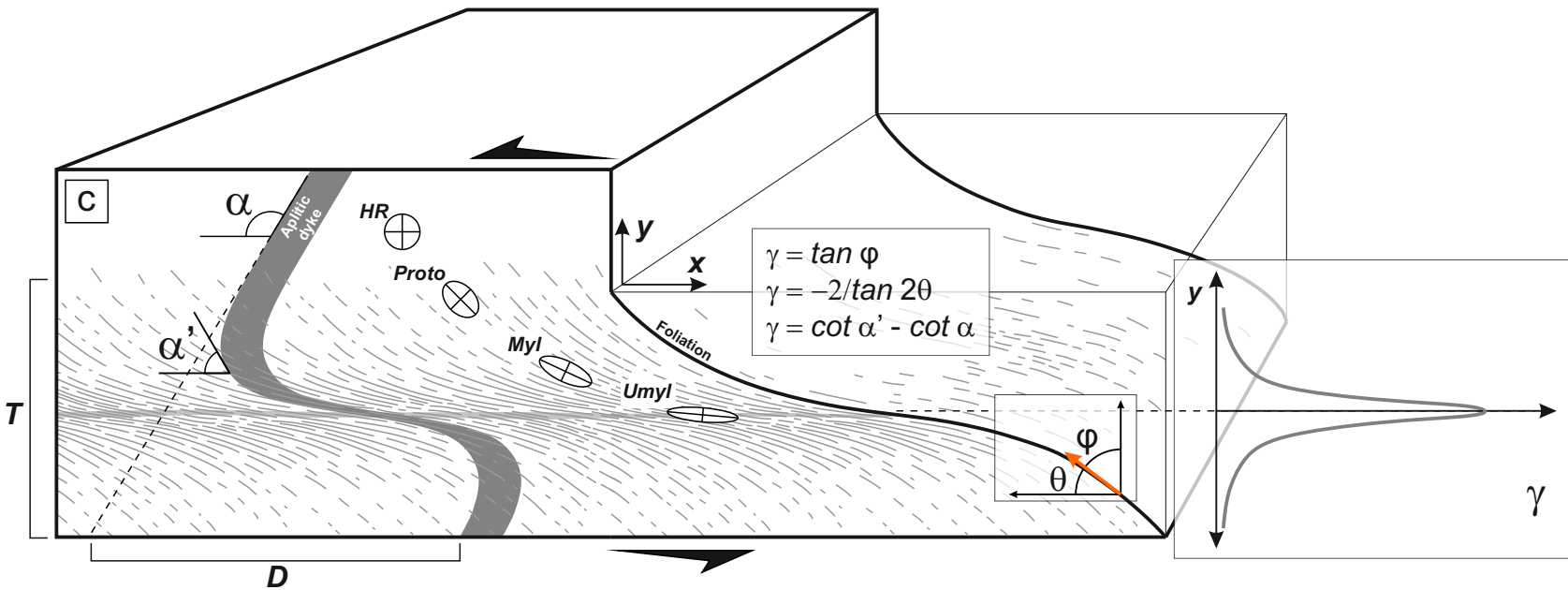
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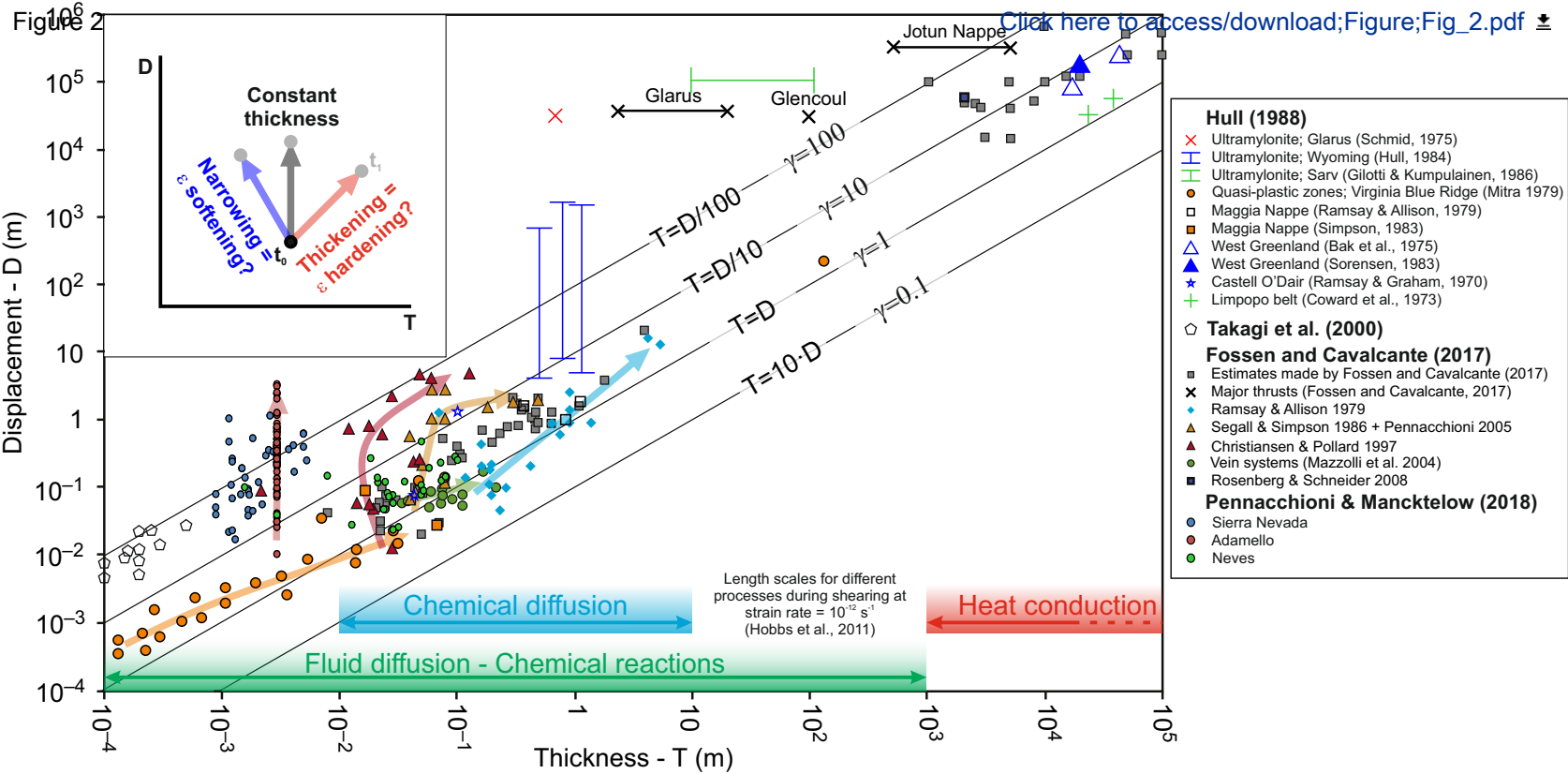
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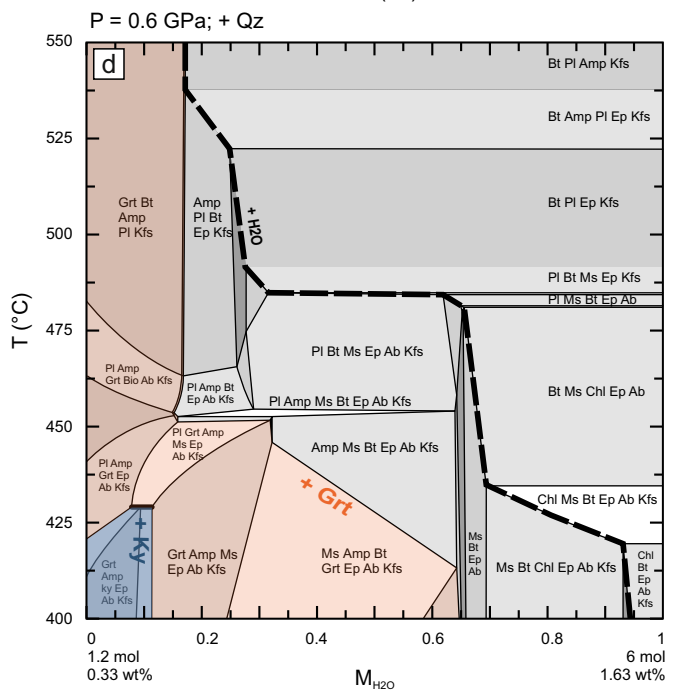
