

Dynamical Stability Indicator based on Autoregressive Moving-Average Models: Critical Transitions and the Atlantic Meridional Overturning Circulation

Marie Rodal,^{1, a)} Sebastian Krumscheid,² Gaurav Madan,³ Joseph Henry LaCasce,³ and Nikki Vercauteren^{3, b)}

¹⁾*FB Mathematik und Informatik, Freie Universität Berlin, Arnimallee 6, 14195 Berlin, Germany*

²⁾*Karlsruhe Institute of Technology, 76131 Karlsruhe, Germany*

³⁾*Section for Meteorology and Oceanography, Department of Geosciences, University of Oslo, Blindernveien 31, Kristine Bonnevis hus, 0371 Oslo, Norway*

(Dated: 27 September 2022)

A statistical indicator for dynamic stability known as the Y indicator is used to gauge the stability and hence detect approaching tipping points of simulation data from a reduced 5-box model of the North-Atlantic Meridional Overturning Circulation (AMOC) exposed to a time dependent hosing function. The hosing function simulates the influx of fresh water due to the melting of the Greenland ice sheet and increased precipitation in the North Atlantic. The Y indicator is designed to detect changes in the memory properties of the dynamics, and is based on fitting ARMA (auto-regressive moving-average) models in a sliding window approach to time series data. An increase in memory properties is interpreted as a sign of dynamical instability. The performance of the indicator is tested on time series subject to different types of tipping, namely bifurcation-induced, noise-induced and rate-induced tipping. The numerical analysis show that the indicator indeed responds to the different types of induced instabilities. Finally, the indicator is applied to two AMOC time series from a full complexity Earth systems model (CESM2). Compared with the doubling CO_2 scenario, the quadrupling CO_2 scenario results in stronger dynamical instability of the AMOC during its weakening phase.

A statistical indicator for dynamic stability is applied to simulation data from an ocean circulation model. The indicator assesses the stability of the time series data and gives indication of approaching tipping points. Three different types of tipping, defined by their causing mechanism, are explored. In addition, the indicator's reaction to the application of colored, as opposed to white, noise is assessed. Finally, the indicator is compared to other statistical early warning indicators.

tipping, or R-tipping, occurs when the system fails to track a continuously changing attractor and hence abruptly leaves the attractor.

Of these three, rate-induced tipping is certainly the least studied, however as demonstrated by Scheffer *et al.*³, Wicczorek *et al.*⁴ and more recently O'Keeffe and Wicczorek⁵, it is an important tipping mechanism that cannot be explained through classical bifurcation theory. Indeed, when the system is unable to track a continuously available quasi-stable state due to the system parameters changing too quickly, it might shift to another available equilibrium state without crossing a bifurcation boundary. There are a few methods available for estimating what exactly "too quickly" means, see Wicczorek and Perryman⁶, Ashwin, Perryman, and Wicczorek⁷, Vanselow, Wicczorek, and Feudel⁸ and O'Keeffe and Wicczorek⁵, but they depend strongly on the time-dependent parameter function; in particular its asymptotic properties. Finding generalizable methods for determining the rate of the parameter drift that induces tipping, will be of great interest going forward. Another issue of great practical importance is the question of how to obtain early warnings for such tipping points, in particular if classical methods for stability analysis also remain valid in the regime of rapid parameter changes.

Ritchie and Sieber⁹ showed that for rate-induced tipping, the most commonly used early-warning indicators, namely increase in variance and increase in autocorrelation, occur not when the equilibrium drift is fastest but with a delay. This suggests that these indicators might not be able to detect tipping before it has already occurred, although their analysis does give indication that the theory behind these indicators, the so-called "critical slowing down", may still hold for rate-induced tipping.

In this paper, we study an indicator for dynamic stability, from now on referred to as the Y indicator, initially proposed

I. INTRODUCTION

Tipping points, or critical transitions, are sudden, drastic changes in a system resulting from initial small perturbations. The study of tipping points is of particular interest to climate scientists and ecologists, as several theoretical studies highlight such tipping for an assortment of climatic and ecological systems, and observations also indicate that abrupt changes are, indeed, common in nature¹.

Ashwin *et al.*² classified tipping points according to the causing mechanism, yielding three classes of tipping points. Bifurcation-induced tipping, or B-tipping, occurs when a steady change in a parameter past a threshold induces sudden qualitative change in the system's behaviour. Noise-induced tipping, or N-tipping, occurs when short-timescale internal variability causes the system to transition between different co-existing attracting states. Finally, rate-induced

^{a)}marie.rodal@fu-berlin.de

^{b)}nikki.vercauteren@geo.uio.no

81 by Faranda *et al.*¹⁰. The Υ indicator uses auto-regressive
 82 moving-average or ARMA(p,q) models to estimate how close
 83 a system is to an equilibrium. It is based on the observation
 84 that the dynamics of an observable arising from a potentially
 85 complex system very close to a stable equilibrium will appear
 86 like a random walk with a tendency to be attracted to a
 87 well-defined equilibrium. When discretized, such dynamics
 88 can be well represented by an ARMA(1,0) process. When
 89 approaching a transition, however, the system may experience
 90 a critical slowing down and diverging memory properties.
 91 The trajectory of the observable hence experiences new
 92 timescales, which can be detected even with a limited dataset
 93 through an increase in the necessary memory lags of fitted
 94 ARMA(p,q) models¹¹. The Υ indicator thus defines a distance
 95 from the limiting random walk-like behaviour as a way to
 96 assess the dynamical stability properties of an observable.
 97 The indicator was applied to atmospheric boundary layer
 98 data by Nevo *et al.*¹² and Kaiser *et al.*¹³ and to atmospheric
 99 circulation data by Faranda and DeFrance¹⁴. They success-
 100 fully demonstrated the indicator's ability to both gauge the
 101 stability of a time series and detect tipping points. However,
 102 the indicator requires some additional testing, in particular
 103 concerning its performance for rate-induced tipping, which
 104 thus far has not been explored. It should be noted that several
 105 different early warning indicators based on ARMA models
 106 have been proposed. In fact, in Faranda, Dubrulle, and
 107 Pons¹¹ the authors propose the sum of the p and q orders of
 108 the model, as well as the sum of the model coefficients as
 109 potential indicators. The sum of the order parameters then
 110 gives an estimate for the memory lag of the process, while
 111 the sum of the model coefficients gives the persistence of this
 112 memory lag. 170

113 To further test the indicator, we have chosen the global
 114 oceanic 3-box model studied by Alkhayuon *et al.*¹⁵, which
 115 in turn is based upon the 5-box model of Wood *et al.*¹⁶
 116 The model represents a simplified Atlantic Meridional
 117 Overturning Circulation (AMOC), which transports warm
 118 surface water from the tropics to North America and Europe,
 119 resulting in a milder climate in these regions than what would
 120 otherwise be expected. Since the current is density driven,
 121 a large influx of freshwater due to the melting of land ice
 122 or increased precipitation in the North Atlantic, would be
 123 expected to result in a reduction in the AMOC flow strength.
 124 The question of whether the AMOC could undergo a sudden
 125 transition from a high flow strength state (the "on" state) to
 126 a state with weak or no overturning (the "off" state), is still
 127 debated. The latest assessment report of the International
 128 Panel for Climate Change (IPCC AR6) concludes that the
 129 AMOC strength will very likely decline in the future, but
 130 states with medium confidence that an abrupt collapse will not
 131 occur in the next century¹⁷. Simple box models, like the one
 132 presented in this paper, show bi-stability, while more realistic
 133 models like the global atmosphere-ocean general circulation
 134 models (AOGCMs) are largely mono-stable, implying that
 135 they do not exhibit the abrupt transition to an "off"-state
 136 so characteristic of the simpler models. However, there is
 137 limited evidence that the more complex models may be too
 138 stable (Weijer *et al.*¹⁸, Hofmann and Rahmsdorf¹⁹ and Liu
196

*et al.*²⁰), in particular that they mis-represent the direction
 of AMOC-induced freshwater transport across the southern
 boundary of the Atlantic (Liu *et al.*²⁰, Huisman *et al.*²¹,
 Liu, Liu, and Brady²², Hawkins *et al.*²³). Liu *et al.*²⁰
 demonstrated that by introducing a flux-correction term into
 the National Center for Atmospheric Research (NCAR)
 Community Climate System Model version 3 (CCSM3), they
 could make the formerly mono-stable system bi-stable.
 In addition, it has been suggested that paleoclimate data is
 consistent with abrupt changes in the surface temperature in
 the North Atlantic region in the past, as might be expected
 with a collapse of the AMOC. Boers²⁴ applied a statistical
 early warning indicator on Earth System Model (ESM)
 outputs, and found significant early-warning signals in eight
 independent AMOC indices. This was interpreted as a
 sign that the AMOC is not only a bistable system, but one
 approaching a critical transition.

Previously, the potential collapse of the AMOC has largely
 been attributed to the crossing of a bifurcation boundary in
 the bi-stable system. However, more recent analysis, see in
 particular Lohman and Ditlevsen²⁵, demonstrate the possi-
 bility of tipping before the bifurcation boundary is reached
 through the mechanism of rate-induced tipping. In addition,
 Lohman and Ditlevsen²⁵ demonstrate that due to the chaotic
 nature of complex systems a well-defined critical rate, i.e., the
 rate of parameter change at which the system tips, cannot be
 obtained, which in turn severely limits our ability to predict
 the long-term behavior of the system. They conclude that due
 to this added level of uncertainty, it is possible that the safe
 operating space with regard to future emissions of CO₂ might
 be smaller than previously thought. This suggests that proper
 evaluation of the probability of rate-induced tipping in the
 different tipping elements of the Earth System is of utmost
 importance in assessing the likelihood of dramatic future
 changes.

Regardless of whether the AMOC in actuality is bi-stable or
 mono-stable, the reduced 5-box model of Alkhayuon *et al.*¹⁵
 is the perfect test case for the Υ indicator as it exhibits both
 bifurcation-induced and rate-induced tipping, provided a time
 dependent hosing function is applied. The hosing function
 represents the influx of fresh water into the ocean due to
 increased precipitation and melting of land and sea ice in
 the North Atlantic region. Alkhayuon *et al.*¹⁵ provide an
 extensive analysis of the tipping mechanisms present in the
 model. Armed with such a well studied theoretical model, we
 will be able to systematically study the indicator's ability to
 not only detect bifurcation-induced and noise-induced, but
 also rate-induced tipping. We will additionally assess the
 indicator's ability to deal with colored noise, something that
 is known to cause issues for other early warning indicators,
 like the increase in variance and auto-correlation²⁴.

In reality, the ocean system has many more degrees of
 freedom than those included in the box models, and ulti-
 mately a mixture of different processes is likely to trigger
 tipping, if occurring. The Coupled Model Intercomparison
 Project (CMIP6), with the Community Earth System Model
 (CESM2)²⁶, provides an alternative AMOC model with

197 many more degrees of freedom. Two scenarios where the
 198 atmospheric CO₂ concentration is abruptly increased will
 199 be considered, providing monthly outputs of geographical
 200 density differences on which the Υ indicator will be applied
 201 In these model scenarios, the abrupt change in CO₂ is
 202 followed by a response of the Earth system, and after 2-3
 203 decades, freshwater eventually circulates in the sub-polar
 204 gyre²⁷. This response hence offers similarities with the
 205 hosing experiments done in the box models. While the two
 206 scenarios are insufficient to assess the potential bistabil-
 207 ity of the AMOC, the Υ indicator will be used to assess the
 208 dynamical stability of the AMOC during its weakening phase.
 209

210 II. THE Υ -INDICATOR FOR EARLY-WARNING SIGNALS

211 In what follows, we will briefly outline the method used to
 212 determine the stability of the time series data. Further details
 213 can be found in Faranda *et al.*¹⁰, Faranda and DeFrance¹⁴
 214 Nevo *et al.*¹² and Kaiser *et al.*¹³

215 The method relies on an accurate representation of a com-
 216 plex dynamical system close to a metastable state by a ran-
 217 dom walk-like behavior with a tendency to be attracted to the
 218 metastable state. Changes in the system's stability are then
 219 characterized as statistically significant deviations from that
 220 local behavior, indicating that the system currently does not
 221 reside close to a metastable state. Indeed, the local dynam-
 222 ics of a continuous-time random dynamical system (i.e., a
 223 stochastic differential equation) near a metastable state come
 224 close to the dynamics of a stochastic spring (i.e., an Ornstein-
 225 Uhlenbeck process), whose discrete-time observations are
 226 well approximated by an ARMA (1,0) process. Here, ARMA
 227 denotes the space of autoregressive moving-average models
 228 with the numbers in parentheses denoting the order of the
 229 model. A time series $x(t)$, $t \in \mathbf{Z}$, is an ARMA(p,q) process
 230 if it is stationary and can be written as

$$231 \quad x(t) = v + \sum_{i=1}^p \phi_i x_{t-i} + \sum_{j=1}^q \theta_j w_{t-j} + w_t \quad (1)$$

232 with constant v , coefficients ϕ_i , θ_j and $\{w_t\}$ being white noise
 233 with positive variance σ^2 (see Brockwell and Davis²⁸ for an
 234 introductory text). In addition, constraints are imposed on the
 235 coefficients ϕ_i and θ_j to ensure that the process in (1) is sta-
 236 tionary and satisfies the invertibility condition. Intuitively, the
 237 variables p and q say something about the memory lag of the
 238 process, while the prefactors ϕ_i and θ_j relate to the persistence
 239 of said memory lag. One expects that the higher the values for
 240 q and p , the longer the system, once perturbed from its equi-
 241 librium state, would need to return to equilibrium. It is this
 242 intuitive notion that the statistical indicator denoted Υ takes
 243 advantage of. Indeed, when approaching a critical transition
 244 the response of the system to perturbations can become in-
 245 creasingly long (referred to as a critical slow down), and this
 246 translates into diverging memory properties of the statistical
 247 signal. Hence, an ARMA(p,q) model will require higher or-
 248 ders to incorporate the memory effects. By fitting the model

(1) repeatedly to a time series data set for varying values of p
 and q , one can, through application of an appropriate informa-
 tion criterion, obtain the values of p and q that best represent
 the time series data. For this purpose, we choose the Bayesian
 information criterion, BIC:

$$\text{BIC} = -2 \ln L(\hat{\beta}) + \ln(\tau)(p + q + 1) \quad (2)$$

where $\hat{\beta}$ denotes the maximum likelihood estimator of $\beta =$
 $(v, \phi_1, \dots, \phi_p, \theta_1, \dots, \theta_q)$, which is obtained by maximising
 the likelihood function L associated with the ARMA(p,q)
 model (1) for a given time series; see Brockwell and Davis²⁸
 for details. The best fitting ARMA(p,q) model is then deter-
 mined as the one that minimizes the BIC. The second term in
 equation (2) punishes complex models with high p and q val-
 ues, and is the reason why we prefer to use the BIC over other
 criteria, such as the perhaps more familiar Akaike Information
 Criterion. Here, τ denotes the number of discrete points in the
 time series to which the ARMA model is fitted. We refer to τ
 as the *window length*.

Finally, the stability indicator is defined as

$$\Upsilon(p, q; \tau) = 1 - \exp\left(\frac{-|\text{BIC}(\bar{p}, \bar{q}) - \text{BIC}(p, q)|}{\tau}\right) \quad (3)$$

where \bar{p} and \bar{q} indicate the order of what we refer to as the
 theorized *base model*. This is the ARMA(p,q) model, charac-
 terized by a specific value of $q = \bar{q}$ and $p = \bar{p}$, to which the
 chosen best fit is compared. The Υ -indicator takes on values
 between 0 and 1, where lower values imply a higher degree
 of stability. The intuition behind using the difference in BIC
 values between the chosen "best" model and a base model is
 that this quantity assesses just how much better the model with
 the lower BIC value approximates the fitted data compared to
 the other. The significance threshold for deviations in the BIC
 values between an ARMA(p,q) and the base model, simply
 denoted as $|\Delta\text{BIC}|$, is $|\Delta\text{BIC}| > 2$. The differences in BIC
 values can be directly related to the Bayes Factor, see Preacher
 and Merkle²⁹, which is another way of quantifying the likeli-
 hood of one model over another.

For the data sets analysed by Faranda *et al.*¹⁰, it was deter-
 mined that the appropriate base model is the ARMA(1,0)
 model, i.e., $\bar{p} = 1$ and $\bar{q} = 0$, which can be viewed as a
 time discretized Langevin process. In later work by Nevo
*et al.*¹² and Kaiser *et al.*¹³ the authors continued to rely
 on ARMA(1,0) as the base model. While Faranda *et al.*¹⁰
 used a statistical argument to justify the choice of the base
 model, Nevo *et al.*¹² and Kaiser *et al.*¹³ argued, as already
 noted above, that the dynamics near a stable state can be ap-
 proximated as that of a stochastic spring, further strengthen-
 ing the case for ARMA(1,0) as the general choice of base
 model. However, due to the additional well-posedness con-
 straints on the autoregressive and moving-average coefficients
 ϕ_i and θ_j in (1), depending on the treatment of constraints by
 the fitting routine one can have cases where the BIC value
 of the ARMA(1,0) process is smaller than the corresponding
 value for the chosen ARMA(p,q) model. In these cases the
 ARMA(1,0) process is rejected as the best fit, despite having
 the lowest BIC value, due to violating the stationarity or in-
 vertibility conditions required for a numerically well behaved

304 fit. Thus, in this scenario it becomes unclear how to determine 354
 305 the 'distance' between the states. To overcome this issue we 355
 306 have chosen to modify the Υ indicator to allow for a second 356
 307 base state, namely the ARMA(0,0) model. This model is just 357
 308 white noise, possibly with a drift, and is guaranteed to sat- 358
 309 isfy all the auxiliary conditions for the obvious reasons that 359
 310 there are no coefficients available to violate them. We con- 360
 311 sider ARMA(0,0) as a special case of ARMA(1,0) in which 361
 312 $\phi_1 = 0$. The use of the ARMA(1,0) process as a base model 362
 313 was partly justified by the image of a particle trapped in a po- 363
 314 tential well, where a restoring force keeps the particle oscil- 364
 315 lating around the equilibrium. The justification for including 365
 316 ARMA(0,0) as a potential base model follows a similar argu- 366
 317 ment, except that in this case the noise amplitude is too low 367
 318 compared to the width of the potential well to feel the restor- 368
 319 ing force. To use both base models, we first introduce 369

$$370 \Delta\text{BIC}_0(p, q) := \text{BIC}(0, 0) - \text{BIC}(p, q) \quad (4)_{371}$$

372 and 373

$$374 \Delta\text{BIC}_1(p, q) := \text{BIC}(1, 0) - \text{BIC}(p, q) \quad (5)_{375}$$

376 With this, the modified Υ -Indicator for the extended base 376
 377 model class can be written as 377

$$378 \Upsilon(p, q; \tau) = 1 - \exp\left(\frac{-\min\{|\Delta\text{BIC}_0(p, q)|, |\Delta\text{BIC}_1(p, q)|\}}{\tau}\right)_{379}$$

(6)₃₈₁

382 In addition, it must be specified that in the cases where the 382
 383 constrained fitting failed for the ARMA(1,0) model so that 383
 384 $\Delta\text{BIC}_1(p, q)$ may be negative, $\Delta\text{BIC}_0(p, q)$ is automatically 384
 385 chosen in practise. For obvious reasons, there cannot be a 385
 386 case where $\Delta\text{BIC}_0(p, q)$ is itself negative. 386

387 Furthermore, following Faranda, Dubrulle, and Pons¹¹, we 387
 388 define the *order*, \mathcal{O} , and *persistence*, \mathcal{R} , of an ARMA(p, q) 388
 389 process as 389

$$391 \mathcal{O} = p + q, \quad (7)_{392}$$

$$393 \mathcal{R} = \sum_{i=1}^p |\phi_i| + \sum_{j=1}^q |\theta_j|, \quad (8)_{394}$$

395 where ϕ_i and θ_j denote the autoregressive and moving- 395
 396 average coefficients, respectively. While the order relates 396
 397 to the memory lag of the process, the persistence relates to 397
 398 the *persistence* of said memory lag, hence the name. When 398
 399 approaching a tipping point, one would expect one out of two 399
 400 things to happen: either both the persistence and the order 400
 401 increase significantly, due to the increased memory of the 401
 402 process, or the order remains constant, and the persistence 402
 403 approaches the value of the order \mathcal{O} , indicating a loss of 403
 404 stationarity. According to Faranda, Dubrulle, and Pons¹¹, the 404
 405 latter alternative corresponds to a case in which the potential 405
 406 landscape of the system does not change considerably when 406
 407 approaching the transition. 407

408 This observation strengthens the case for the modified Υ 408
 409 indicator in contrast to excluding windows of the time series 409
 410 where $\Delta\text{BIC}_1(p, q)$ is negative, as these periods are indicative 410
 411 of an instability resulting from the loss of stationarity of the 411

ARMA(1,0) process.

To apply the method to a time series data set, one first 370
 371 has to ensure stationarity of the data. This can be done in 371
 372 two ways, depending on the nature of the time series. In 372
 373 some cases, it is sufficient to split the time series into small 373
 374 enough intervals, so that within each interval the time series 374
 375 is approximately stationary. To check for stationarity one 375
 376 runs a Kwiatkowski–Phillips–Schmidt–Shin (KPSS) tests on 376
 377 the intervals. This way, one also obtains an upper bound on 377
 378 the length of the intervals; see Kaiser *et al.*¹³. The other 378
 379 option is to not assume stationarity from the outset, and 379
 380 instead allow for application of a differencing routine to the 380
 381 separate intervals, achieving stationarity that way. In that 381
 382 case, a KPSS test is run on each interval, and if the interval 382
 383 is found to not be stationary, differencing is applied. This 383
 384 process is then repeated until stationarity is achieved. The 384
 385 KPSS test is to be preferred over the unit root test due to the 385
 386 danger of over-differencing (Hyndman and Khandakar³⁰). 386
 387 As we wish to study rate induced tipping phenomena, which 387
 388 yields highly non-stationary time series even for very small 388
 389 interval lengths, the latter method is to be preferred. By this 389
 390 choice we go from an ARMA to an ARIMA model, in which the 390
 391 I stands for "integrated" in reference to the differencing 391
 392 routine used to ensure the stationarity of the time series. 392

393 Provided one can select sufficiently long time series intervals 393
 394 where the process is approximately stationary, one can fit 394
 395 ARMA(p, q) models to available observations during these 395
 396 intervals, and through the Υ indicator obtain an estimate for 396
 397 how close any given interval is to an equilibrium state. To 397
 398 determine the best fit, we use the `auto.arima` function found 398
 399 in the `FORECAST` R package, setting `BIC` as the information 399
 400 criterion used for model selection. Since we will not assume 400
 401 stationarity of the time series, `auto.arima` first determines 401
 402 the correct differencing order before continuing with the fitting 402
 403 procedure; the details of said procedure can be found in 403
 404 Hyndman and Khandakar³⁰. 404

405 It is clear that the method is strongly dependent upon the size 405
 406 of the intervals, which we will refer to as the window length, 406
 407 τ . This is not only due to the inclusion of the $1/\tau$ factor in 407
 408 the exponential, but also due to the inherent τ -dependence of 408
 409 $\text{BIC}(p, q)$ and $\text{BIC}(1, 0)$. In fact, the rationale for including 409
 410 the $1/\tau$ factor in the definition of Υ is to attempt to remove 410
 411 or reduce this dependence. From equation (2) one might 411
 412 conclude that the correct scaling would be $1/\ln(\tau)$, as 412
 413 opposed to $1/\tau$. However, we do not only want to remove the 413
 414 dependence on τ , but also include the significance threshold 414
 415 for ΔBIC , such that the Υ value of any point where ΔBIC is 415
 416 below 2 is suppressed relative to other points. 416

III. APPLICATION TO THE GLOBAL OCEANIC 3-BOX MODEL

To determine the validity of the Υ -indicator as a measure 417
 418 of stability, as well as its ability to detect different types of 418

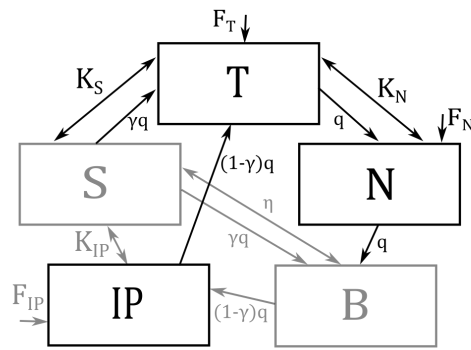


FIG. 1: Sketch of the 5-box model for the Atlantic Meridional Overturning Circulation (AMOC). Here, a light gray coloring is used to denote the two boxes whose salinities do not change, as well as all the arrows indicating terms which do not appear in the equations describing the dynamics of the 3-box model. Adapted from Alkhayuon *et al.*¹⁵.

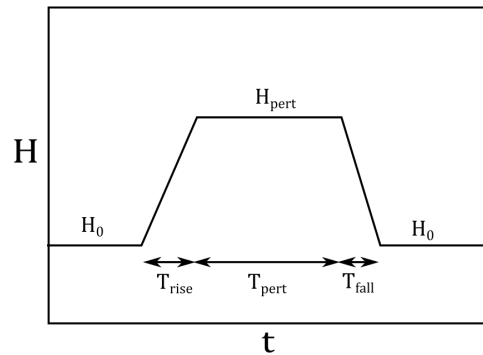


FIG. 2: Schematic illustration of the piece-wise linear hosing function used to simulate the influx of fresh water. Adapted from Alkhayuon *et al.*¹⁵.

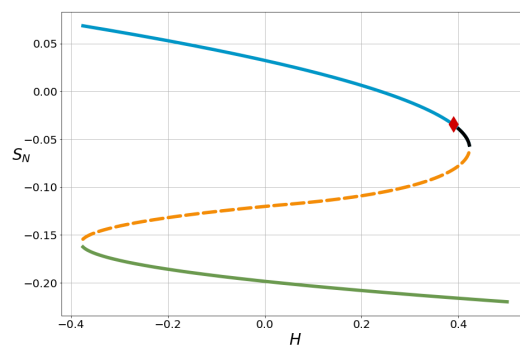


FIG. 3: Bifurcation diagram for S_N , for the 3-box model of the AMOC. The dashed line denotes the unstable equilibrium branch. The red diamond denotes the location of the hopf-bifurcation.

410 tipping points, we start by applying the method to the global
 411 oceanic 3-box model discussed by Alkhayuon *et al.*¹⁵. The
 412 3-box model of Alkhayuon *et al.*¹⁵ is a simplification of the
 413 5-box model of Wood *et al.*¹⁶ in which the salinity of the
 414 Southern Ocean (S) and the Bottom waters (B) is assumed to
 415 be approximately constant. The model thus consists of 5 separate
 416 boxes, of which only 3 boxes, namely the North Atlantic
 417 (N), Tropical Atlantic (T) and Indo-Pacific (IP) boxes have
 418 varying salinities S . A schematic illustration of the model
 419 is shown in Figure 1. See Alkhayuon *et al.*¹⁵ or Wood *et al.*¹⁶
 420 for a detailed exposition of the box model. We note that the
 421 parameters of the box model are tuned using the full complexity
 422 FAMOUS AOGCM model, with varying levels of CO_2 .
 423 The parameters used in this paper are for the case $2\times\text{CO}_2$ as
 424 compared to pre-industrial times.
 425 We denote salinity by S_i , the volume by V_i and the fluxes by
 426 F_i , where $i \in \{N, T, S, IP, B\}$ denotes the respective boxes.
 427 Let Γ denote the AMOC flow defined by

$$\Gamma = \lambda \left[\alpha(T_S - T_0) + \frac{\beta}{100}(S_N - S_S) \right] \quad (9)$$

429 The model approximates a buoyancy-driven flow, with a trans-
 430 port proportional to the density difference between the boxes,
 431 assuming a linearized equation of state. The evolution equa-
 432 tions for the salinities S_N and S_T are

$$433 \quad \frac{V_N}{Y} \frac{dS_N}{dt} = \Gamma(S_T - S_N) + K_N(S_T - S_N) - 100F_N S_0 \quad (10)$$

$$434 \quad \frac{V_T}{Y} \frac{dS_T}{dt} = \Gamma[\gamma S_S + (1 - \gamma)S_{IP} - S_T] + K_S(S_S - S_T) + K_N(S_N - S_T) - 100F_T S_0 \quad (11)$$

435 for $\Gamma \geq 0$, and

$$436 \quad \frac{V_N}{Y} \frac{dS_N}{dt} = |\Gamma|(S_B - S_N) + K_N(S_T - S_N) - 100F_N S_0 \quad (12)$$

$$437 \quad \frac{V_T}{Y} \frac{dS_T}{dt} = |\Gamma|(S_N - S_T) + K_S(S_S - S_T) + K_N(S_N - S_T) - 100F_T S_0 \quad (13)$$

438 for $\Gamma < 0$, where S_B and S_S are regarded as fixed parameters 475
 439 and $Y = 3.15 \times 10^7$, which converts the time unit from sec-
 440 onds to years. S_0 is a reference salinity, and K_i are coefficients 476
 441 associated with the gyre strengths. We note that all the salinity
 442 values are given as perturbations from a background state, see 477
 443 Appendix A of Alkhayuon *et al.*¹⁵ for details on the transfer 478
 444 mation. Since the total salinity is assumed to be conserved, 479
 445 the salinity of the Indo-Pacific (IP) box, S_{IP} , can be computed 480
 446 from S_N and S_T .

447 The values of the assorted parameters can be found in Table 482
 448 and Table 2. 483

449 The fluxes, F_N and F_T , are linear functions of the hosing func- 484
 450 tion $H(t)$ which simulates the influx of fresh water. In the case 485
 451 of $2 \times \text{CO}_2$ the fluxes are (see Wood *et al.*¹⁶) 486

$$452 \quad F_N = 0.486 \times 10^6 + H(t) 0.1311 \times 10^6 \quad (14) \quad 487$$

$$453 \quad F_T = -0.997 \times 10^6 + H(t) 0.6961 \times 10^6 \quad (15) \quad 488$$

454 where all fluxes are given in units of Sverdrup (Sv). 490
 455 The values for the case of $1 \times \text{CO}_2$ can be found in Table 5 of 491
 456 Alkhayuon *et al.*¹⁵. 492

457 Figure 3 shows the bifurcation diagram for S_N ; for S_T 493
 458 we refer to Alkhayuon *et al.*¹⁵ The bifurcation diagram 494
 459 for the flow strength Γ is qualitatively similar, since all 495
 460 other parameters in Eq. 9 are kept constant. The diagram 496
 461 clearly shows that this is a bi-stable system with two stable 497
 462 equilibrium branches connected by an unstable branch. 498

463 The upper equilibrium branch loses stability, not at the 499
 464 saddle-node bifurcation, but rather due to a Hopf-bifurcation, 500
 465 indicated by a red diamond in the diagram. Thus, part of the 501
 466 upper equilibrium branch, denoted in black, is in fact unstable. 502

467 To simulate the influx of fresh water we apply a time 504
 468 dependent, piece-wise linear hosing function, $H(t)$ (see 505
 469 Figure 2), to equations (10)-(13). Here 506

$$471 \quad H(t) = \begin{cases} H_0 & t < 0, \\ H_0 + \alpha(t) & t \in [0, T_{rise}], \\ H_{pert} & t - T_{rise} \in [0, T_{pert}], \\ H_{pert} - \beta(t) & t - T_{rise} - T_{pert} \in [0, T_{fall}], \\ H_0 & t \geq T_{rise} + T_{pert} + T_{fall}, \end{cases} \quad (16) \quad 510$$

472 where $\alpha(t)$ and $\beta(t)$ are linear functions ensuring continuity 514
 473 of $H(t)$. If we define the rise and fall rates, as 515

$$474 \quad r_{rise} = \frac{|H_{pert} - H_0|}{T_{rise}} \quad \text{and} \quad r_{fall} = \frac{|H_{pert} - H_0|}{T_{fall}} \quad (17) \quad 517$$

then

$$\alpha(t) = r_{rise}t \quad \text{and} \quad \beta(t) = r_{fall}(t - T_{rise} - T_{pert}) \quad (18)$$

As demonstrated by Alkhayuon *et al.*¹⁵, whether the sys-
 tem undergoes a transition from one stable state to the other,
 is dependent not only on the value of H_{pert} , but on the rise
 and fall rates, r_{rise} and r_{fall} , as well as the perturbation time
 T_{pert} . In particular, they demonstrate that even when H_{pert} is
 above the bifurcation value that destabilizes the upper equilib-
 rium branch, the system may still return to this equilibrium,
 provided T_{fall} is short enough; a process which they termed
avoided B-tipping. In addition, they showed that if T_{pert} is too
 short, the system will not tip, but return to the initial equilib-
 rium branch.

In what follows, we will apply the Y indicator as described
 in the previous section to time series data generated by the 3-
 box model. We will separately study time series undergoing
 rate-, noise- and bifurcation-induced tipping, while attempt-
 ing to assess the indicator's ability to gauge the stability of the
 time series as it approaches the tipping point. Before proceed-
 ing, we should clarify one point regarding noise-induced tip-
 ping, and what is meant by an early warning indicator in this
 context. Noise-induced tipping is inherently unpredictable,
 and hence one might conclude that any attempt at predicting
 such transitions is doomed to fail based on a single time se-
 ries. In contrast, assuming the underlying model is known,
 one could use ensembles of realizations to estimate the like-
 lihood of noise-induced transitions. Examples of these statisti-
 cal approaches are discussed in Thompson and Sieber³¹.
 Although one cannot expect to develop an *early warning* indi-
 cator for these types of transitions, one should at the very least
 be able to tell, from time series data, once such a transition has
 occurred, i.e., when the unstable equilibrium branch has been
 crossed and the system is approaching a different equilibrium.
 The objective should then be to develop an indicator that is
 able to identify this induced instability as soon as possible af-
 ter the transition.

Finally, we note that, while it is possible to extend ARMA
 fitting to multivalued time series data, we have chosen to not
 go down that route, and instead only apply the indicator to a
 single time series for the salinity values from the North At-
 lantic basin, S_N . The reason for choosing S_N over S_T is that
 within the 3-box model, the equilibrium branches of S_N are
 that much further apart, making the transitions easier to see.
 Such a simplification might at first glance seem rather con-

The Υ indicator for Early Warning

7

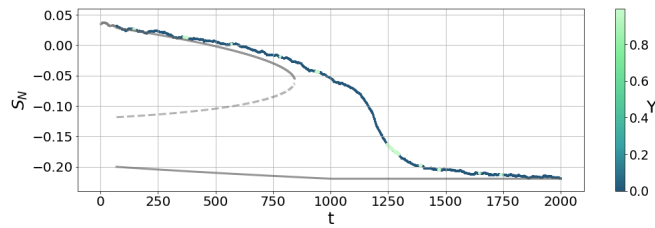


FIG. 4: Bifurcation-induced tipping, color coded according to the value of Υ with window length, $\tau = 350$. The gray lines denote the equilibrium branches, with the dashed line corresponding to the unstable branch. We clearly see several brightly colored points corresponding to a high values of Υ , which should be indicative of a high degree of instability and an approaching tipping point.

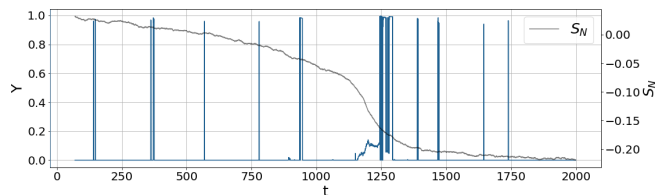


FIG. 5: Υ as a function of time for a time series of S_N undergoing B-tipping.

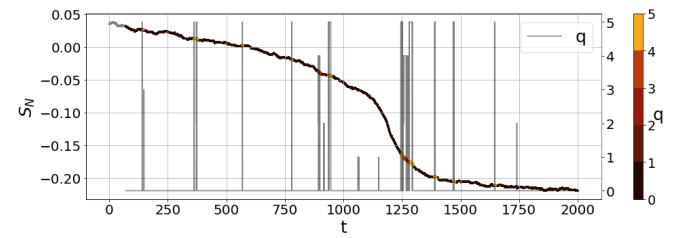
519 trived, however we argue that, as the goal of any indicator is
 520 to be used on real-world time series data in which the connec-
 521 tion to other time series is largely unknown, it is reasonable
 522 to only concentrate on one time series, despite the underlying
 523 system being multidimensional.

524 A. Bifurcation-induced Tipping

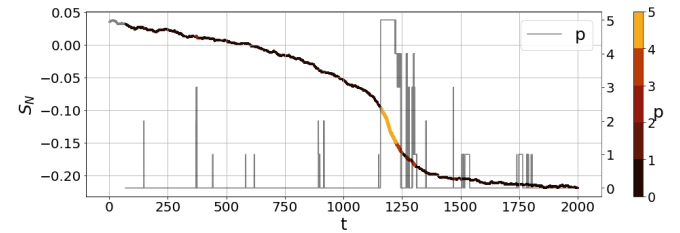
525 To induce B-tipping in the 3-box model, we gradually
 526 change $H(t)$ according to equation (16), with $H_0 = 0$, $H_{pert} = 0.5$, $T_{rise} = 1000$. This corresponds to an increase in the fresh-
 527 water fluxes F_T and F_N , corresponding to the flux into the
 528 tropical and North Atlantic boxes, by approximately 34% and
 529 13%, respectively. This, in turn, corresponds to roughly a 0.1-
 530 0.2 Sv increase, in line with freshwater "hosing" experiments
 531 of the North Atlantic³². We let T_{pert} go to infinity, such that
 532 $H(t)$ never returns to its initial value. As $H(t)$ changes, S_N
 533 follows the upper equilibrium branch as sketched in Figure
 534 3, until it reaches the hopf-bifurcation (around $H = 0.4$), at
 535 which point the upper equilibrium branch becomes unstable
 536 and S_N starts approaching the lower equilibrium branch. We
 537 choose a window length of 350 points corresponding to about
 538 70 years.

539 Figure 4 shows the time series of S_N color coded according
 540 to the value of Υ , with brighter colors corresponding to higher
 541 values of Υ and hence a greater degree of instability. Figure
 542 5 shows Υ as a function of time, with clear peaks corresponding
 543 to brightly colored points in Figure 4.

544



(a)



(b)

FIG. 6: Bifurcation induced tipping of $S_N(t)$, color coded according to the value of the best-fit ARMA model orders (a) q and (b) p (scatter plot). The line plots additionally show the same values for q and p as functions of time in (a) and (b), respectively.

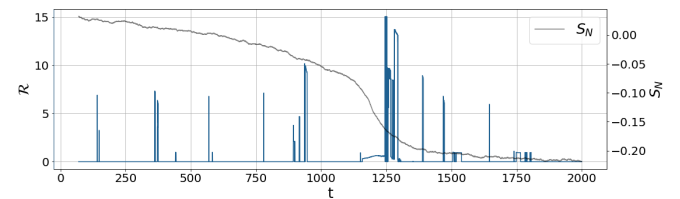


FIG. 7: Plot of the persistence \mathcal{R} (Eq. 8) as a function of time for a time series of S_N undergoing B-tipping.

It should be noted that low amplitude white noise is also applied to facilitate ARIMA model fitting. The noise intensity is kept small enough to avoid noise-induced tipping.

Figures 4 and 5 clearly indicate that there are several points on the time series as it approaches the transition, which are deemed to have a high degree of instability. We further note that, although the result is not shown here, the high Υ values in Figures 4 and 5 correspond to intervals for which $\Delta\text{BIC}_1(p, q)$ is negative, indicating that, as discussed previously, the ARMA(1,0) model would, when only considering BIC values, be the better fit, but it violates the auxiliary conditions, indicating a loss of stationarity. Hence, at these points ARMA(1,0) is excluded as a possible model, implying that ARMA(0,0) is the chosen base model.

In addition, we look at the order of the best-fit ARMA model, namely the q and p values, as well as the persistence, to gain further insight into the stability properties of the time series. Figure 6 shows the time series of S_N color coded according to the values of q and p . When comparing with Figure 4, this seems to indicate that the high values of Υ appearing before the transition are primarily associated with an increase in the

545

546

547

The Y indicator for Early Warning

8

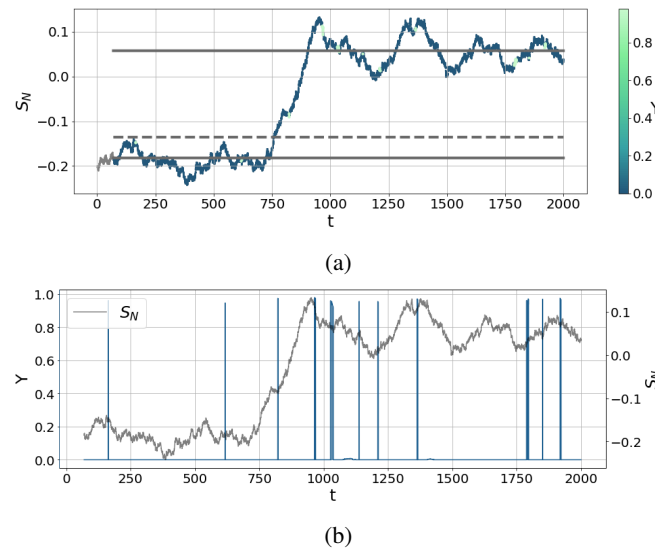


FIG. 8: (a) Noise-induced tipping, color coded according to the value of Y . The gray lines denote the equilibria, with the dashed line denoting the unstable equilibrium branch. Transition from the lower to the upper equilibrium branch for $H = -0.25$, $\tau = 350$. (b) Plot of Y as a function of time. Note how the peaks correspond to the brightly colored points in (a).

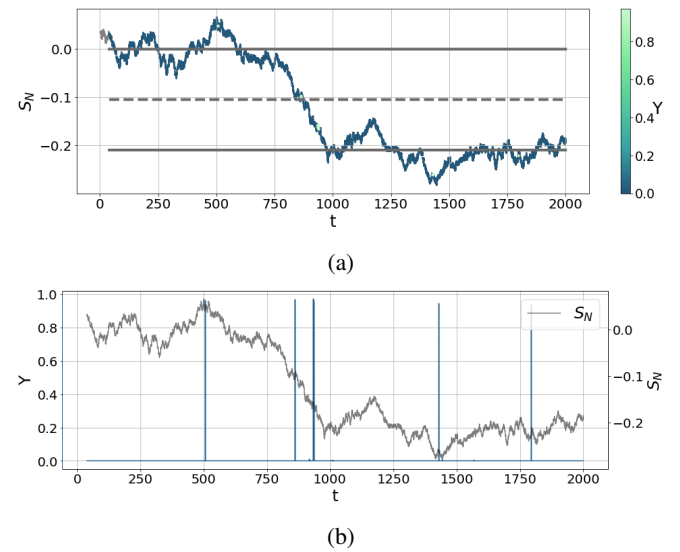


FIG. 9: (a) Noise-induced tipping, color coded according to the value of Y . The gray lines denote the equilibria, with the dashed line denoting the unstable equilibrium branch. Transition from the upper to the lower equilibrium branch for $H = 0.24$, $\tau = 200$. (b) Plot of Y as a function of time. Note how the peaks correspond to the brightly colored points in (a).

569 q -values. This is not unexpected, as it is primarily the change
570 in the properties of the noise which is expected to give an
571 indication of an approaching transition. Figure 7 shows the
572 persistence plotted as a function of time t . We see a clear in-
573 crease in the persistence directly preceding the tipping point
574 around $t = 1000$.

575 We make a final comment regarding Figure 6 and its relation
576 to our choice of ARMA(1,0) and ARMA(0,0) as base models.
577 In Faranda *et al.*¹⁰ this choice was guided by the fact that for
578 the time series under consideration the order, i.e. $p + q$, of the
579 intervals was clustered around 1, and as the authors explicitly
580 excluded pure moving-average processes, they concluded that
581 ARMA(1,0) was the appropriate base model. However, from
582 Figure 6 we see that for the time series currently under con-
583 sideration, the order is clustered around 0. This observation
584 further strengthens the case for using ARMA(0,0) as an ad-
585 ditional base model. We hypothesize that the dominance of
586 ARMA(0,0) is related to the low degree of noise in the sys-
587 tem, which makes the restoring force that returns the system
588 to equilibrium less prominent, hence obscuring tendency of
589 the random-walk to be attracted to a metastable state.

590 B. Noise-induced Tipping

591 To induce N-tipping, we fix the hosing parameter H and
592 apply additive white noise to all the equations equally. The
593 noise term is added equally to (10)-(13), with the same noise
594 amplitude in all cases. We look at transitions from the upper
595 branch to the lower branch and *vice versa*. In either case,

it is convenient to choose a value for H that is close to the
bifurcation point, as the probability of transitioning is much
higher in these regions, and hence one does not need high
amplitude noise to induce transitions between the branches.

Figures 8 and 9 show two time series undergoing noise in-
duced tipping, one going from the lower to the upper branch,
while the other going the other way around. In the first case
 $H = -0.25$, while in the second $H = 0.24$. The amplitude
of the additive white noise is the same in both cases. For the
window length τ , we have chosen a length of 350 and 200
points, corresponding to about 70 and 41 years, respectively.
The window length is chosen so that it is at most half as long
as the transition time, which is taken to be the time for the
system to arrive at the other equilibrium once it has crossed
the unstable branch. Of course, when dealing with simulation
data such as this, we have the advantage of knowing where
the stable and unstable branches are, which is an advantage
that anyone dealing with real-world data does not have. In
principle one could use the clustering methods proposed by
Kaiser *et al.*¹³ to approximate the window length, although
this method also requires that one knows how many clusters,
i.e., equilibrium states, one should look for. The clustering
method works particularly well for noise induced transitions,
as one can repeatedly induce transitions back and forth, to
gain an ensemble of transitions, yielding a higher degree of
accuracy.