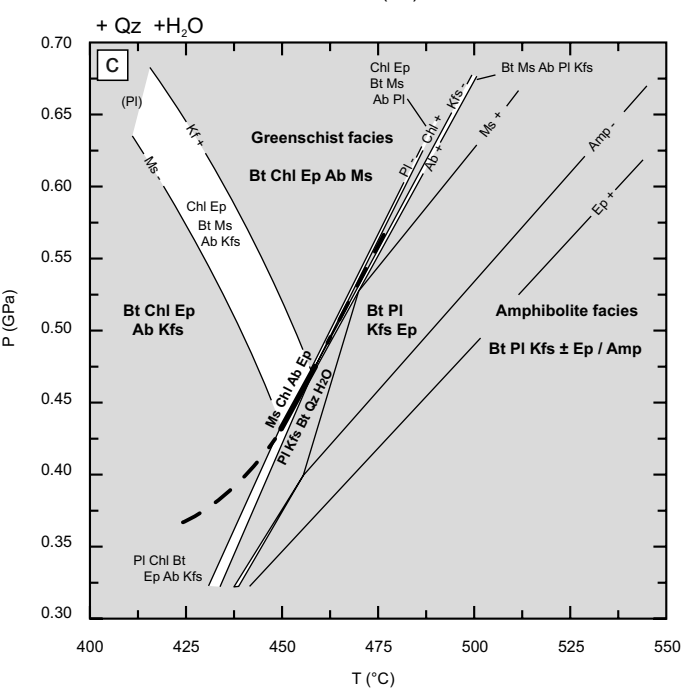
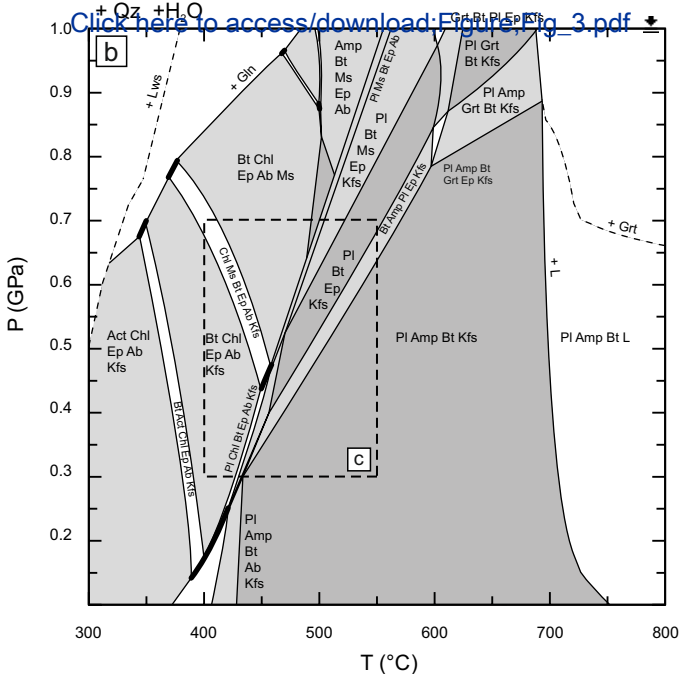
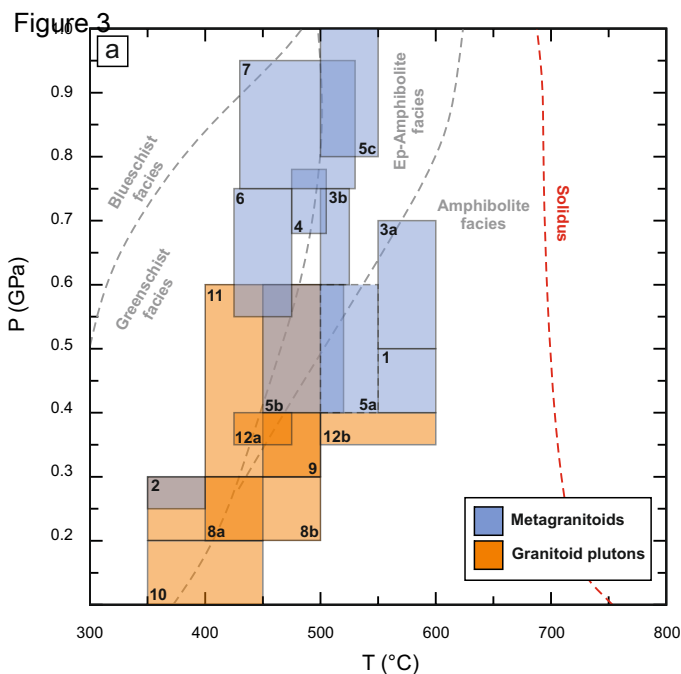
1108 Yardley, B.W.D., Valley, J.W., 1997. The petrologic case for