In previous works, the choice of τ has largely been guided by
a desire to ensure the stationarity of the time series intervals.
However, as we are not requiring the individual time series
segments to be stationary *a priori*, we are permitted to use

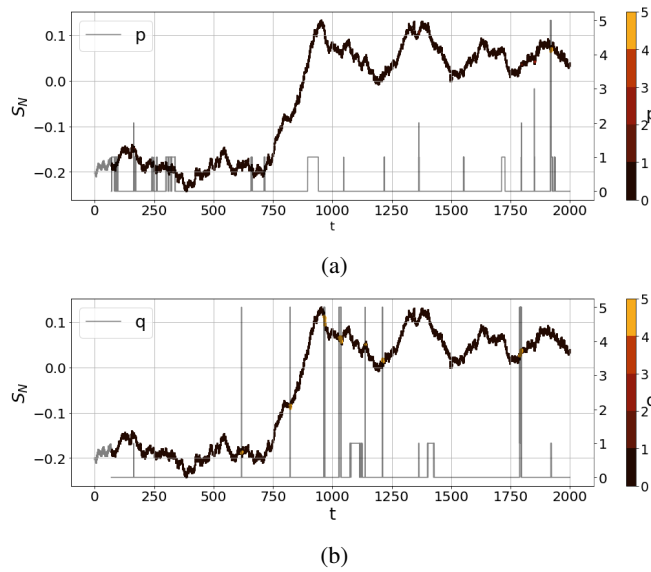


FIG. 10: Noise-induced tipping of $S_N(t)$ for $H = -0.25$, $\tau = 350$, color coded according to the value of (a) p and (b) q . For clarity we have also plotted p and q as functions of time in (a) and (b), respectively.

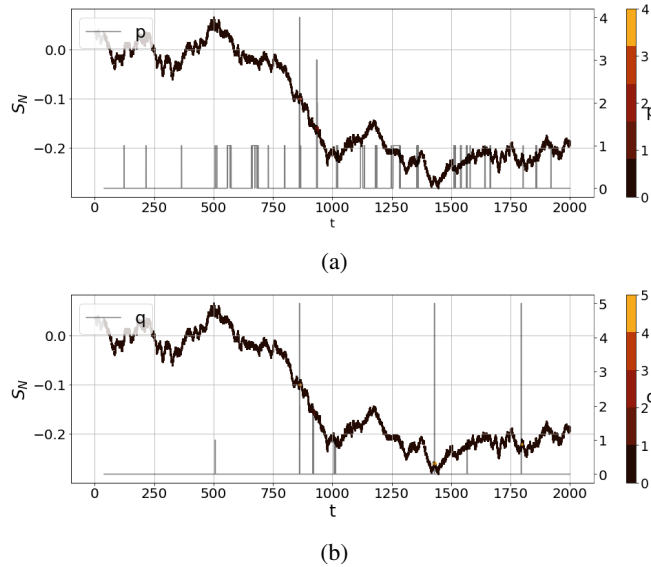


FIG. 11: Noise-induced tipping of $S_N(t)$ for $H = 0.24$, $\tau = 200$, color coded according to the value of (a) p and (b) q . For clarity we have also plotted p and q as functions of time in (a) and (b), respectively.

much longer time series intervals. In the world of ARIMA fitting a time series of length above 200 points would generally be considered a very long series, however, we should keep in mind that the sampling frequency of our simulated data is quite high; in fact, there are 5 points per time unit (i.e. year), yielding a total of 10000 points for the 2000 years of simulations. An interval consisting of 200 points corresponds to around 40 years, which is not an unreasonably long time

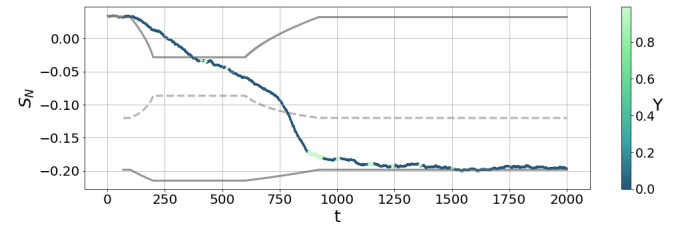


FIG. 12: Rate-induced tipping of S_N , color coded according to the value of Υ . The moving equilibria are plotted in gray, with the dashed line denoting the unstable branch. Compare this figure to Figure 14a, which shows the same time series, but color coded according to the value of q .

interval for the dynamics of the AMOC. When fitting an ARIMA model to a time series, one wishes to avoid too long time series to avoid including events from the past that no longer have any relevance for the future. This, and not the inherent inaccuracy of the fit itself, is the primary reason for limiting the length of a time series.

Returning to Figures 8 and 9, we note that there are a few brightly colored points indicating a high degree of instability. There are for example, in both cases, several points in the middle of the gap between the two stable branches, indicated by solid gray lines in the figure. This is consistent with the results of Kaiser *et al.*¹³. In addition, for the transition from the lower to the upper branch, Figure 8, there are several brightly colored points just after the system has reached the upper equilibrium branch. Although it is not so clear in the figure due to the presence of noise, any time S_N returns to the upper equilibrium branch it initially overshoots and then oscillates around the equilibrium value with continuously decreasing amplitude (see Figure 13 for a clearer example of this behavior). This is probably due to the presence of an unstable limit cycle, and the aforementioned sub-critical hopf bifurcation. Hence, we see it as an encouraging sign that the indicator seems to be able to identify these points as well. We further note that, although the result is not shown, the high Υ value points in figure 8 and 9 correspond to points where $\Delta_1 \text{BIC}(p, q)$ is negative, as was the case for the B-tipping example in the previous section.

Looking at the p and q values in Figures 10 and 11, it is clear that high values of Υ correspond to high values of q , while the connection between p and Υ remains uncertain. However, we note that the high Υ values appearing around the transition correspond to high values of both p and q , and consequently also of persistence (result not shown).

C. Rate-induced Tipping

To induce R-tipping we fix H_{pert} below the bifurcation value, ensuring that both equilibria still exist and are stable, and vary T_{fall} . We set $T_{rise} = 100$ and $T_{pert} = 400$, while $H_{pert} = 0.37$. This corresponds to an increase in the freshwater fluxes F_T and F_N , corresponding to the flux into the tropical

The Υ indicator for Early Warning

10

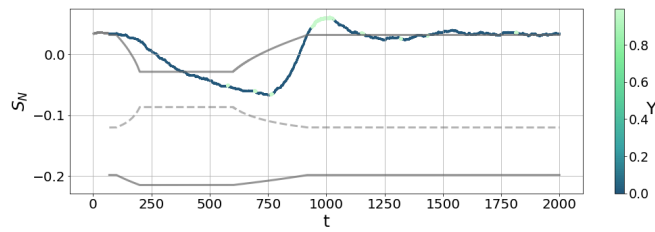
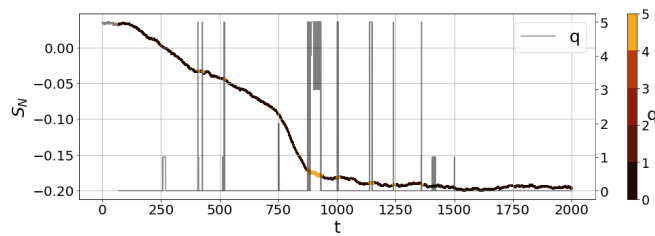
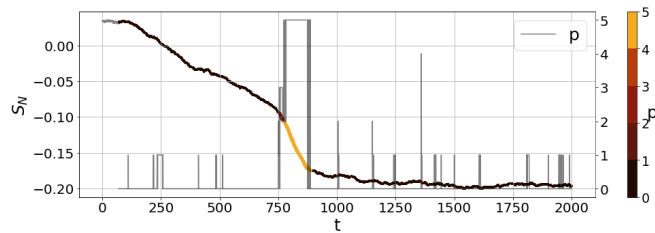


FIG. 13: S_N as a function of time, color coded according to the value of Υ for $T_{fall} = 280$. With these parameter values, the system does not tip, but returns to the upper equilibrium branch after some time. Note that the system initially overshoots the stable branch upon return. This is probably due to the presence of the unstable limit cycle. The equilibrium branches are plotted in gray, with the dashed line denoting the unstable branch.

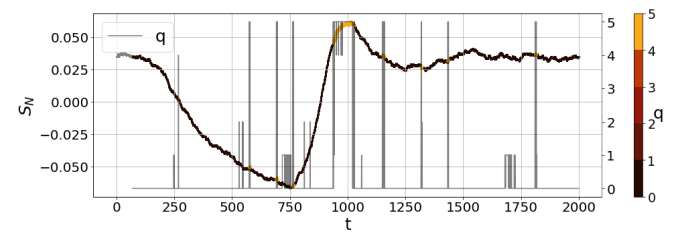


(a)

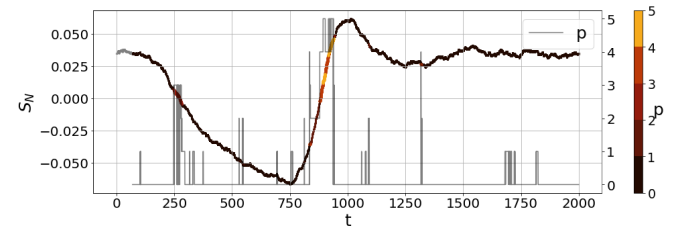


(b)

FIG. 14: Rate-induced tipping of $S_N(t)$, color coded according to the value of (a) q and (b) p . The value for q and p are also plotted as functions of time in (a) and (b), respectively.



(a)



(b)

FIG. 15: S_N as a function of time, color coded according to the value of (a) q and (b) p , for $T_{fall} = 280$. For these parameter values, the system does not tip, but returns to the initial equilibrium after some time t . For clarity, p and q are also plotted as functions of time in (a) and (b), respectively. It is instructive to compare these plots to Figure 13.

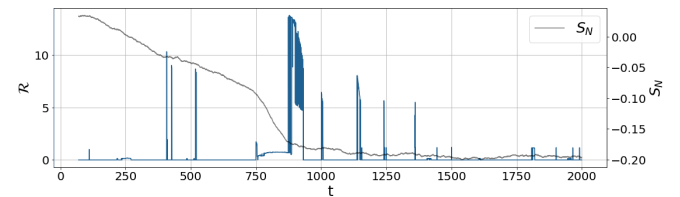


FIG. 16: Persistence of a time series undergoing rate-induced tipping, plotted as a function of time. The underlying series is the time series shown in Figure 12. We see several high persistence values, corresponding with a high value for the order, $q + p$ (compare with Figure 14), appearing before the potential tipping point around $t = 500$.

674 and North Atlantic boxes, by approximately 25% and 10%, re-690
675 spectively. Next, we observe that for $T_{fall} = 280$ the system 691
676 returns to the upper equilibrium branch, while for $T_{fall} = 320$, 692
677 the system transitions to the lower branch. The transition 693
678 happens even though the bifurcation boundary has not been 694
679 crossed. Again, we note that some additive white noise has 695
680 been applied to allow for ARIMA fitting. 696
681 Figure 12 shows a time series undergoing rate-induced tip-697
682 ping, with the color coding corresponding to the values of Υ . 698
683 Again, we have chosen $\tau = 350$ points, corresponding to 700
684 years. We see several brightly colored points, indicating a 700
685 high degree of instability, before the system transitions. These 701
686 points occur initially as the system approaches the unstable 702
687 branch (between approximately $t = 350$ and $t = 500$). These 703
688 points do not appear for the time series that does not tip, Fig-704

ure 13, despite the fact that within this time interval, the two
time series are virtually identical, and could therefore be an
indication of an approaching tipping point. However, again
looking at Figure 13 we see some brightly colored points, cor-
responding to large Υ , in the interval $t = 600$ to $t = 750$, and it
is unclear what approaching instability these points would be
indicative of, and thus might be regarded as false signals.

Looking at Figure 14, it becomes clear that the high values of
 Υ found in Figure 12 correspond to high values of q , while a
comparison with Figure 16, gives the same indication for the
persistence. In other words, high values of Υ primarily corre-
spond to high values of persistence and q .

From Figure 13, we can also see how the indicator correctly
identifies the unstable limit cycle, which we have argued
causes the overshoot when returning to the upper equilibrium
branch. Figure 15 shows the same time series as in Figure 13,

705 color coded according to the values of q and p . While high
706 values of q seem to be associated with increased instability,
707 the high values of p primarily occur as the system returns to
708 the equilibrium. We would therefore suggest that high val-
709 ues of the autoregressive order, p , should be interpreted as an
710 indication that the system is following a moving equilibrium
711 branch. Comparing Figures 16 and 14a it becomes clear that
712 the points with high q value around $t = 1000$, correspond to
713 particularly high values of persistence, even when compared
714 to other points of similar order. We also note that, as in the
715 previous two tipping scenarios, the high Υ values, or equiva-
716 lently high p values,

717 We end this section with a brief comment on the rate-induced
718 tipping example presented in this section. In this example the
719 system is, as it undergoes rate-induced tipping, approaching a
720 bifurcation boundary. It would be instructive to study a case in
721 which this is not the case to ensure that the detected instabil-
722 ity is not merely due to the approaching bifurcation boundary.
723 However, as one would need to look at different model exam-
724 ples than those presented here, this is outside the scope of the
725 current work.

726 IV. COMPARISON WITH OTHER EARLY WARNING 727 INDICATORS

728 As briefly alluded to in the introduction, it is well estab-
729 lished that bifurcation-induced tipping is generally preceded
730 by an increase in lag 1 autocorrelation and variance (Lenton
731 *et al.*³³, Dakos *et al.*³⁴, Boers²⁴). The intuition behind
732 this is that as the system approaches a bifurcation point,
733 the potential well flattens out, reducing the speed at which
734 the system recovers from a perturbation, so called "critical
735 slowing down", which should manifest as an increase in the
736 variance and autocorrelation of the time series. However, the
737 variance and autocorrelation might also increase for other
738 reasons, in particular if the properties of the noise changes.
739 What happens to the autocorrelation and variance when the
740 system approaches a rate-induced tipping point is thus far
741 unclear, although it is conceivable that the "critical slowing
742 down" hypothesis still holds for this type of tipping, see
743 Ritchie and Sieber⁹. Obviously, it does not hold true for time
744 series undergoing purely noise induced tipping, as there is no
745 change in the potential well. However, the autocorrelation
746 and variance of the time series will dramatically change as
747 the system crosses the unstable equilibrium branch and enters
748 a different potential well.

749 In what follows, we will compare these classical indicators
750 to the Υ indicator for rate-induced and bifurcation-induced
751 tipping in the AMOC 3-box model. It is instructive to just
752 look at the part of the time series prior to the transition,
753 in general one wishes to be able to detect early signs of the
754 transition *before* it happens. For the time series undergoing
755 bifurcation-induced tipping (Figure 4) we chose a segment
756 consisting of the points between approximately $t = 200$ and
757 $t = 1100$. For the time series undergoing rate-induced tipping
758 (Figure 13), we choose a segment consisting of the points be-
759 tween $t = 200$ and $t = 700$. This segment is in all probability

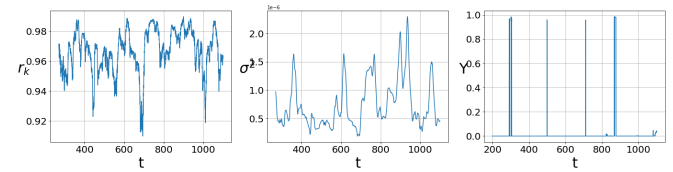


FIG. 17: Autocorrelation, Variance and Υ plotted as functions of time for a time series undergoing B-tipping. The increase in the variance as one approaches the tipping point is clear, while the increase in autocorrelation is less clear.

too long, meaning that it also contains the transition itself, as opposed to only points prior to the transition. However, this is the inherent difficulty with rate induced tipping; there is currently no way to analytically determine *when* the transition happens, and one largely has to guess. Based on Figures 12 and 13, one could potentially conclude that the tipping point is found somewhere between $t = 400$ and $t = 600$, but this is pure guess work. For this reason we have included points up until $t = 700$.

770 Given a set of measurements Y_1, Y_2, \dots, Y_N the sample
771 variance is defined as

$$\sigma^2 = \frac{1}{N} \sum_{i=1}^N (Y_i - \bar{Y})^2 \quad (19)$$

while the lag k autocorrelation is given by

$$r_k = \frac{1}{N\sigma^2} \sum_{i=1}^{N-k} (Y_i - \bar{Y})(Y_{i+k} - \bar{Y}) \quad (20)$$

where \bar{Y} denotes the sample mean of the series Y_1, Y_2, \dots, Y_N (see for example chapter 2 of Box, Jenkins, and Reinsel³⁵). Although time does not enter explicitly in the formulas, it is assumed that the measurements are taken at regular intervals.

When computing the variance and autocorrelation it is essential that the signal is properly detrended; otherwise any trend will immediately obscure the relevant dynamics. As for the Υ indicator, one generally employs a rolling window approach, with an appropriately chosen window length τ . Lenton *et al.*³³ demonstrated that detrending can be done within each time window, as opposed to on the whole time series at once, without significantly changing the result. We have chosen this same approach, using linear detrending, as opposed to quadratic or higher order detrending methods, to remove the trend. The window length τ was set to 350 points, corresponding to 70 years.

Figures 17 and 18 show the autocorrelation, variance and Υ plotted as functions of time. The peaks in Υ preceding the transition are clear, as is the increase in variance and autocorrelation, at least in the case of R-tipping, provided the tipping point is approximately at $t = 450$. For B-tipping, there appears to be a clear increase in the variance preceding the tipping point, provided the tipping point happens around

The Υ indicator for Early Warning

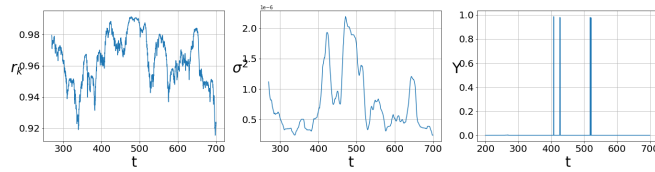


FIG. 18: Autocorrelation, Variance and Υ plotted as functions of time for a time series undergoing R-tipping. Assuming that the tipping point is around $t=450$, one can clearly see an increase in both autocorrelation and variance prior to the tipping point.

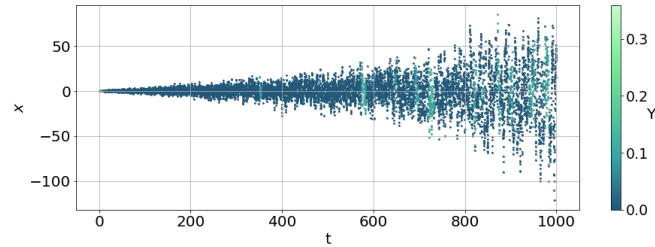


FIG. 19: Time series with colored noise but no tipping points, color coded according to the value of Υ .

801 $t = 850$ (see Figure 4 for comparison). The expected increase
802 in autocorrelation is, however, less clear.
803 It is possible that the high degree of autocorrelation in the
804 3-box model, as observed in Figures 17 and 18 is correlated
805 to the frequent failure of the ARMA(1,0) model, whereby
806 failure we mean that the autoregressive coefficient, sometimes
807 referred to as the AR1 coefficient, violates the stationarity
808 condition, and resulting in ARMA(1,0) being excluded as a
809 possible candidate model.
810 As already noted, the upper equilibrium branch does not
811 lose stability due to a saddle node bifurcation, but rather
812 loses stability due to a sub-critical Hopf bifurcation. It
813 is possible that classical indicators are struggling to pick
814 up on this. Furthermore, the noise amplitude is kept low
815 to avoid noise-induced tipping, which might make it dif-
816 ficult for the indicators to pick up on changes in the dynamics.

817
818 The autocorrelation and variance of a time series can in-
819 crease for reasons that have nothing to do with an approaching
820 tipping point. Hence, we wish to see how the Υ indicator re-

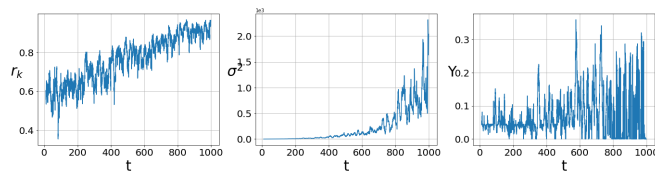
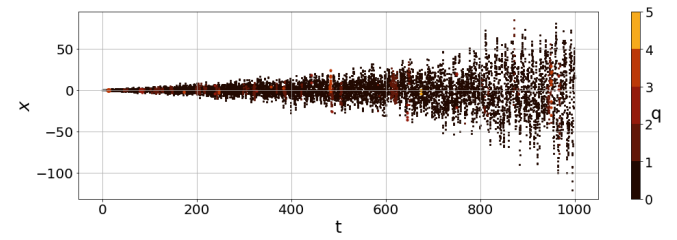
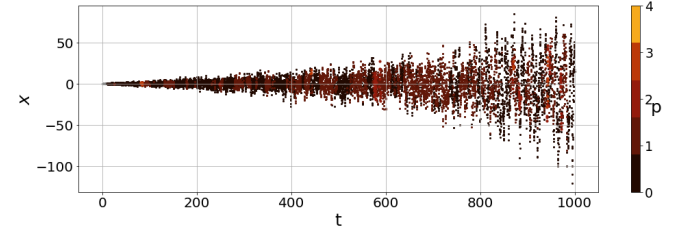


FIG. 20: Autocorrelation, Variance and Υ plotted as functions of time for a time series with colored noise but no tipping points. All three indicators show a dramatic increase, falsely suggesting an upcoming tipping point.



(a)



(b)

FIG. 21: Time series with colored noise and no tipping points, corresponding to equation (21), color coded according to the value of (a) q and (b) p .

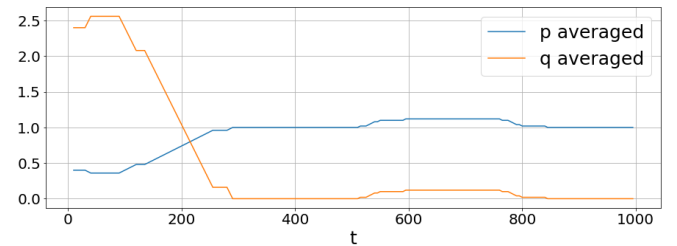


FIG. 22: The values of p and q for the colored noise time series, averaged with a window length of 50 points, corresponding to 25 non-dimensional time units.

sponds to colored noise, whose variance and autocorrelation increases with time t . To this end, we construct an artificial time series of the form

$$\frac{dx}{dt} = -5x + \xi(t) \quad (21)$$

where $\xi(t)$ is autocorrelated colored noise. $\xi(t)$ is in effect modelled as an ARMA(1,0) process whose AR1 coefficient increases linearly in time. In addition, the variance of this process also increases linearly in time. This is equivalent to the example presented in Boers²⁴. Applying the Υ indicator to this time series yields the result shown in Figure 19. Figure 20 shows a comparison between the autocorrelation, variance and value of Υ for the same time series. All three indicators show a dramatic increase, despite there being no approaching tipping point. However, looking at the plot of the time series when color coded according to the values of p and q , Figure 21, a curious pattern emerges: the increase in Υ is largely associated with increased p value. Looking at Figure 22 the trend becomes even clearer: here we have computed the rolling average of the p and q values with a window length

840 of 50 points corresponding to 25 non-dimensional time units.
 841 We see that while the average value of q goes towards zero
 842 for large t , the average value of p settles around one. The
 843 general trend is independent of the choice of window length,
 844 provided the window length is between 30 and 300 points.
 845 This behavior is unlike what was observed for the 3-box
 846 model. The high values of Υ were associated with a high value
 847 of q . We thus argue that high values of q were associated
 848 with increased instability, while high values of p were more
 849 indicative of the system following a moving equilibrium.
 850 Thus, one would, through the distinction between q and p
 851 values, potentially have a way of distinguishing the effect of
 852 colored noise from real early warning signals. However, it
 853 is conceivable that the result for the artificial colored noise
 854 time series is a consequence of how we have constructed the
 855 colored noise, so further studies on this are warranted.

856
 857 Finally, we note that the constructed colored noise time
 858 series is a very artificial example of colored noise, as the
 859 noise amplitude increases by a probably unrealistic amount,
 860 and when applied to any reasonable time series it would
 861 obscure the dynamics altogether. This is to say that although
 862 we can likely assume that the noise in real-world data is
 863 autocorrelated, it will be much more subtle, and not result in
 864 equally high values of Υ .

865 V. APPLICATION TO SIMULATION DATA FROM CESM2

866 So far, we have only applied the dynamic stability indica-
 867 tor to data from a very simplified model. The actual ocean
 868 has many more degrees of freedom and the response could
 869 be quite different. Nevertheless, it is of interest to see how
 870 the indicator responds when applied to such a system. To this
 871 end, we employ data from the earth systems model CESM2
 872 under two climate scenarios: one in which the atmospheric
 873 CO₂ concentration is abruptly doubled and another in which
 874 it is abruptly quadrupled. Both simulations were initialized
 875 using a pre-industrial control run (*piControl*) and then run for
 876 500 years. The CO₂ was then increased, at $t = 6000$ months.
 877 The data was saved at monthly intervals and the seasonal cy-
 878 cle was removed prior to the analysis. Such an abrupt change
 879 in CO₂ represents an extreme forcing, and contrasts with the
 880 ramped-up hosing employed with the idealized model. How-
 881 ever, the oceanic response is not instantaneous, but requires 2-
 882 3 decades for freshwater to circulate in the model's sub-polar
 883 gyre²⁷. We consider this more hereafter.

884 A. Abrupt $4 \times \text{CO}_2$

885 The time series of a monthly-mean density difference, $\delta\rho$,
 886 and AMOC strength, ψ_{AMOC} , are shown in Figure 23 for the
 887 case of abrupt $4 \times \text{CO}_2$. The density difference, a measure dy-
 888 namically linked to the AMOC strength (Madan *et al.*²⁷), is
 889 calculated from the difference in surface densities averaged in
 890 boxes to the north and south of the North Atlantic Current.
 891 The surface density is calculated using the thermodynamic

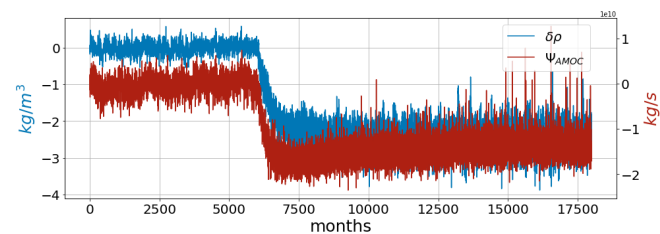


FIG. 23: CESM2 model with abrupt $4 \times \text{CO}_2$, where the monthly density difference (blue) is plotted together with the maximum AMOC flow strength (red). Note that the CO₂ was increased at $t=6000$ months.

equation of state of seawater as per UNESCO 1983 Report³⁶. The AMOC strength is calculated as the monthly maxima of meridional overturning stream function between 20°N - 60°N and below 450 m depth.

Shortly after the quadrupling of CO₂, there is an abrupt transition followed by a dramatic increase in the variance. We will apply the indicator to the density difference time series, although one could of course apply the same analysis to the AMOC strength.

We choose a window length of 250 data points, corresponding to exactly 20 years of monthly data. Figure 24 shows the density difference, $\delta\rho$, color coded according to the values of Υ . We only display the part of the time series close to the transition, as this is of primary interest. The point at which the CO₂ concentration is abruptly increased, at $t = 6000$ months, is indicated by a dashed line.

The increase in Υ during the early part of the AMOC weakening process is apparent. Note in particular the three sharp peaks shortly after time $t=6000$. Figure 25 again shows the time series, now color coded according to the values of q and p . The latter are also plotted for further clarification. From this plot, it becomes clear that the most common fit prior to the transition is the ARMA(1,0) process, which aligns with the observations of Faranda *et al.*¹⁰. After the weakening phase, the value of p is generally an order higher, presumably related to the dramatic increase in the variance. The three sharp peaks in the plot of Υ appearing around time $t = 6300$ correspond to high values of q . The gradual increase in Υ preceding these peaks is presumably due to the increase in the persistence (not shown). The q component exhibits peaks prior to $t = 6000$, when the forcing is applied and these are reflected in small peaks in Υ . These are obviously not connected to the AMOC weakening. Following the initial weakening phase, the value for Υ remains high, probably a result of the increase in the p value. However, the values of Υ do not go above 0.4 which is considerably smaller than the values found for the 3-box model. In addition, from our previous discussion on the response of the Υ indicator to colored noise, it is conceivable that the increase in Υ observed from in the CESM2 data is primarily caused by changes in the noise amplitude, and not as a consequence of inherent instability of the underlying dynamics.

Furthermore we note that, although the result is not explicitly shown, for the CESM2 data ΔBIC_1 is always smaller than

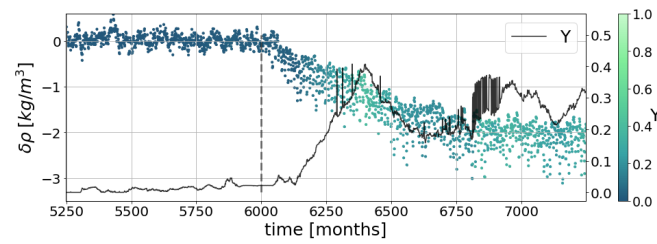
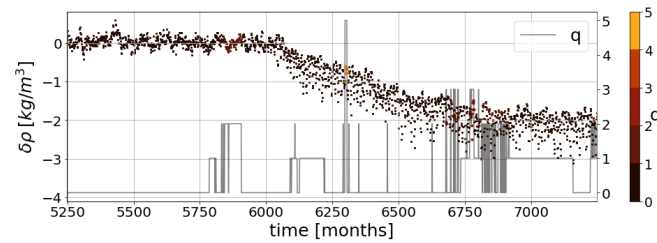
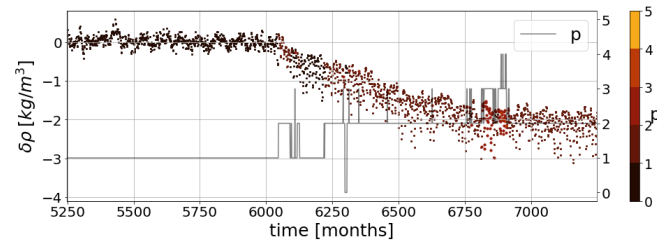


FIG. 24: Time series of monthly density changes for abrupt $4 \times \text{CO}_2$, color coded according to the value of Y . The window length is 250 points, corresponding to exactly 20 years. The dashed line indicates the point when the CO_2 concentration abruptly changes.



(a)



(b)

FIG. 25: Time series of monthly density changes for abrupt $4 \times \text{CO}_2$, color coded according to the value of (a) q and (b) p . The value for q and p are also plotted as functions of time in (a) and (b), respectively.

ΔBIC_0 , and the ΔBIC_1 values are at no point negative, implying that the autoregressive coefficient in the ARMA(1,0) model always satisfy the stationarity constraints. This differs from what was observed in the 3-box model and is presumably related to the difference in the observed Y values. However, we emphasize that it is not clear if one can compare values of Y between datasets. For the autocorrelation and the variance it is typically assumed that it is the change *within* the dataset that is significant, rather than the absolute numerical values. For completeness, we have included a comparison between Y and two other statistical early warning indicators, namely autocorrelation and variance. This is shown in Figure 26. In all cases, the window length is 250 points, corresponding to approximately 20 years. All three indicators show a clear increase shortly after time $t = 6000$.

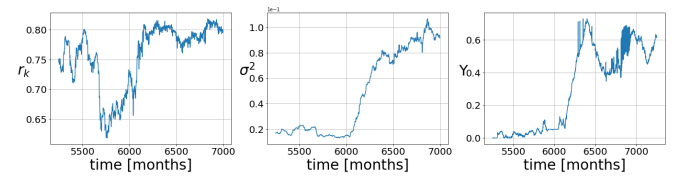


FIG. 26: Autocorrelation, variance and Y plotted as functions of time for the case of abrupt $4 \times \text{CO}_2$.

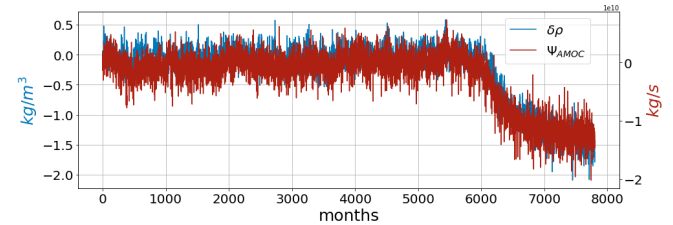


FIG. 27: CESM2 model with abrupt $2 \times \text{CO}_2$, where the monthly density difference (blue) is plotted together with the maximum AMOC flow strength (red).

B. Abrupt $2 \times \text{CO}_2$

The time series of the monthly density difference, $\delta \rho$, and AMOC strength, Ψ_{AMOC} , in the case of abrupt $2 \times \text{CO}_2$ is shown in Figure 27. Again, we only apply the indicator to the density difference data, and choose the same window length as in the case of abrupt $4 \times \text{CO}_2$. Figure 28 shows an excerpt of the density difference time series close to the initial weakening, as well as a plot of the Y values. A weakening is clearly seen in the model's own AMOC measure, and is also accurately captured with the measure based on the density difference across the Gulf Stream (Fig. 27). The first thing to note is how small the Y values are compared to what we have seen previously; on the order of 10^{-2} . It should, however, be noted that the ΔBIC values are well above the significance threshold²⁹. Figure 29 shows the density difference time series color coded according to the value of q and p . From this, we again see that prior to the increase in CO_2 , the most common fit is the ARMA(1,0) process, while after the initial weakening phase the p values show a clear increase. The q value, on the other hand, does not exceed 2, indicating a very low degree of memory in the noise term. Since we have by now clearly demonstrated a correlation with the value of Y and the value of q , this should provide an explanation as to why we see such low values of Y . From this analysis, one would conclude the system does not appear to be approaching a tipping point. Indeed, the measure suggests that the weakening in the overturning in this case with reduced forcing is not associated with a loss of dynamical stability. Once more we have, as shown in Figure 30, included a comparison with other early warning indicators. The autocorrelation and variance show a dramatic increase around time $t=6000$, which corresponds to the appearance of the cluster of sharp peaks in the time series plot for Y .

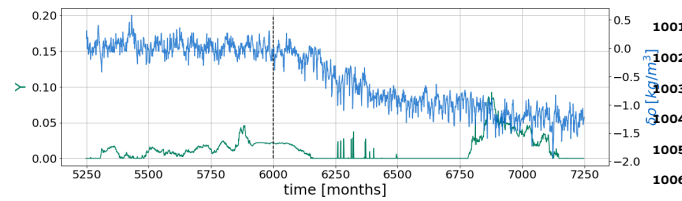


FIG. 28: Monthly density changes, $\delta\rho$, for abrupt $2 \times \text{CO}_2$ (blue) and the value of Υ (green) plotted as functions of time. The dashed line indicates the point when the CO_2 concentration abruptly changes.

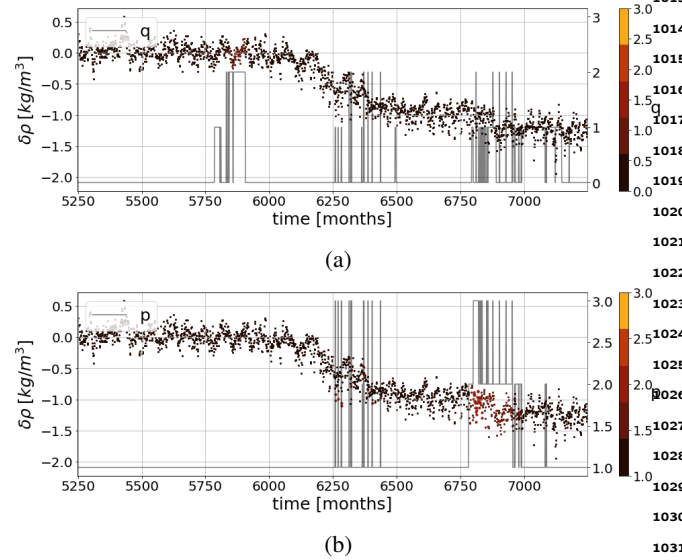


FIG. 29: Time series of monthly density changes for abrupt $2 \times \text{CO}_2$, color coded according to the value of (a) q and (b) p . The value for q and p are also plotted as functions of time in (a) and (b), respectively.

993 VI. DISCUSSION

994 In summary, we analysed an indicator for dynamical
 995 stability based on ARMA modelling as a way to detect
 996 transitions in complex systems. A detected need for higher
 997 order terms in the ARMA model fitted to moving windows of
 998 a timeseries is related to diverging memory properties, which
 999 are expected to arise when approaching a transition to a new
 1000 equilibrium state. The rationale behind this indicator is that

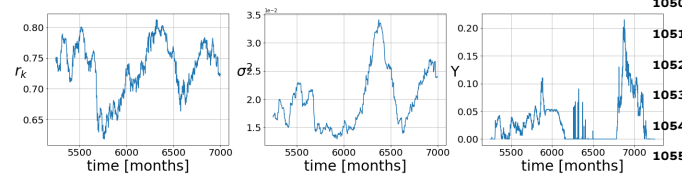


FIG. 30: Autocorrelation, variance and Υ plotted as functions of time for the case of abrupt $2 \times \text{CO}_2$

it uses a broad family of linear statistical models that can be fitted even on short time series and which have proven their utility in many contexts (see Brockwell and Davis²⁸). That the underlying models do not require long time series is an advantage when employing a sliding window approach on limited data sets. The method generalizes classical metrics of instability, and allows one to extract more global dynamical information from the time series data.

The indicator was tested on time series data from a 3-box model of the AMOC, where three categories of critical transitions, namely B-, N-, and R-tipping, were explored. In all cases the transition is identified by the indicator, albeit it is not always easy to interpret the signal. In the rate-induced tipping scenario a comparison between the avoided tipping and the tipping cases shows a response of the indicator prior to the transition only in the tipping case although the time series are nearly identical at this stage. The indicator also successfully identifies the unstable limit cycle when returning to the upper equilibrium branch. We similarly see fairly clear signals in the bifurcation-induced tipping scenario prior to the transition. For the case of noise-induced tipping, the signal is less clear, obscured by the high amplitude noise. However, when going from the lower to the upper equilibrium branch the indicator signals an increased degree of instability in accordance with the presence of the unstable limit cycle.

The primary drawback of the Υ indicator is that it is computationally quite expensive, at least compared to the autocorrelation and variance, and that, due to its complexity, the results can be harder to interpret. We therefore suggest that the indicator should be applied with care, and preferably in combinations with other measures of instability, like the increase in the order, $p + q$, and the persistence. Although the current scaling with τ , see equation (3), seems to yield reasonable results, it is certainly possible that another scaling would be preferred. It is also possible that this is problem-dependent. This uncertainty regarding the correct scaling is certainly a drawback, but we argue that this problem can largely be circumvented by including an examination of the persistence and order values. However, it would still be advantageous to have an indicator whose values were to have a clear meaning in terms of the stability of the system, and it is not clear if the Υ indicator as it stands achieves this, partly due to the aforementioned issue with the choice of the correct scaling. Although we have attempted to make some comparison to other early warning indicators, like the increase in autocorrelation and variance, we are not claiming that the Υ indicator is in any way better than these other indicators, rather that it can act as a complementary approach, as it can allow one to extract more information from time series data. For example, we have suggested, that it might be helpful in identifying the effects of colored noise, something the other indicators struggle with.

Furthermore, we note that it is conceivable that one would wish to exclude white noise and pure moving-average, MA(1), processes when doing the fitting, as was done in the

1059 earlier studies by Faranda *et al.*¹⁰. In such a scenario the
 1060 modified definition of the Υ indicator would of course no
 1061 longer be valid, as the ARMA(0,0) process is excluded, and
 1062 thus cannot be used as a base model. In this case one might
 1063 argue that the points where $\Delta_1 \text{BIC}$ are negative should either
 1064 be ignored completely, or one should assume that the best fit
 1065 is in fact the ARMA(1,0) process and the algorithm is being
 1066 too strict in its enforcement of the auxiliary conditions on
 1067 the fitting parameters. This would of course lead to different
 1068 results than what has been presented here, and is an option
 1069 worth considering.

1070
 1071 When considering a full complexity AMOC model as
 1072 arising from a global climate model (CESM2) many more
 1073 degrees of freedom are involved. This has two consequences
 1074 firstly, the pure categories of tipping cannot really be expected
 1075 anymore and secondly, the tipping behaviour might disappear
 1076 altogether as the added degrees of freedom may stabilize the
 1077 system.

1078 When applied to the CESM2 data, the results were mixed.
 1079 The measure exhibited a significant increase in Υ under the
 1080 more severe $4\times\text{CO}_2$ forcing but much less variability with
 1081 the weaker $2\times\text{CO}_2$ forcing. Hence the measure only registers
 1082 larger changes in AMOC as associated with dynamically un-
 1083 stable behavior. Indeed, it is possible that the model AMOC
 1084 experiences a continuously shifting steady state, rather than
 1085 making a transition between two distinct states as in low-
 1086 dimensional models. The results from the doubling CO_2
 1087 experiment seems to support this hypothesis. Other members
 1088 of the CMIP6 ensemble exhibiting very different AMOC
 1089 weakening from the same forcing, with some declining by
 1090 only 15% and others falling by 80%²⁷, and this suggests a
 1091 continuum of different responses.

1092 While the results for $4\times\text{CO}_2$ suggest a loss of dynamical
 1093 stability during the AMOC weakening phase, concluding on
 1094 the tipping behaviour would require a more in depth analysis
 1095 along the lines done in Hawkins *et al.*²³; in this paper the
 1096 bi-stability is clearly demonstrated by exploring a range of
 1097 hosing experiments. Although we are confident that the
 1098 indicator can be used to assess the stability of such complex
 1099 systems, as was already demonstrated in previous works by
 1100 Nevo *et al.*¹², concluding on the ability to detect critical
 1101 transitions would require a full analysis of the hysteresis
 1102 behaviour of the system.

1103 ACKNOWLEDGMENTS

1104 This research has been partly funded by the Deutsche
 1105 Forschungsgemeinschaft (DFG) through grant CRC 1114
 1106 „Scaling Cascades in Complex Systems“, Project Number
 1107 235221301.
 1108 LaCasce was supported in part by the Rough Ocean project
 1109 number 302743, from the Norwegian Research Council.
 1110 The computations for CESM2 data were performed on re-
 1111 sources provided by Sigma2 - the National Infrastructure for
 1112 High Performance Computing and Data Storage in Norway.
 1113 The authors thank Davide Faranda for stimulating discussions

and Amandine Kaiser for help with the development of the
 code used for the numerical analysis.

DATA AVAILABILITY STATEMENT

The data that support the findings of this study are available
 from the corresponding author upon reasonable request.

- ¹T. M. Lenton, “Environmental Tipping Points,” Annual Review of Environ-
 ment and Resources **38**, 1–29 (2013), publisher: Annual Reviews.
- ²P. Ashwin, S. Wicczorek, R. Vitolo, and P. Cox, “Tipping points in
 open systems: bifurcation, noise-induced and rate-dependent examples
 in the climate system,” Phil. Trans. R. Soc. A **370**, 1166–1184 (2012),
 doi:10.1098/rsta.2011.0306.
- ³M. Scheffer, E. H. Van Nes, M. Holmgren, and T. Hughes, “Pulse-driven
 loss of top-down control: the critical-rate hypothesis,” Ecosystems **11**,
 226–237 (2008), doi:0.1007/s10021-007-9118-8.
- ⁴S. Wicczorek, P. Ashwin, C. Luke, and P. M. Cox, “Excitability in
 ramped systems: The compost-bomb instability,” Phil. Trans. R. Soc. A
467, 1243–126 (2011).
- ⁵P. E. O’Keeffe and S. Wicczorek, “Tipping phenomena and points of no
 return in ecosystems: Beyond classical bifurcations,” SIAM Journal on Ap-
 plied Dynamical Systems **19**, 2371–2402 (2020).
- ⁶S. Wicczorek and C. Perryman, “Adapting to a changing environment:
 Non-obvious thresholds in multi-scale systems,” Proc. R. Soc. A **470**,
 20140226 (2014).
- ⁷P. Ashwin, C. Perryman, and S. Wicczorek, “Parameter shifts for nonau-
 tonomous systems in low dimension: bifurcation- and rate-induced tip-
 ping,” Nonlinearity **30**, 2185–2210 (2017).
- ⁸A. Vanselow, S. Wicczorek, and U. Feudel, “When very slow is too fast-
 collapse of a predator-prey system,” Journal of theoretical biology **479**, 64–
 72 (2019), doi: 10.1016/j.jtbi.2019.07.008.
- ⁹P. Ritchie and J. Sieber, “Early-warning indicators for rate-induced tipping,”
 Chaos **26**, 093116 (2016), doi:10.1063/1.4963012.
- ¹⁰D. Faranda, F. M. E. Pons, E. Giachino, S. Vaienti, and B. Dubrulle, “Early
 warnings indicators of financial crises via auto regressive moving average
 models,” Communications in Nonlinear Science and Numerical Simulation
29, 233–239 (2015), doi:10.1016/j.cnsns.2015.05.002.
- ¹¹D. Faranda, B. Dubrulle, and F. M. E. Pons, “Statistical early-warning indi-
 cators based on autoregressive moving-average models,” Journal of Physics
 A: Mathematical and Theoretical **47**, 252001 (2014), doi:10.1088/1751-
 8113/47/25/252001.
- ¹²G. Nevo, N. Vercauteren, A. Kaiser, B. Dubrulle, and D. Faranda,
 “Statistical-mechanical approach to study the hydrodynamic stability of the
 stably stratified atmospheric boundary layer,” Phys. Rev. Fluids **2**, 084603
 (2017), doi:10.1103/PhysRevFluids.2.084603.
- ¹³A. Kaiser, D. Faranda, S. Krumscheid, D. Belusic, and N. Vercauteren,
 “Detecting regime transitions of the nocturnal and polar near-surface
 temperature inversion,” Journal of the Atmospheric Sciences, AMS **77**,
 2921–2940 (2020), doi:10.1175/JAS-D-19-0287.1.
- ¹⁴D. Faranda and D. DeFrance, “A wavelet-based approach to detect climate
 change on the coherent and turbulent component of the atmospheric cir-
 culations,” Earth Syst. Dynam. **7**, 517–523 (2016), doi:10.5194/esd-7-517-
 2016.
- ¹⁵H. Alkhuayon, P. Ashwin, L. C. Jackson, C. Quinn, and R. A. Wood,
 “Basin bifurcations, oscillatory instability and rate-induced thresholds for
 atlantic meridional overturning circulation in a global oceanic box model,”
 Proc. R. Soc. A **475** (2019), doi:10.1098/rspa.2019.0051.
- ¹⁶R. A. Wood, J. M. Rodriguez, R. S. Smith, L. C. Jackson, and E. Hawkins,
 “Observable, low-order dynamical controls on thresholds of the atlantic
 meridional overturning circulation,” Climate Dynamics **53**, 6815–6834
 (2019), doi:10.1007/s00382-019-04956-1.
- ¹⁷V. Masson-Delmotte, P. Zhai, A. Pirani, S. Connors, C. Péan, S. Berger,
 N. Caud, Y. Chen, L. Goldfarb, M. Gomis, M. Huang, K. Leitzell, E. Lon-
 noy, J. Matthews, T. Maycock, T. Waterfield, O. Yelekçi, R. Yu, and B. Z.