a dry lower crust. *Journal of*  
1109 *Geophysical Research: Solid Earth*. <https://doi.org/10.1029/97jb00508>



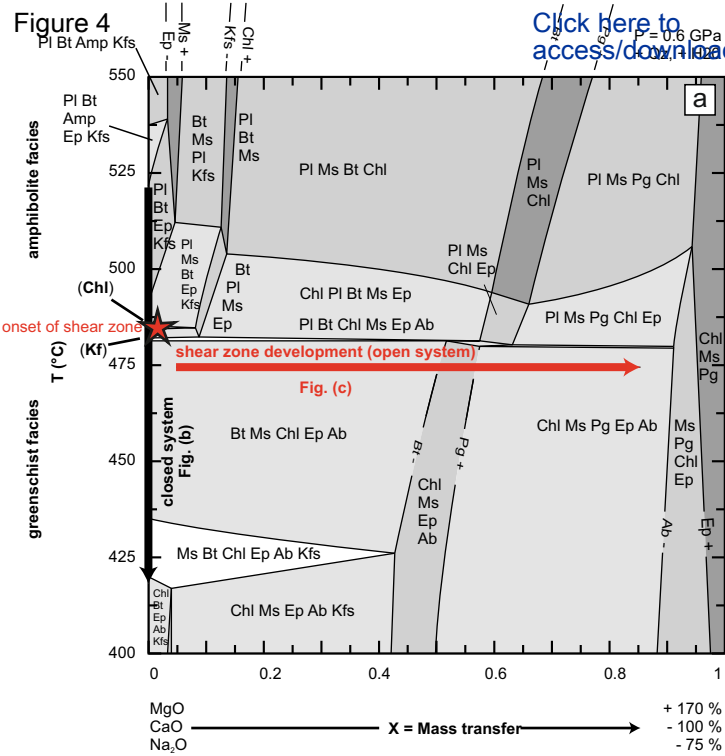
Figure 1



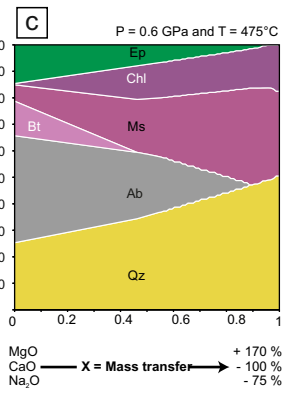