TABLE I: Adapted from Alkhayuon *et al.*¹⁵

	Volume	Salinity	Flux
North Atlantic	$V_N = 0.3683 \times 10^7 \text{ m}^3$	$S_N = 0.034912$	$F_N = 0.486 \text{ Sv}$
Tropical Atlantic	$V_T = 0.5418 \times 10^7 \text{ m}^3$	$S_T = 0.035435$	$F_T = -0.997 \text{ Sv}$
Southern Ocean	$V_S = 0.6097 \times 10^7 \text{ m}^3$	$S_S = 0.034427$	$F_S = 1.265 \text{ Sv}$
Indo-Pacific	$V_{IP} = 1.4860 \times 10^7 \text{ m}^3$	$S_{IP} = 0.034668$	$F_{IP} = -0.754 \text{ Sv}$
Bottom Ocean	$V_B = 9.9250 \times 10^7 \text{ m}^3$	$S_B = 0.034538$	

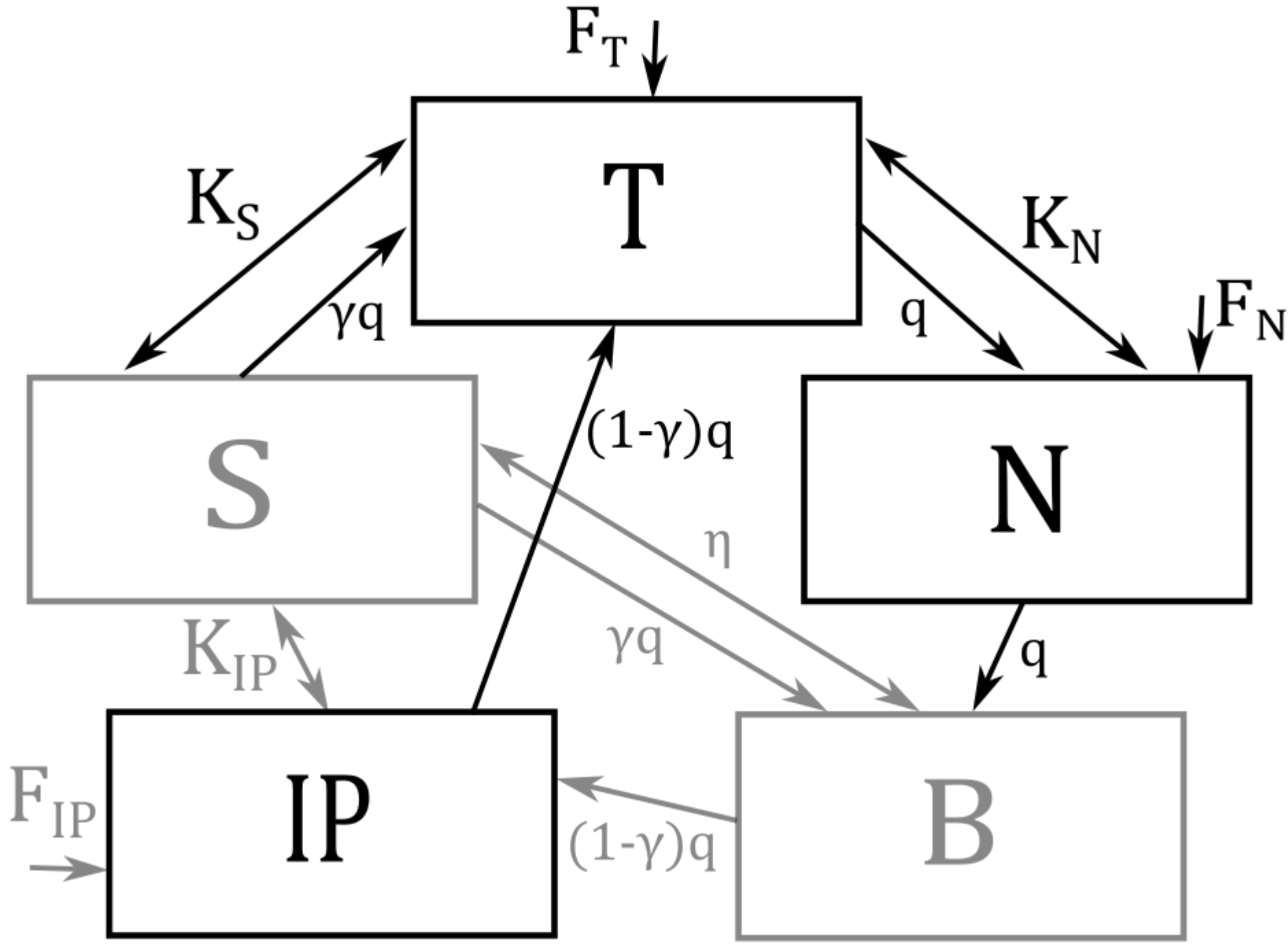
TABLE II: Adapted from Alkhayuon *et al.*¹⁵

name	default value	units	name	default value	units
α	0.12	kg/(m ³ °C)	K_N	1.762	Sv
β	790.0	kg/m ³	K_S	1.872	Sv
S_0	0.035		λ	1.62×10^7	m ⁶ /(kg s)
T_S	7.919	°C	γ	0.36	
T_0	3.870	°C			

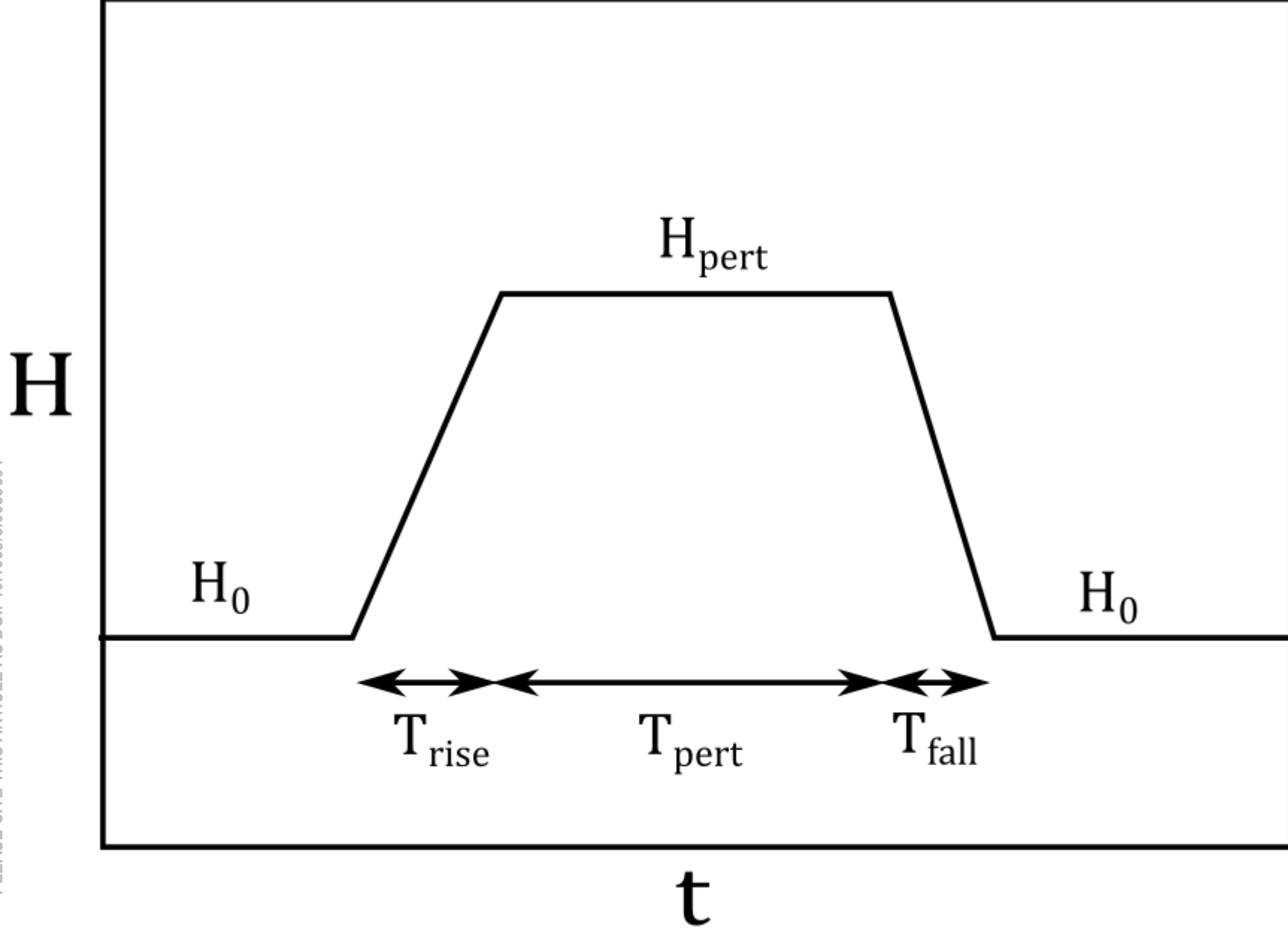
1178 (eds.), "IPCC 2021: Climate change 2021: The physical science basis. contribution of working group I to the sixth assessment report of the intergovernmental panel on climate change," Cambridge University Press (2021), doi:10.1017/9781009157896.
 1179
 1180
 1181
 1182
 1183
 1184
 1185
 1186
 1187
 1188
 1189
 1190
 1191
 1192
 1193
 1194
 1195
 1196
 1197
 1198
 1199
 1200
 1201
 1202

680–688 (2021).
 1203
 1204
 1205
 1206
 1207
 1208
 1209
 1210
 1211
 1212
 1213
 1214
 1215
 1216
 1217
 1218
 1219
 1220
 1221
 1222
 1223
 1224
 1225
 1226
 1227
 1228
 1229
 1230
 1231
 1232
 1233
 1234
 1235
 1236
 1237
 1238
 1239
 1240
 1241
 1242
 1243

25 J. Lohman and P. D. Ditlevsen, "Risk of tipping the overturning circulation due to increasing rates of ice melt," PNAS **118** (2021), doi:10.1073/pnas.201798911.
 26 G. Danabasoglu, "Near cesm2 model output prepared for cmip6 cmip," (2019).
 27 G. Madan, A. Gjermundsen, S. Iversen, and J. H. LaCasce, "Weakening of the atlantic meridional overturning circulation under extreme climate change." (2022), (In prep.).
 28 P. J. Brockwell and R. A. Davis, *Introduction to Time Series and Forecasting*, 2nd ed. (Springer, 2002).
 29 K. J. Preacher and E. C. Merkle, "The problem of model selection uncertainty in structural equation modeling," Psychological Methods **17**, 1–14 (2012), doi:10.1037/a0026805.
 30 R. J. Hyndman and Y. Khandakar, "Automatic time series forecasting: The forecast package for R," Journal of Statistical Software **27** (2008).
 31 J. M. T. Thompson and J. Sieber, "Climate tipping as a noisy bifurcation: a predictive technique," IMA Journal of Applied Mathematics **76**, 27–46 (2011).
 32 D. Roche, D. Paillard, T. Caley, and C. Waelbroeck, "Lgm hosing approach to heinrich event 1: results and perspectives from data-model integration using water isotopes," Quaternary Science Reviews **106**, 247–261 (2014).
 33 T. M. Lenton, V. N. Livina, V. D. Dakos, E. H. van Nes, and M. Scheffer, "Early warning of climate tipping points from critical slowing down: comparing methods to improve robustness," Phil. Trans. R. Soc. A **370**, 1185–1204 (2012).
 34 V. Dakos, S. R. Carpenter, W. A. Brock, A. M. Ellison, V. Guttal, A. R. Ives, S. Kefi, V. Livina, D. A. Seekell, E. H. van Nes, and M. Scheffer, "Methods for detecting early warnings of critical transitions in time series illustrated using simulated ecological data," PLoS ONE **7**, e41010 (2012), doi:10.1098/rsta.2011.0304.
 35 G. E. Box, G. M. Jenkins, and G. C. Reinsel, *Time Series Analysis: Forecasting and Control*, 4th ed. (WILEY SONS, INC., PUBLICATION, 2008).
 36 N. P. Fofonoff and R. Millard Jr, "Algorithms for the computation of fundamental properties of seawater." (1983).
 37 P. D. Ritche, J. J. Clark, P. M. Cox, and C. Huntingford, "Overshooting tipping point thresholds in a changing climate," Nature **592**, 517–523 (2021), doi:10.1038/s41586-021-03263-2.
 38 V. Laitinen, V. Dakos, and L. Lahti, "Probabilistic early warning signals," Ecology and Evolution **11**, 14101–14114 (2021), doi:10.1002/ece3.8123.



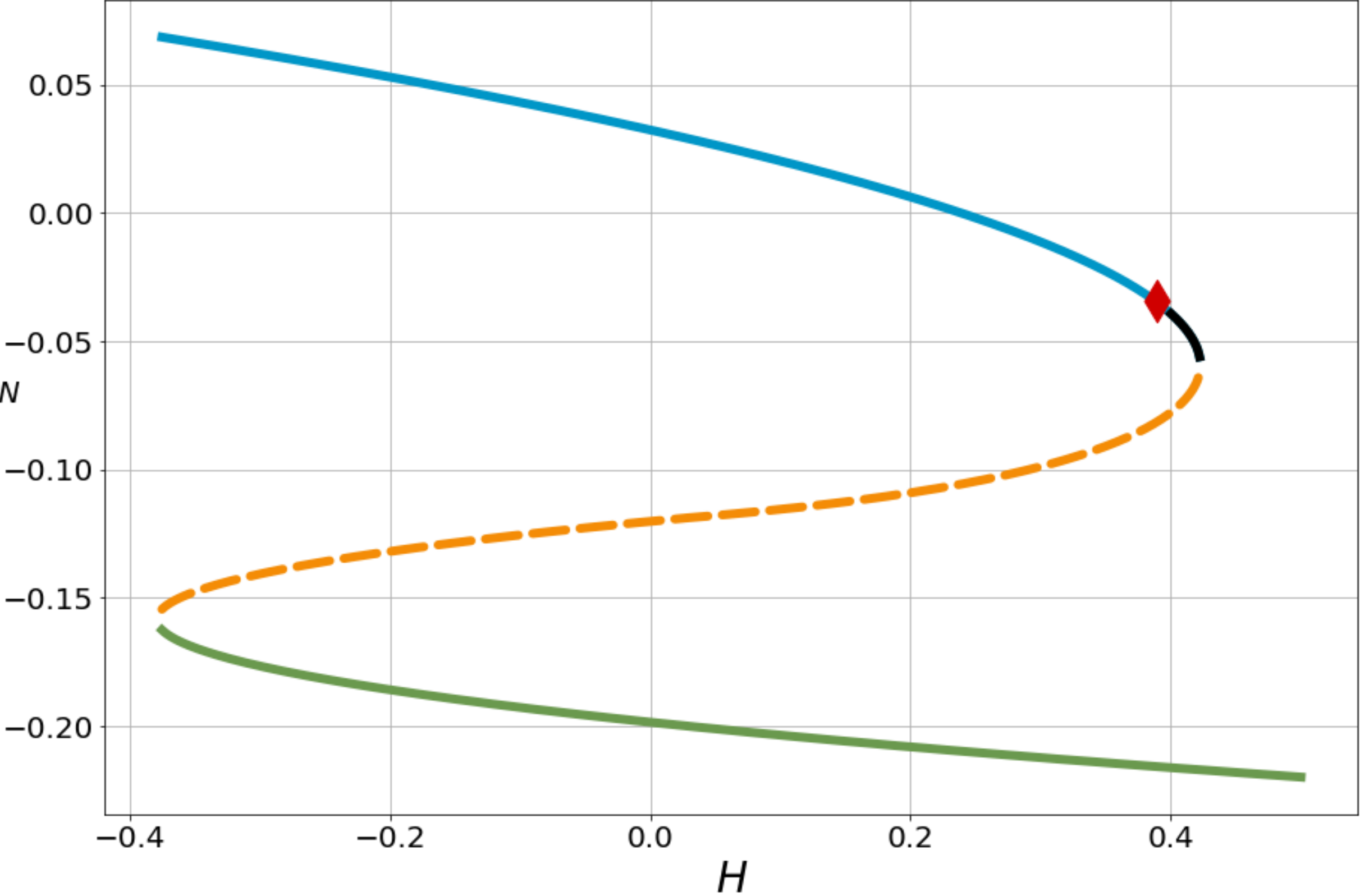
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



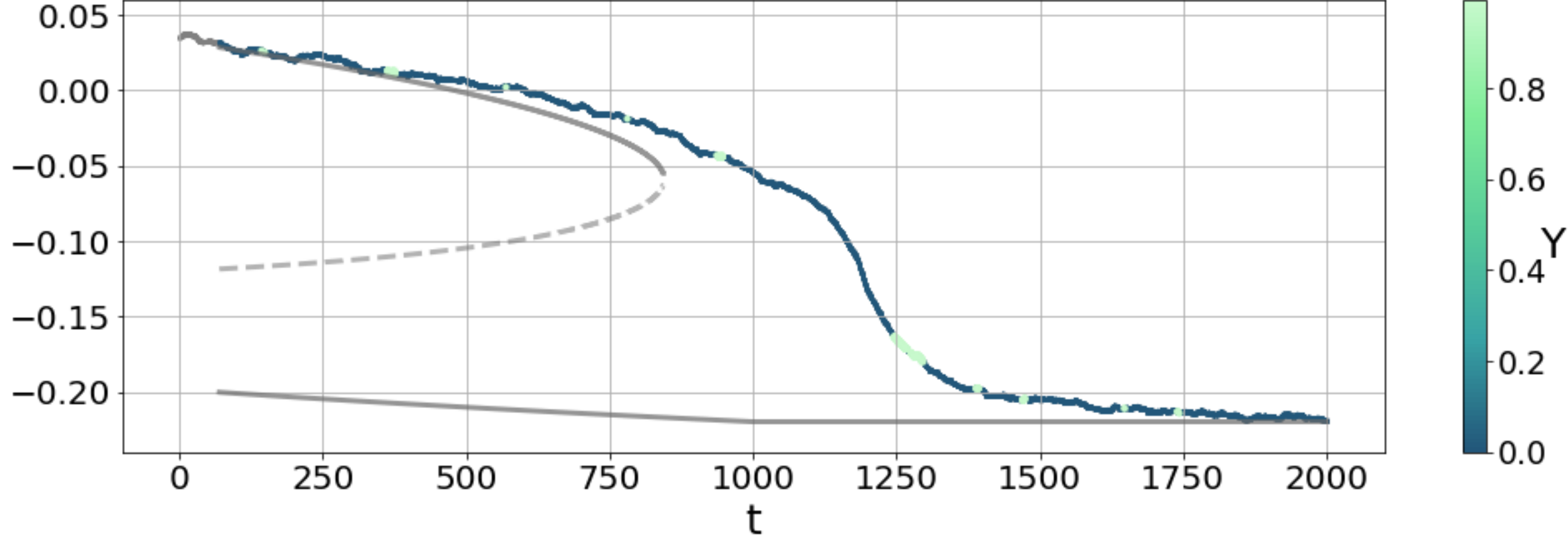
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

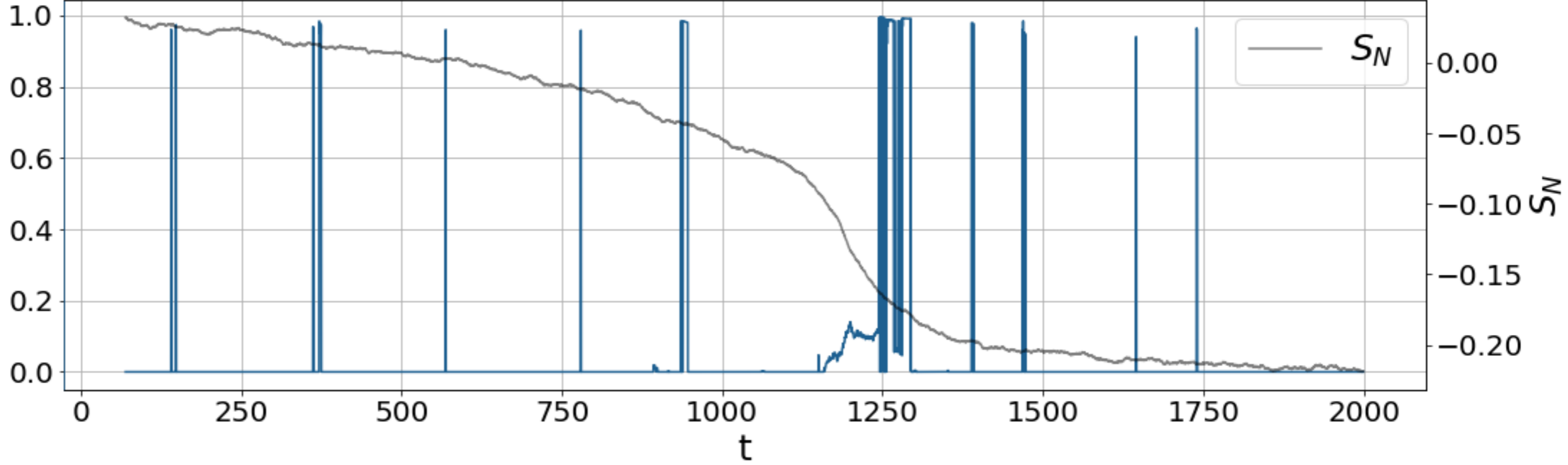
S_N



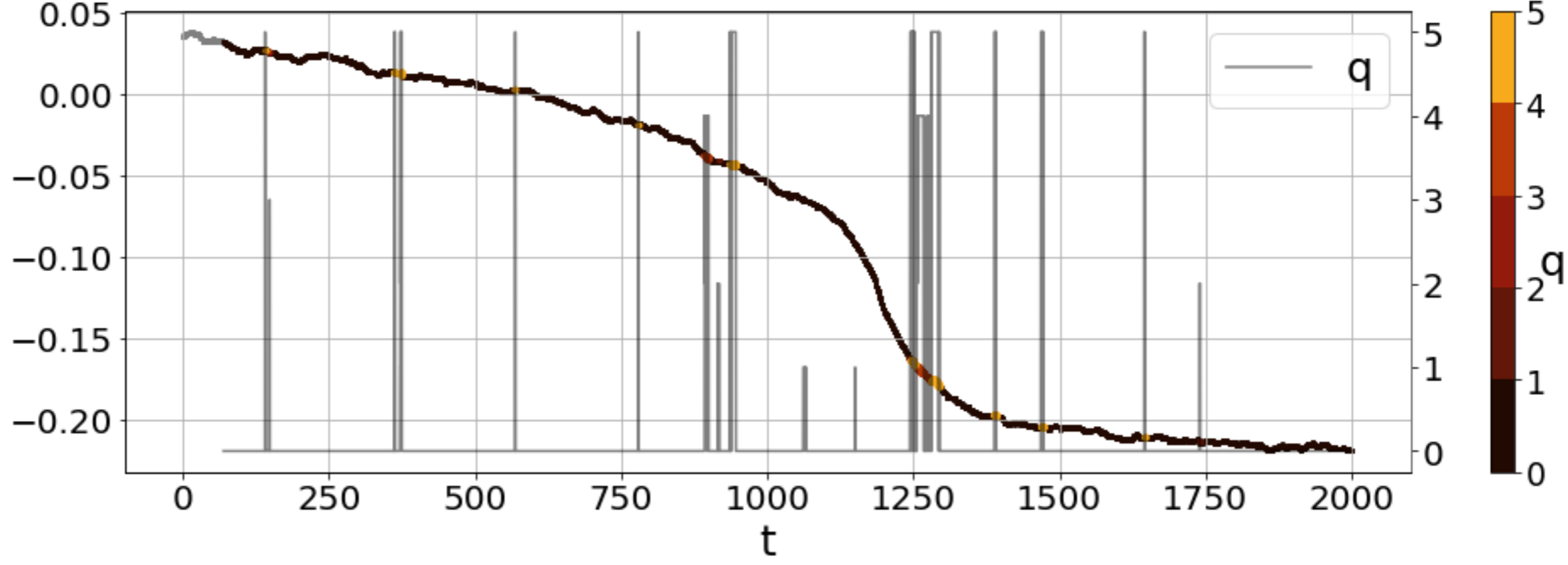
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



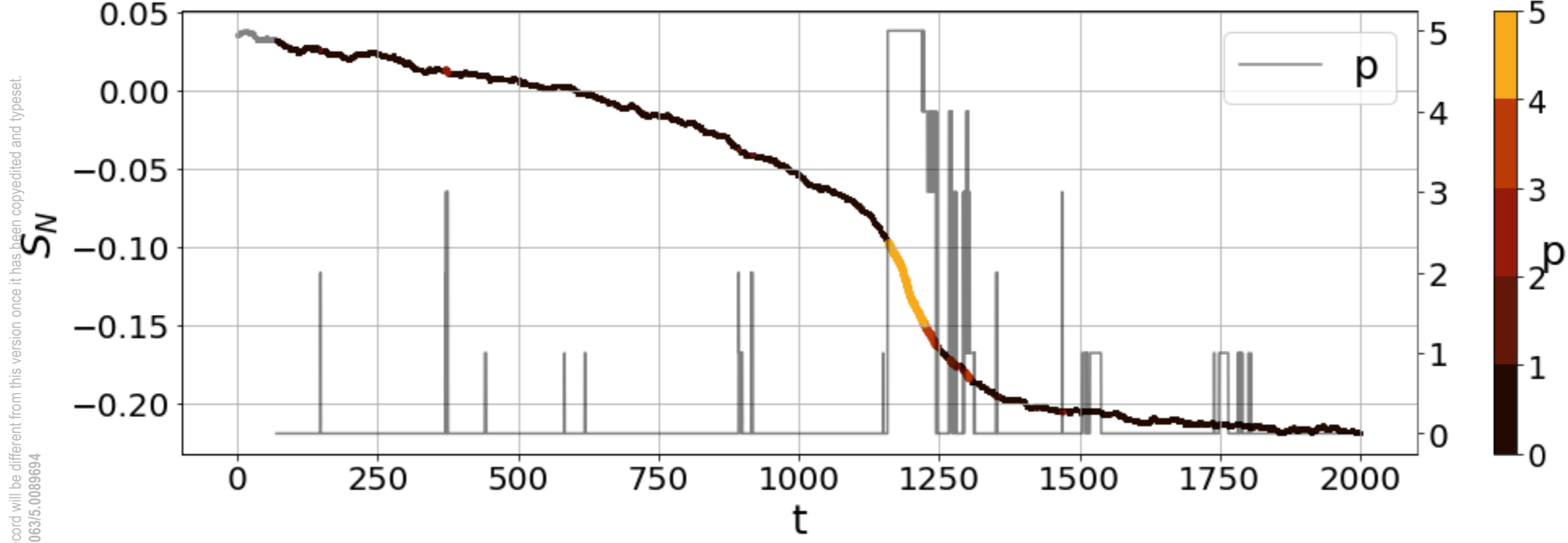
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/1.50089694



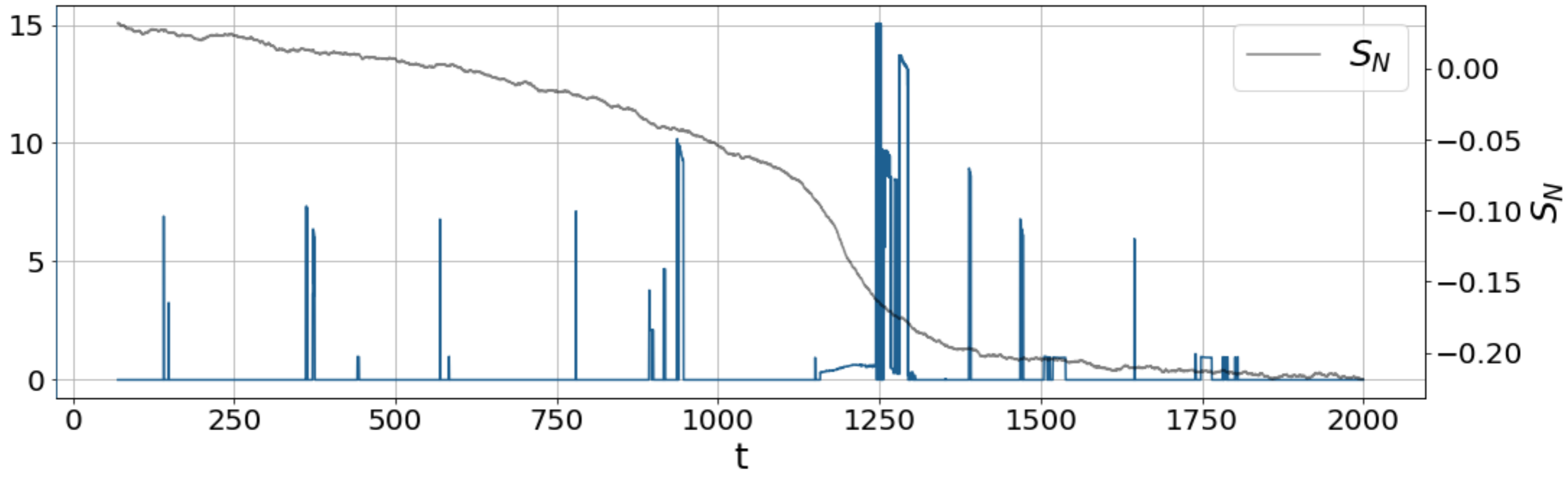
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



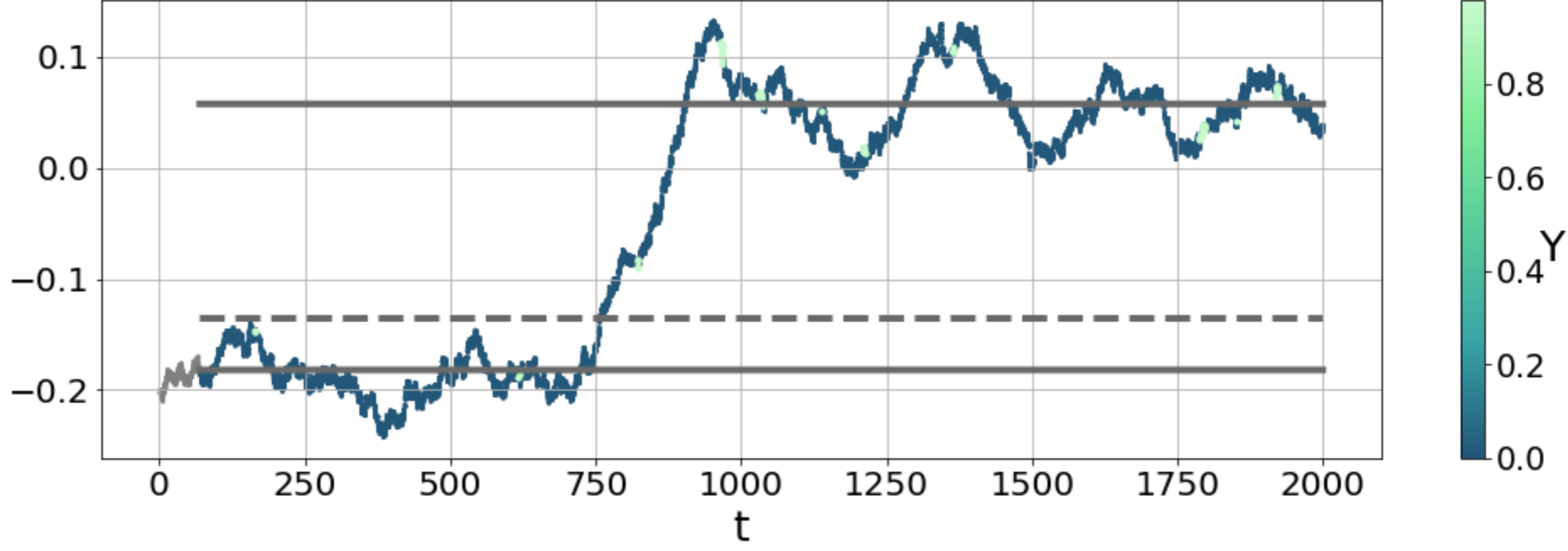
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



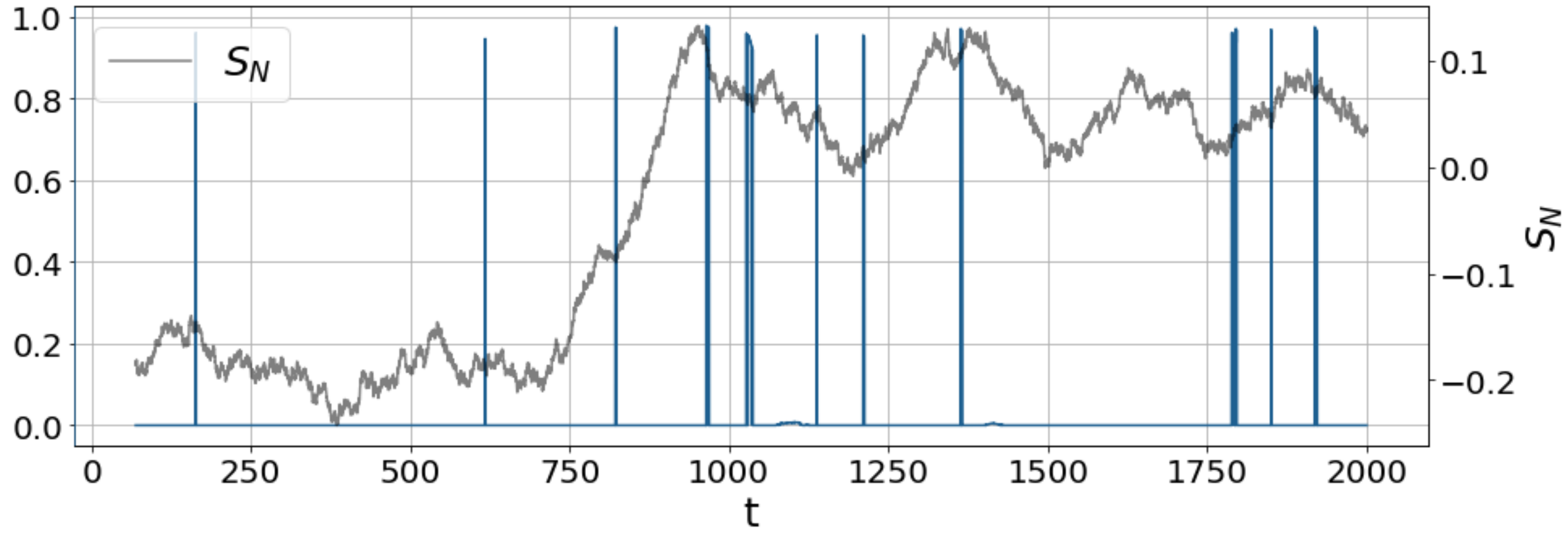
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/1.5008964



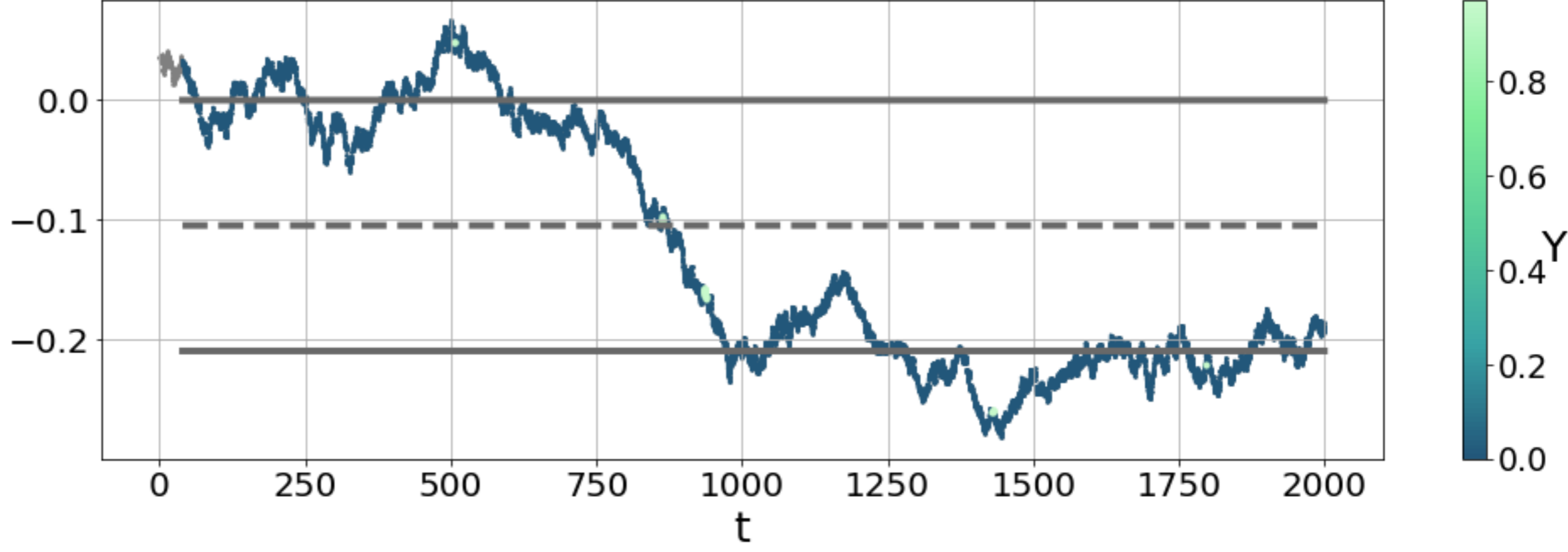
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



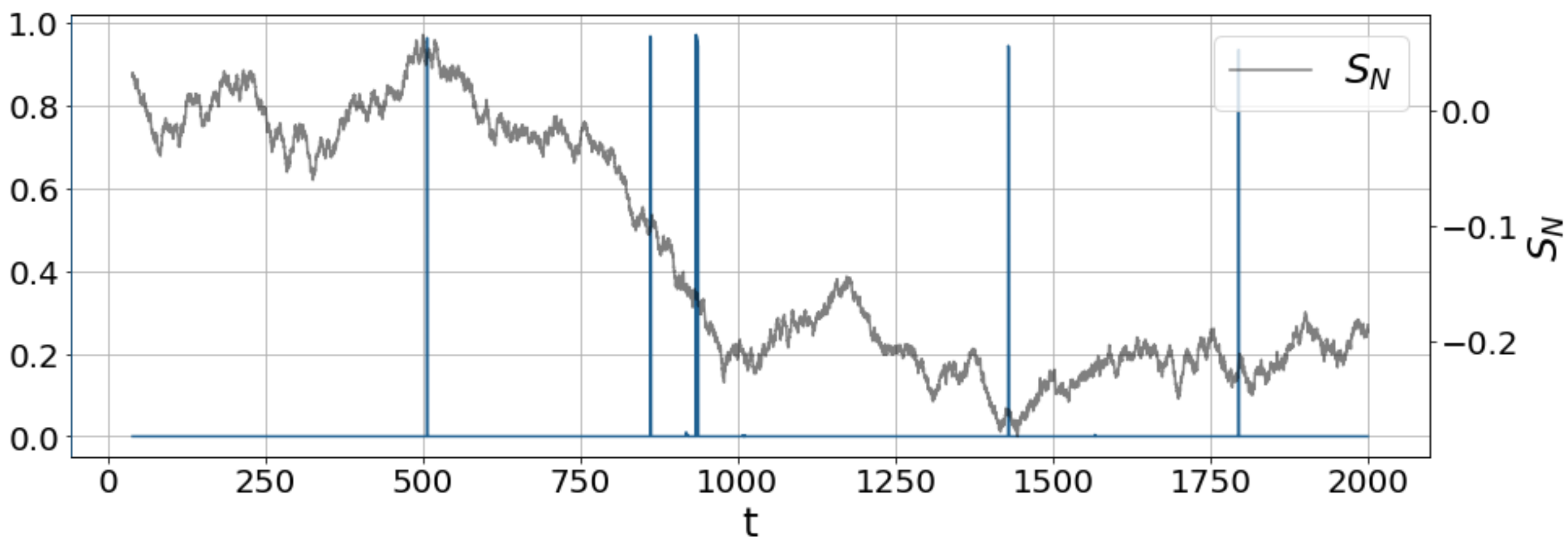
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



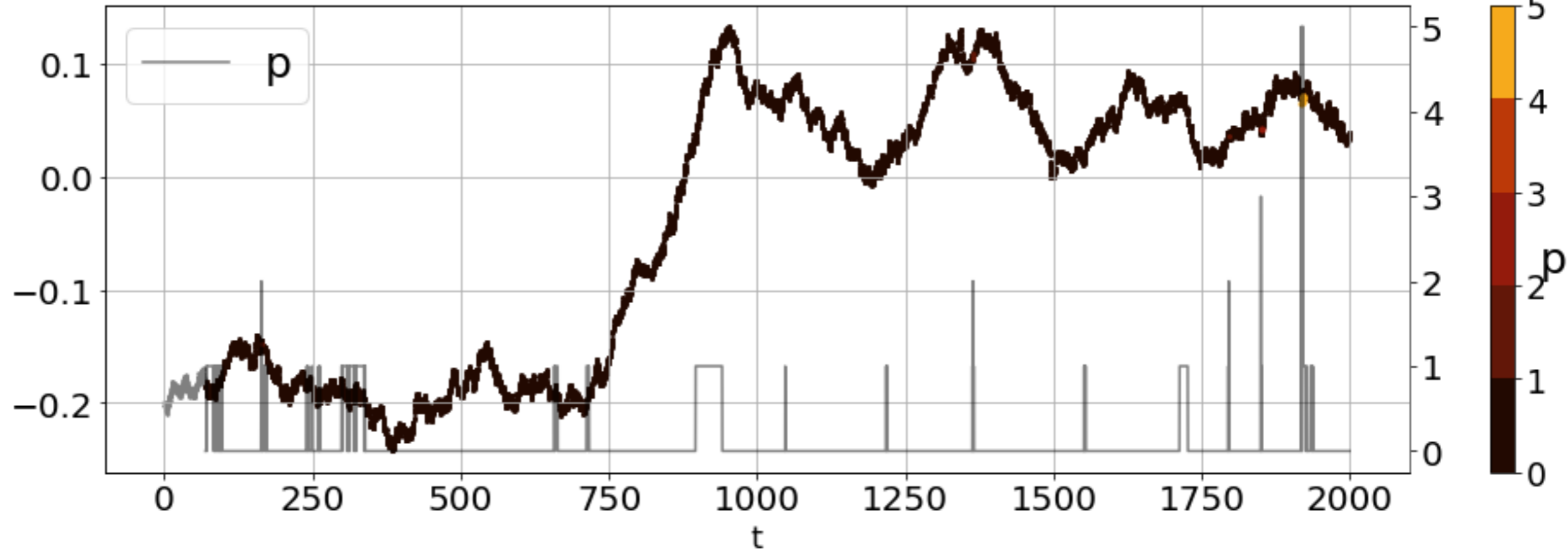
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



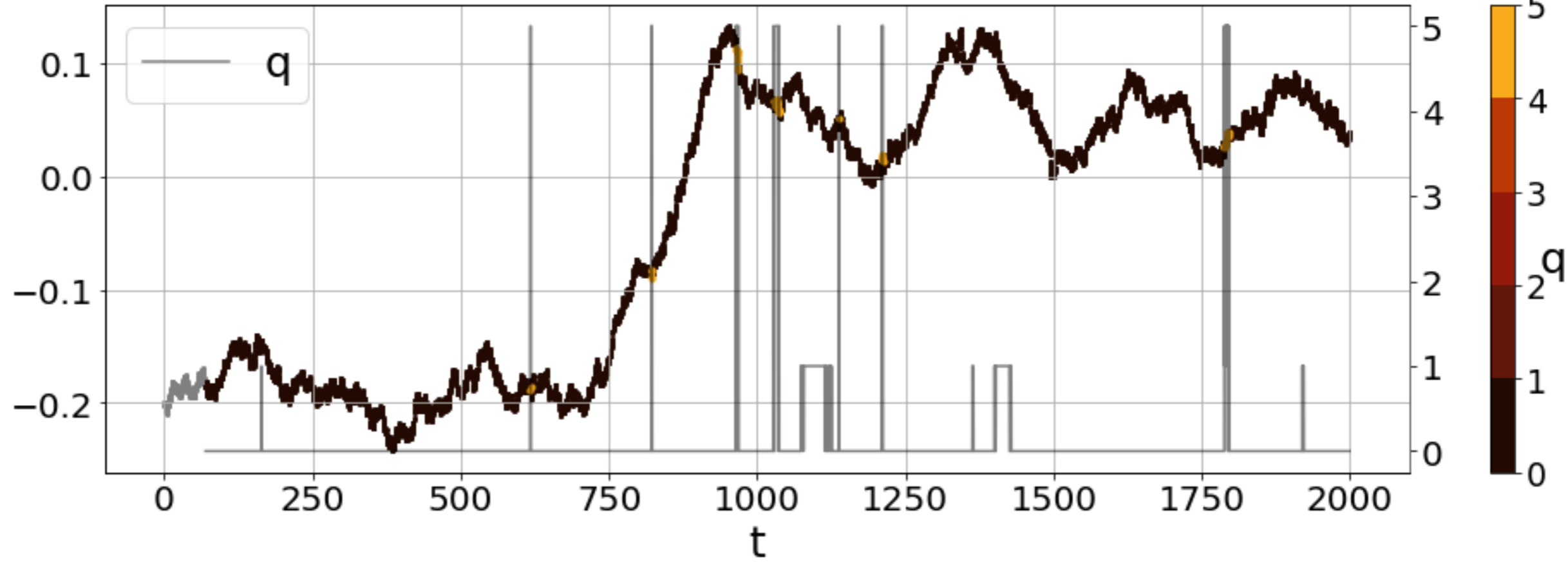
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