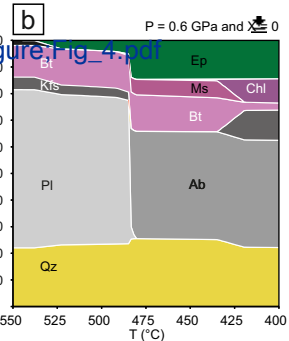




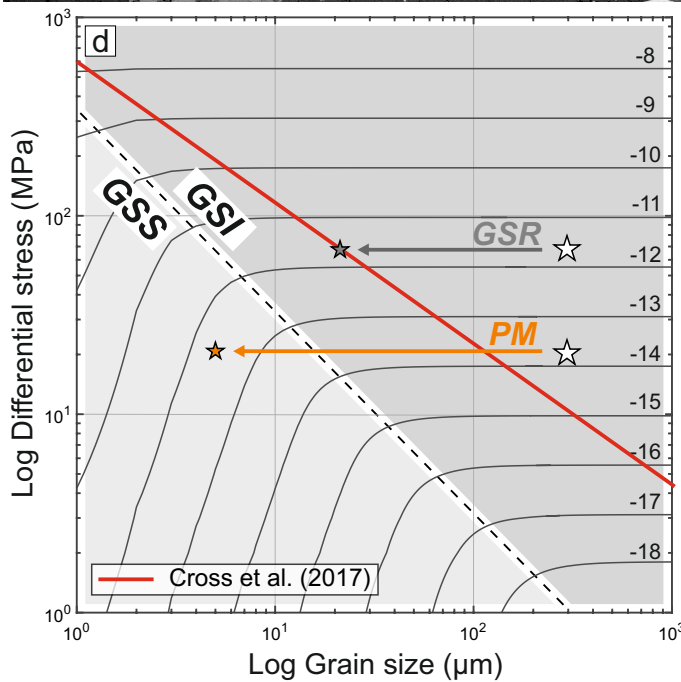
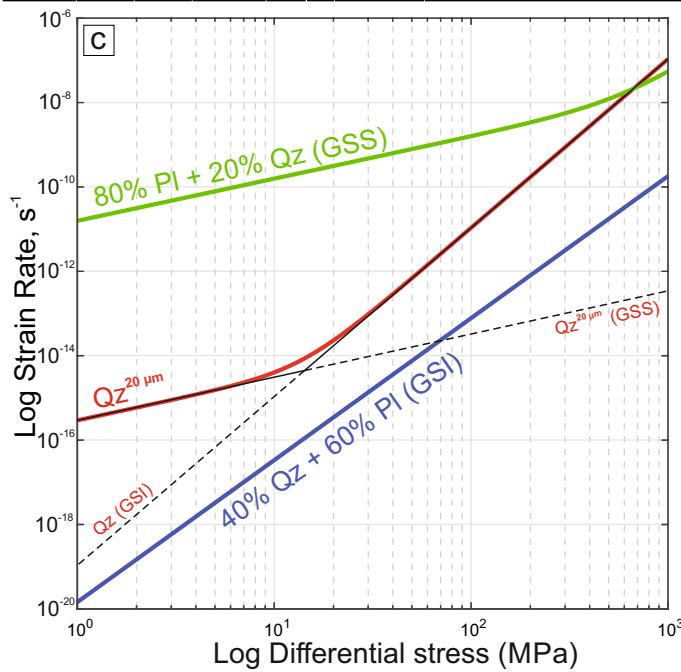
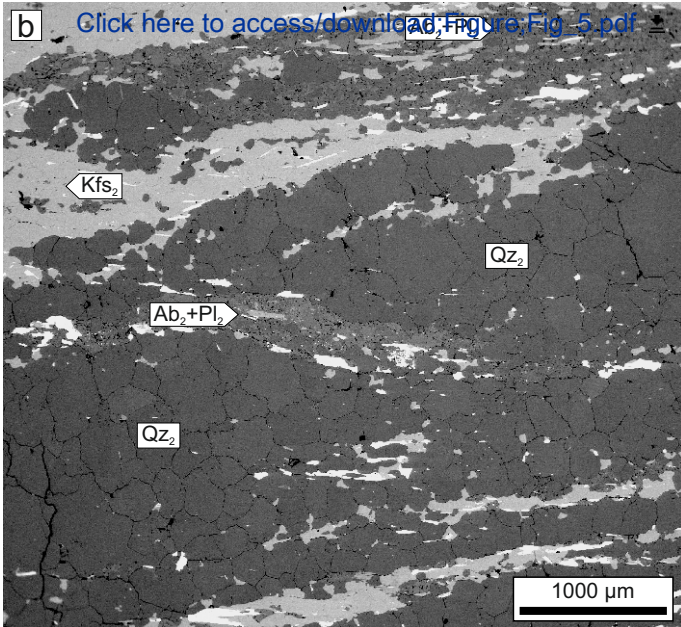
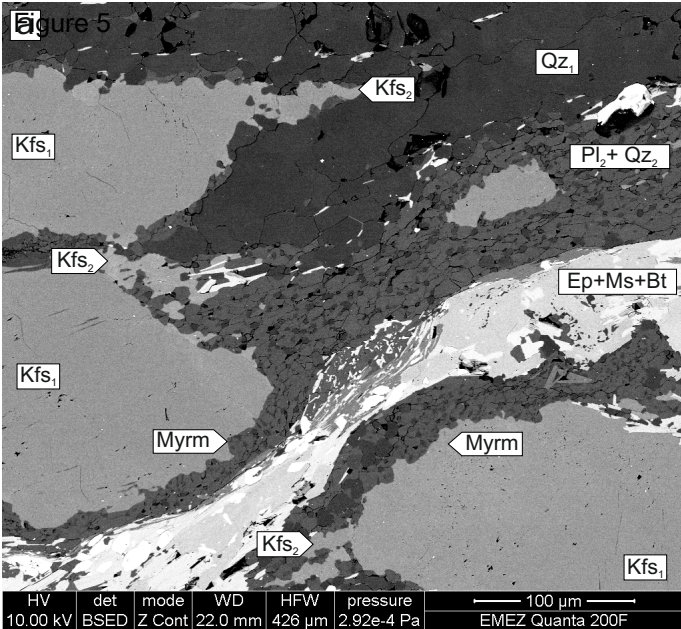