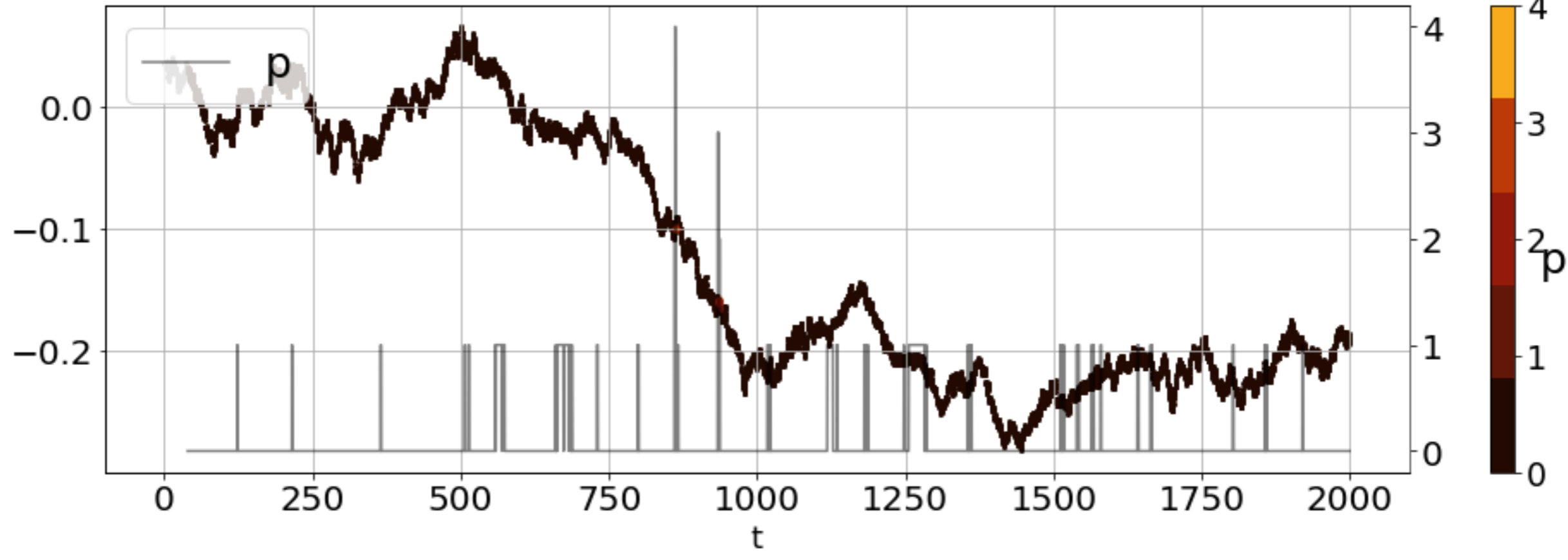
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



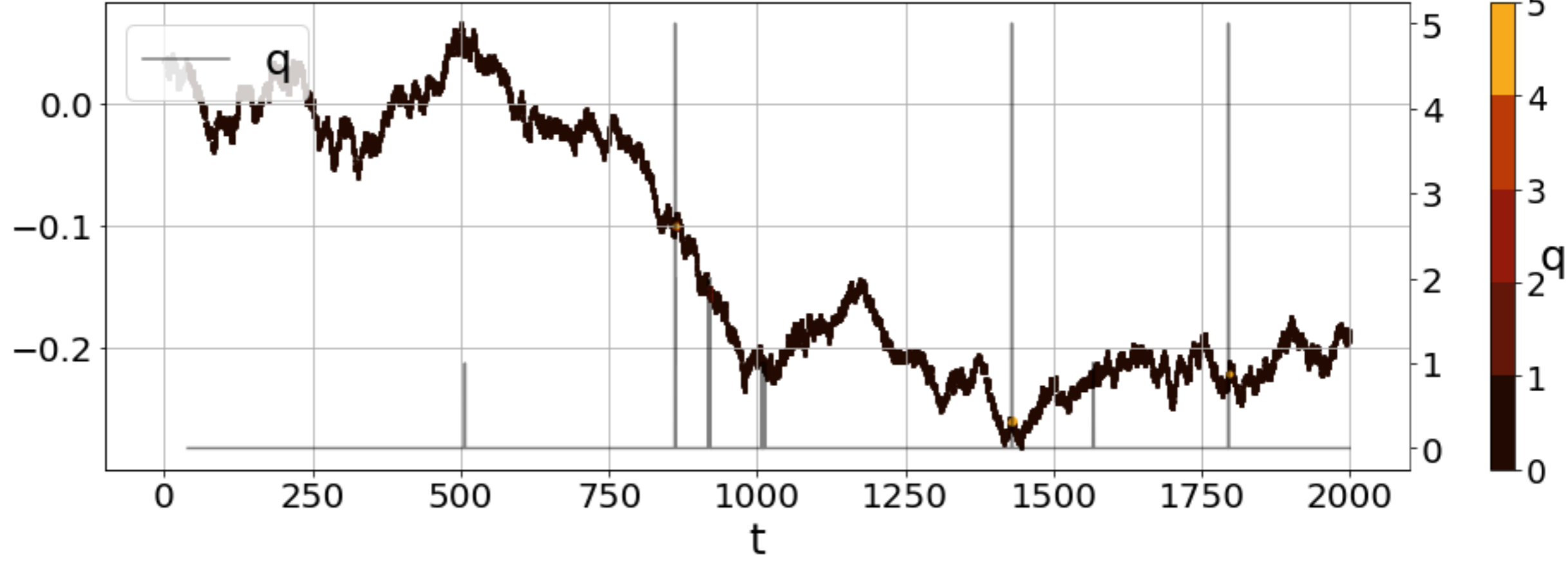
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



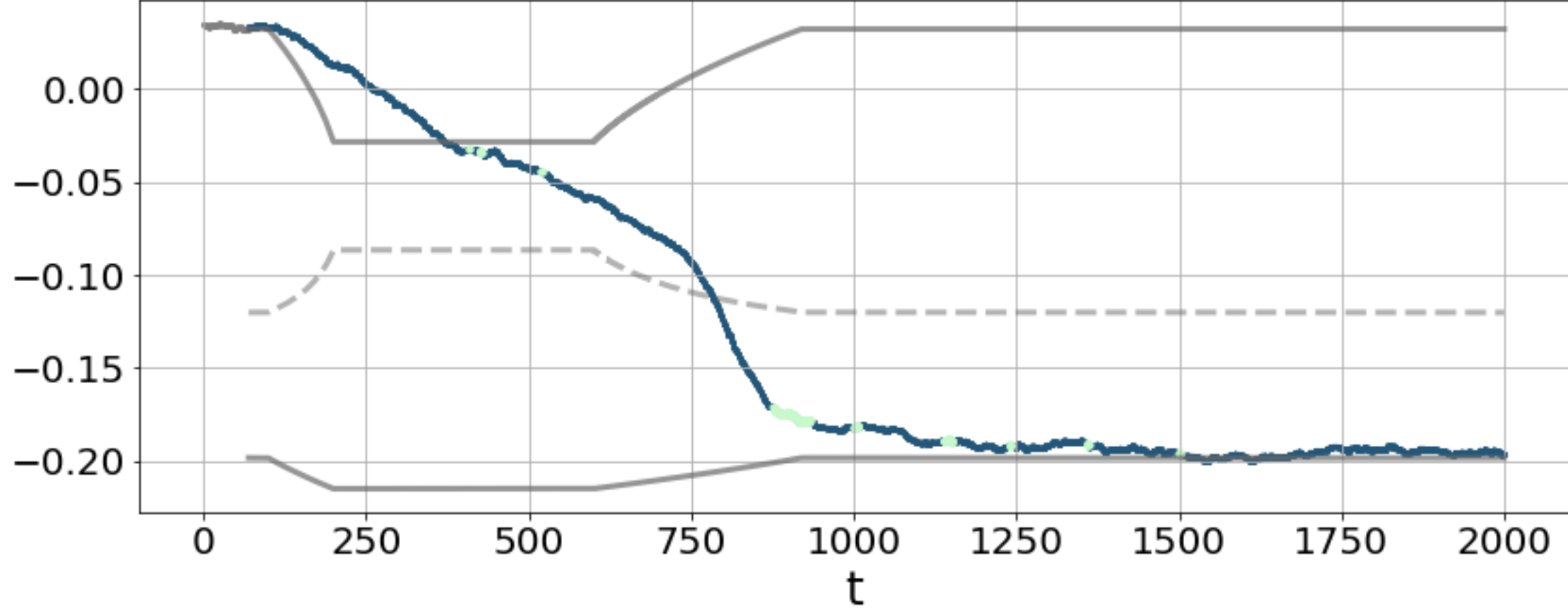
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



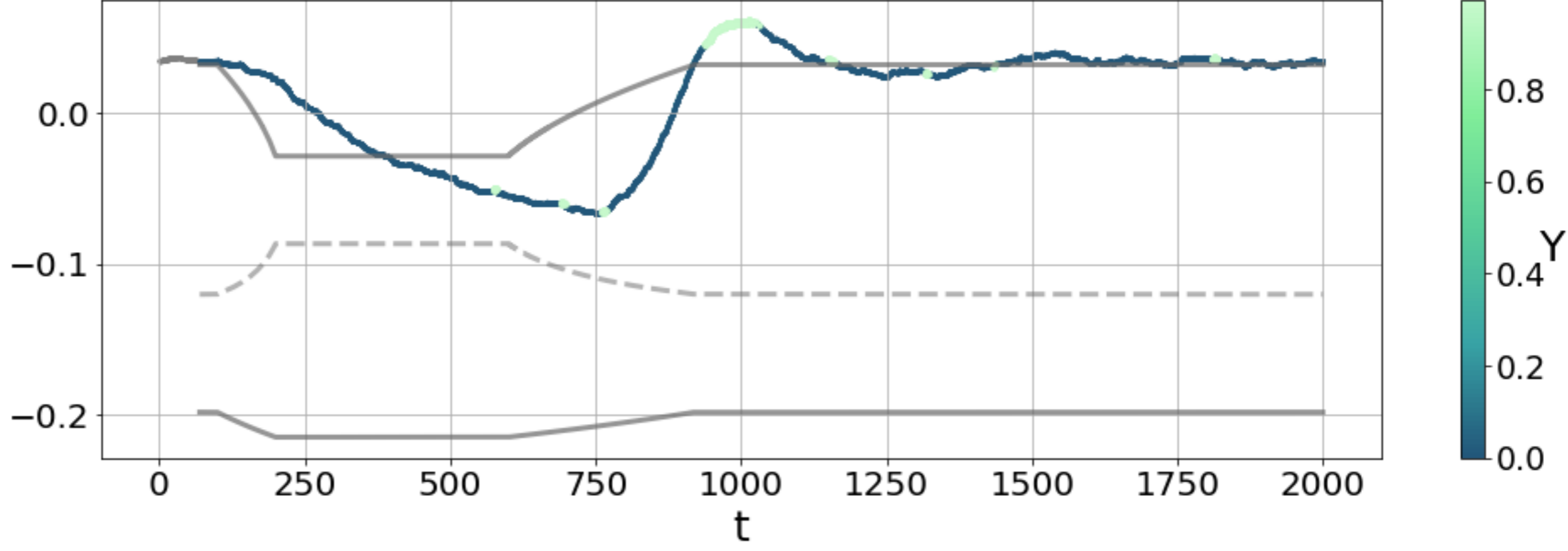
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



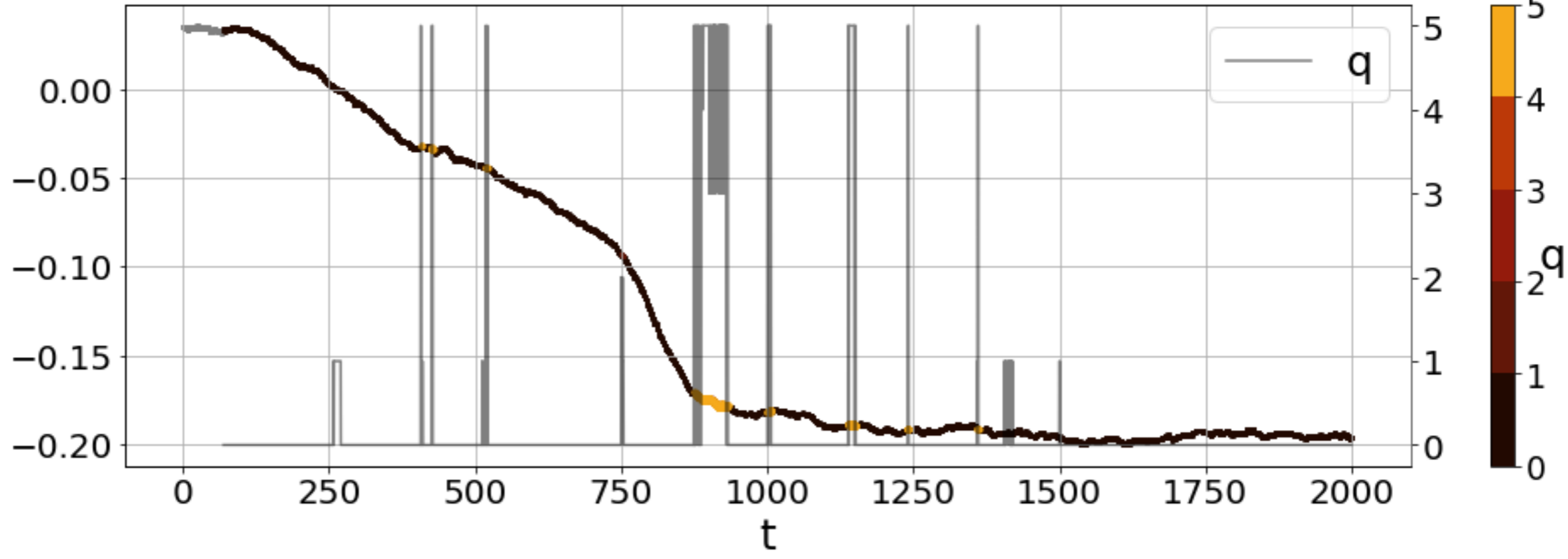
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

 S_N 

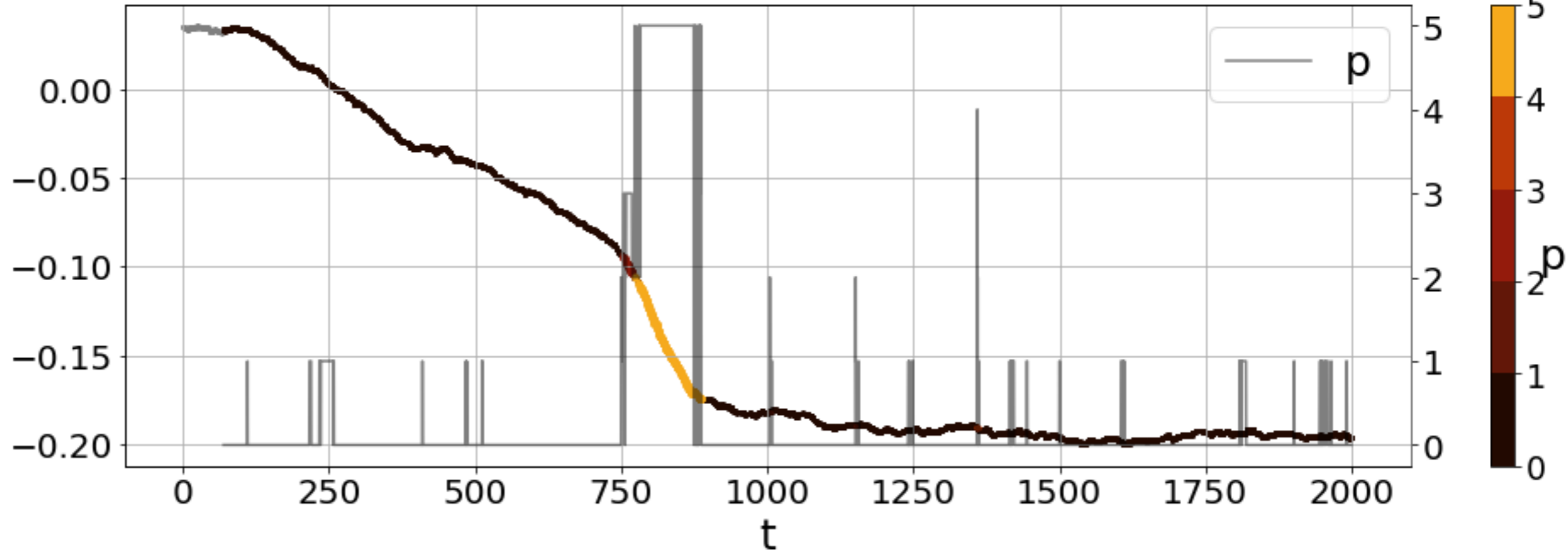
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



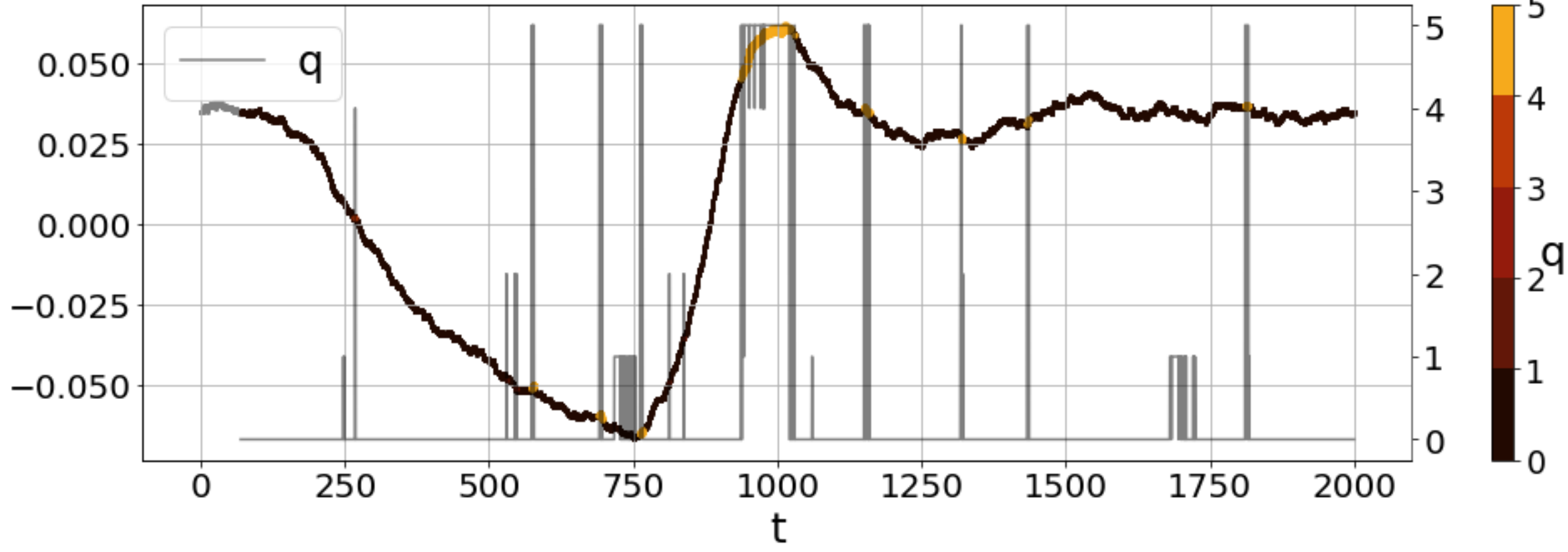
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



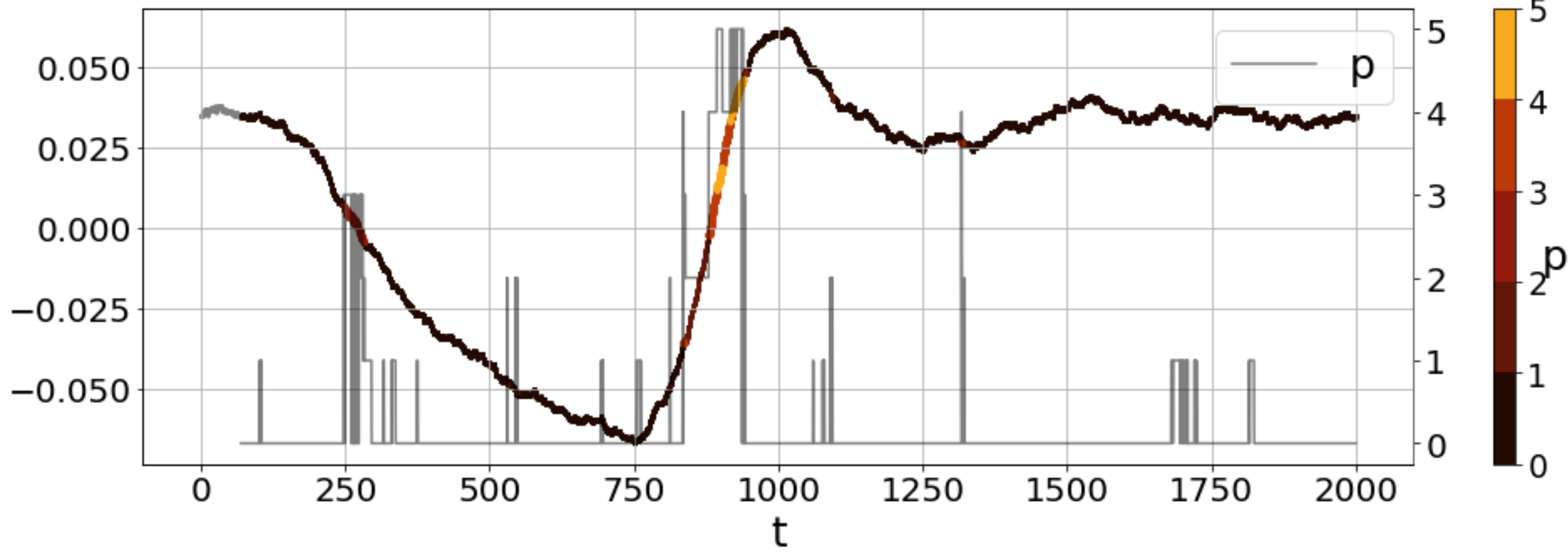
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