**Figure 4**

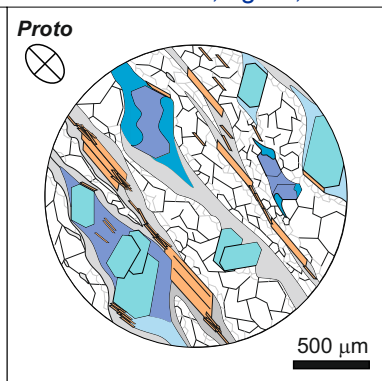
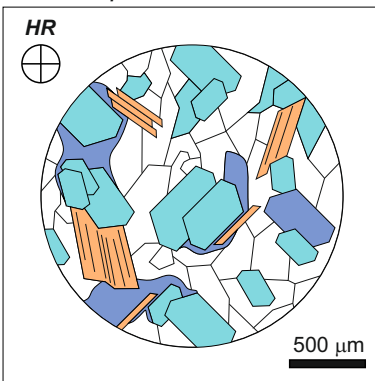


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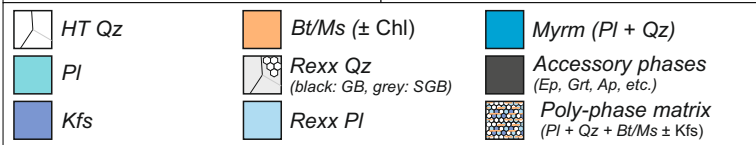
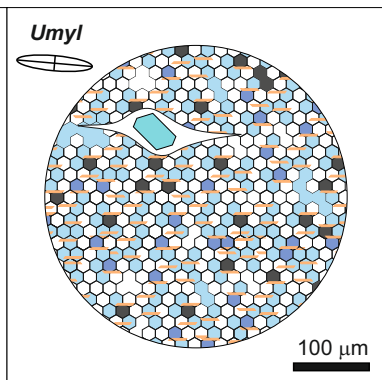
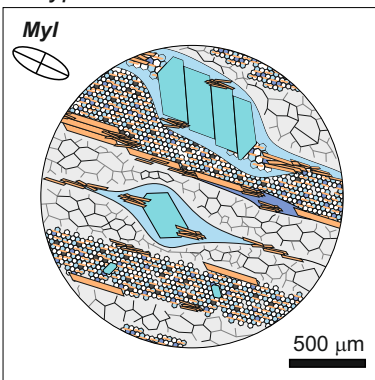






Mylonite (>50% rexx)  
Polyphase IWL

Ultramylonite (>90% rexx)



Grain-boundary sliding, phase nucleation  
and dynamic dilatancy

**b**