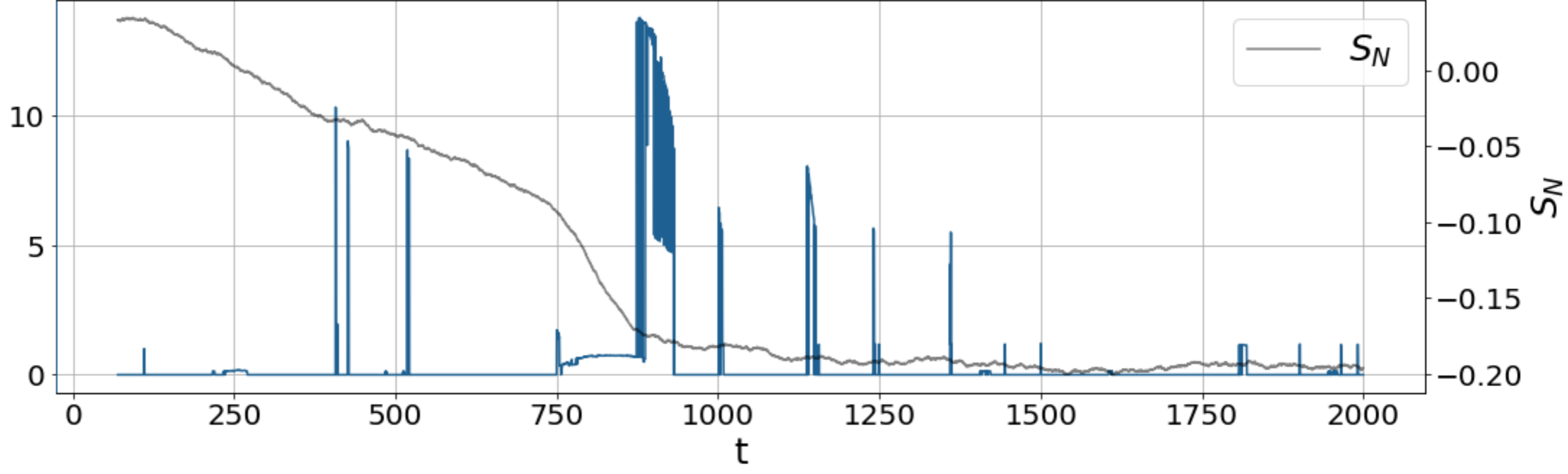


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

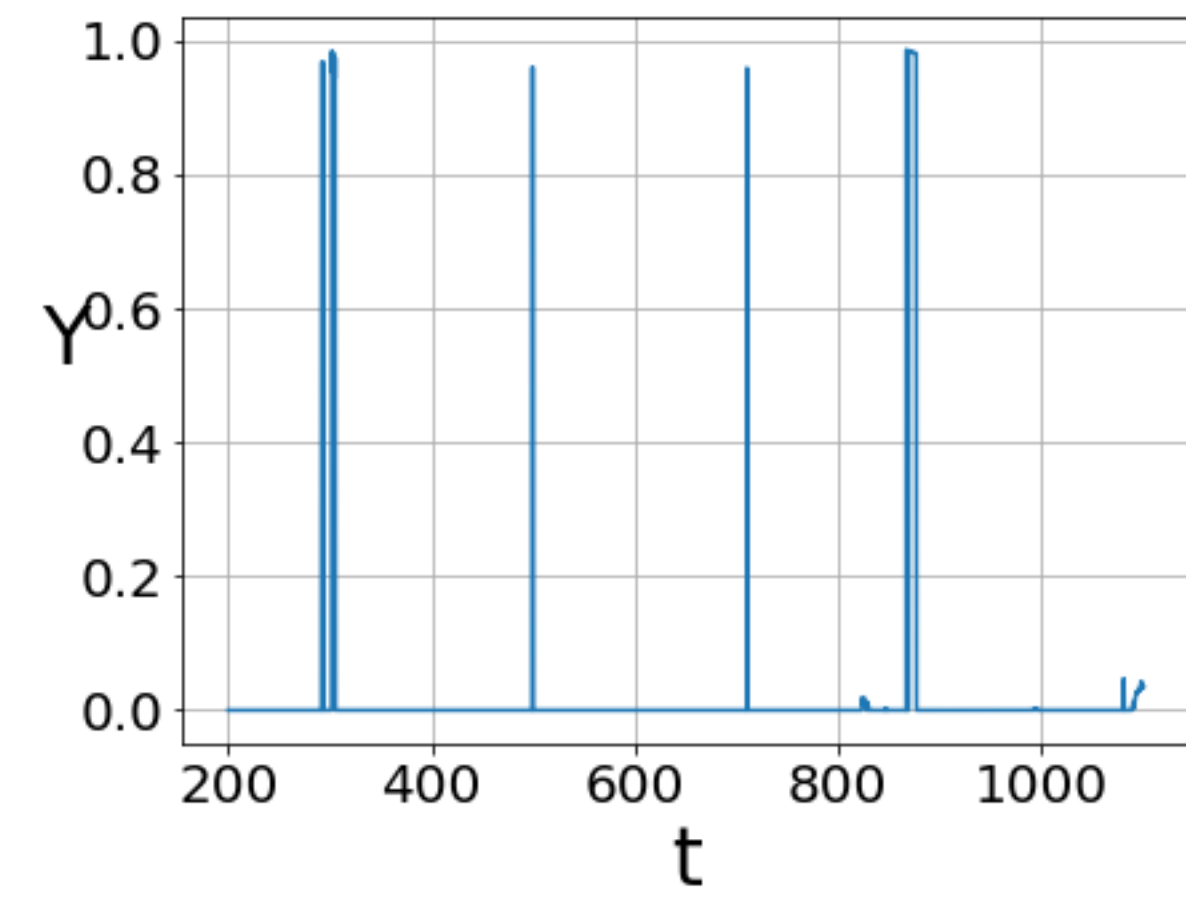
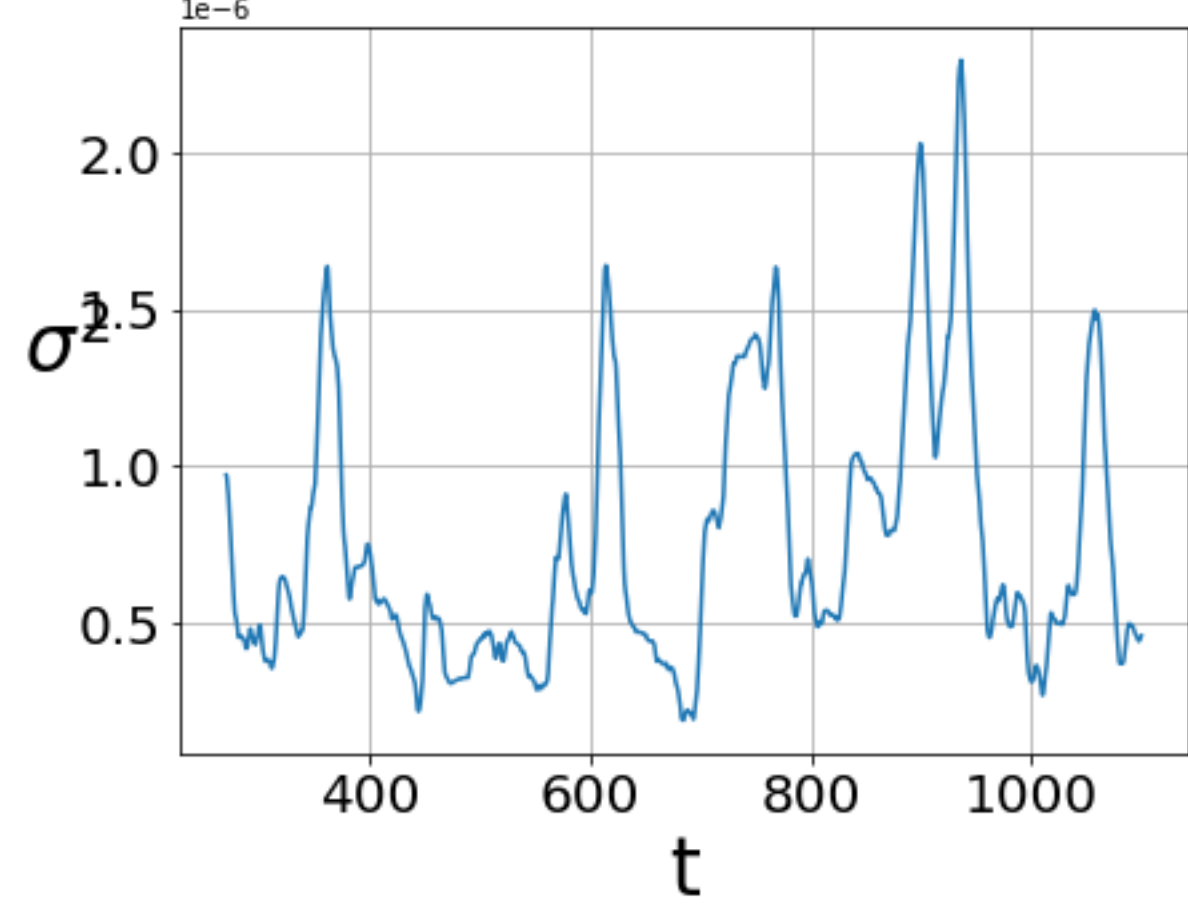
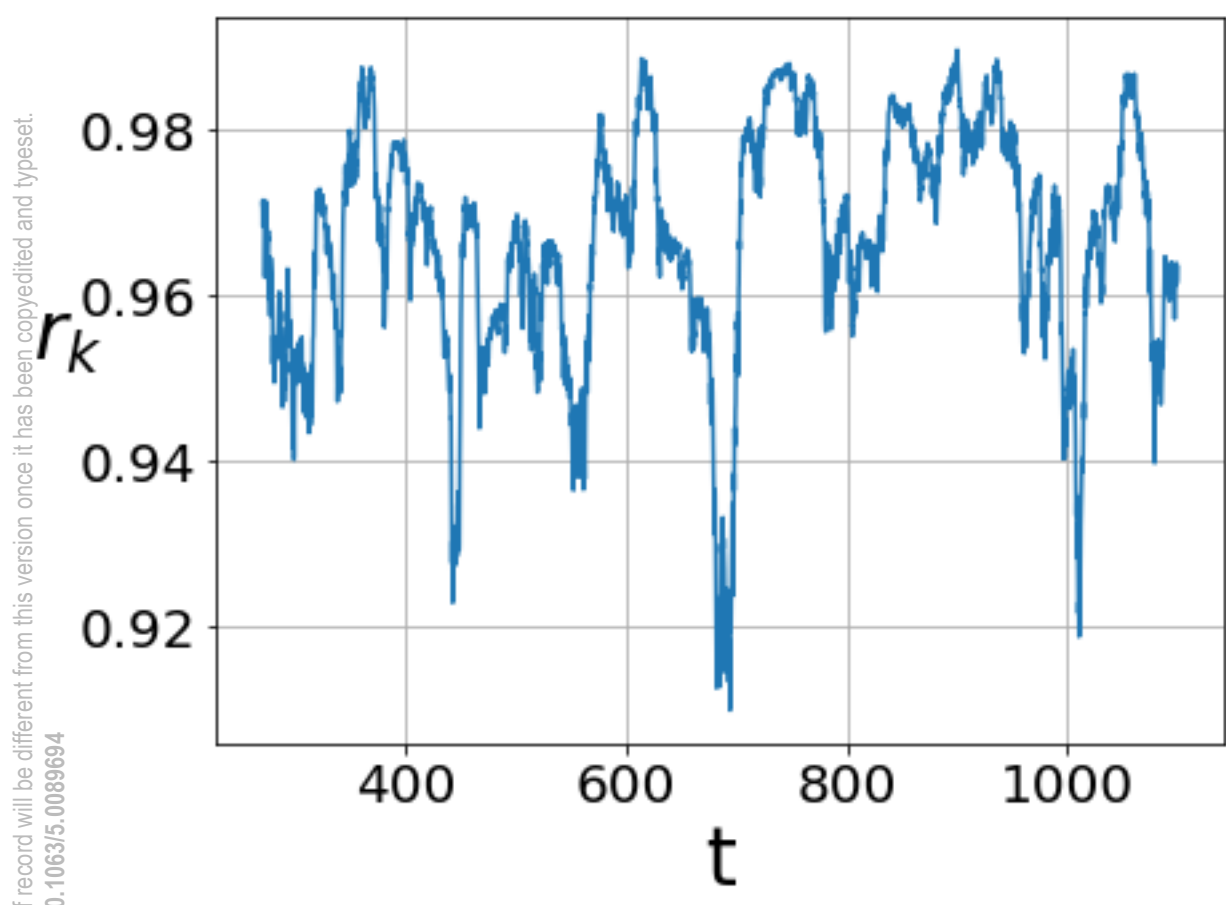


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/1.5008964

\mathcal{R}

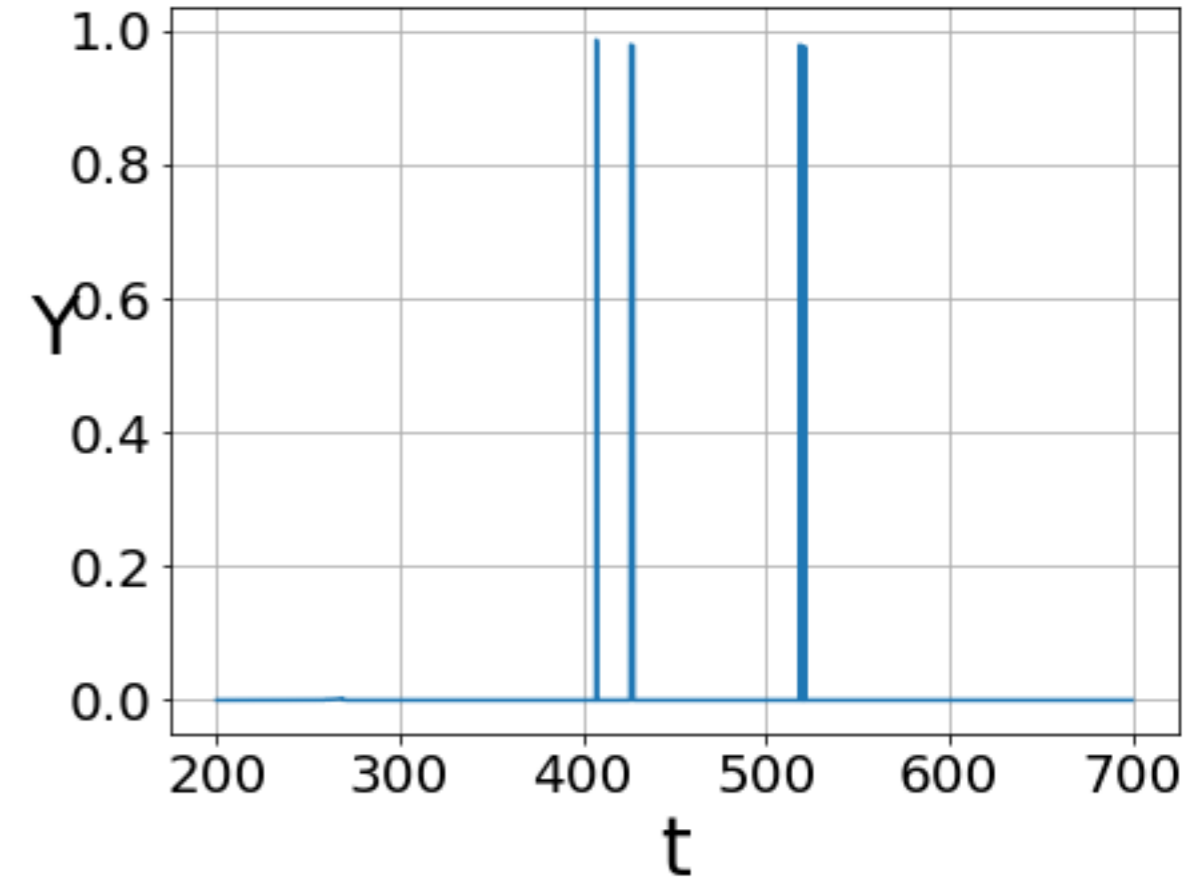
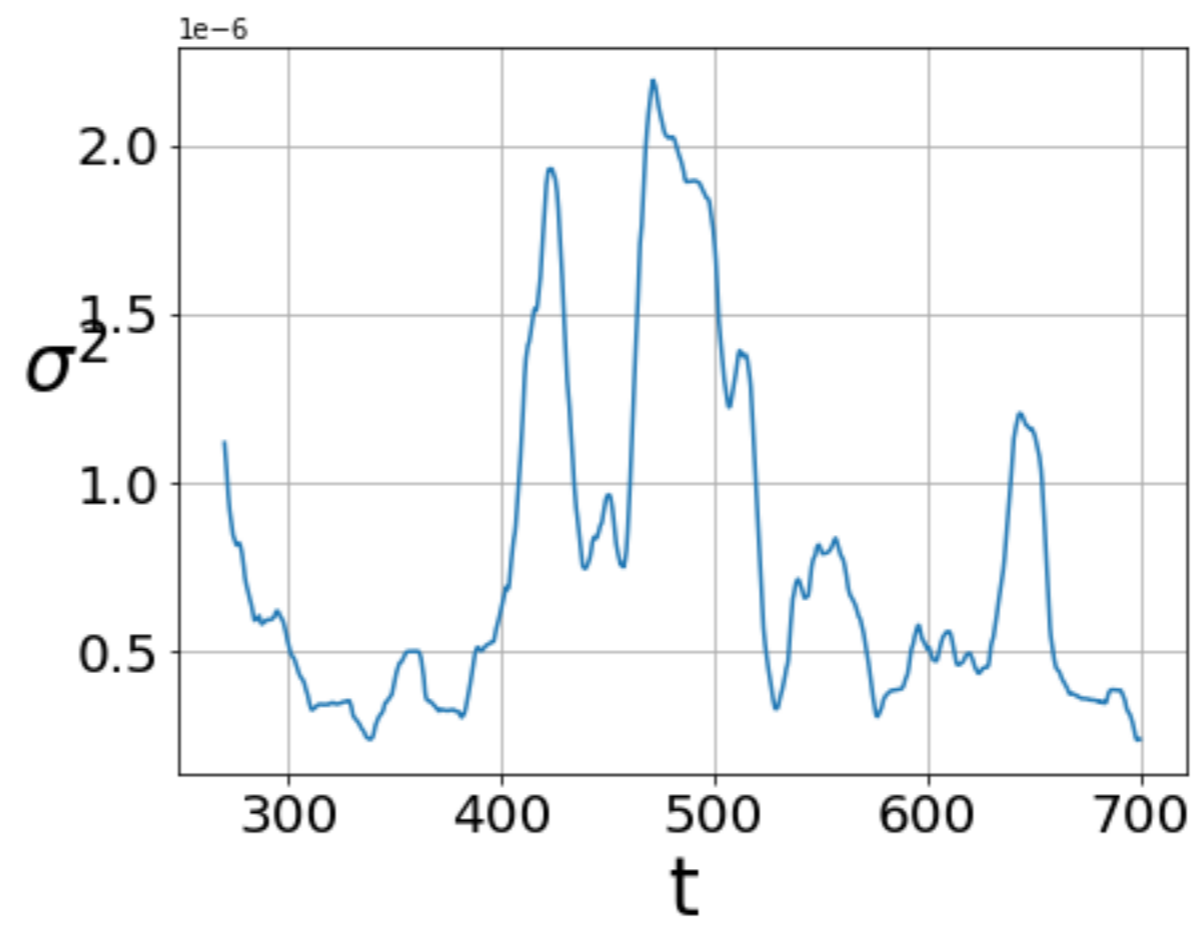
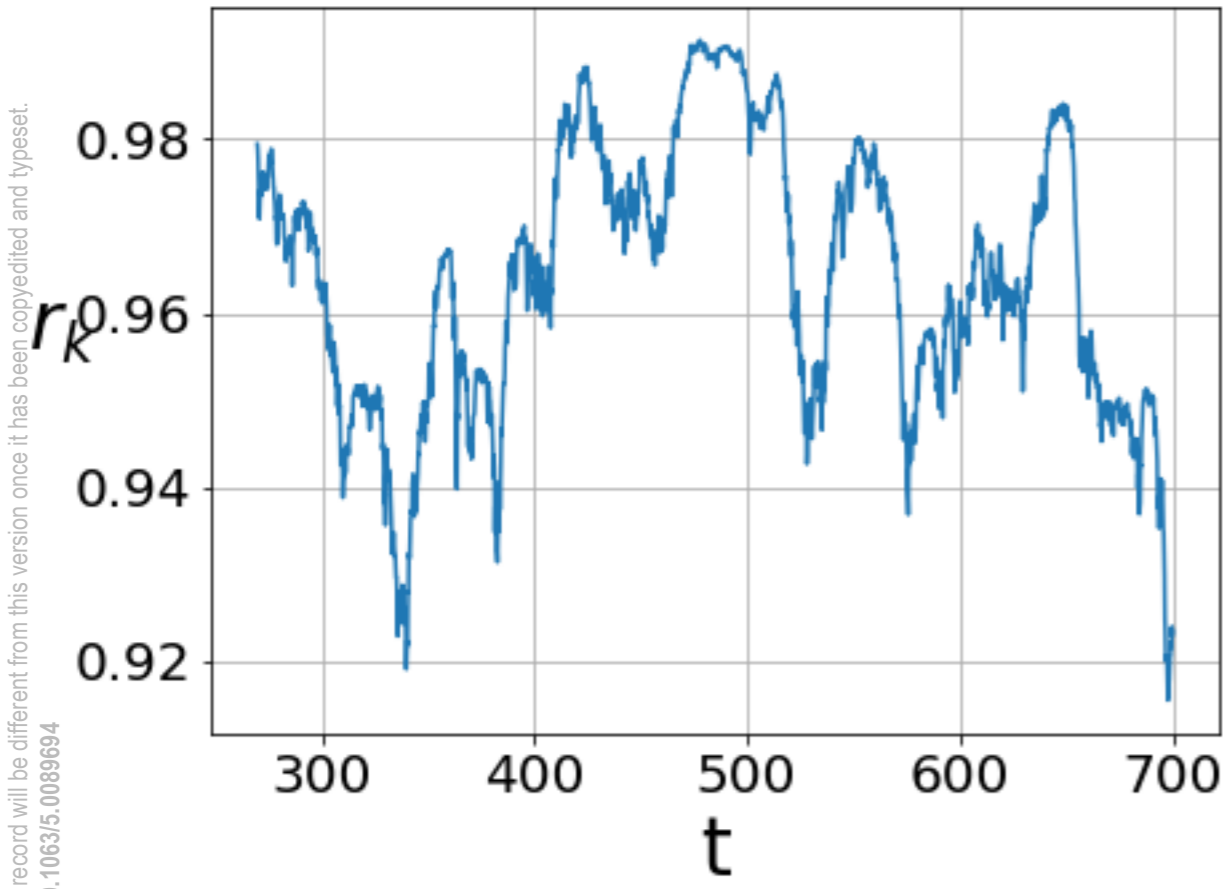


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

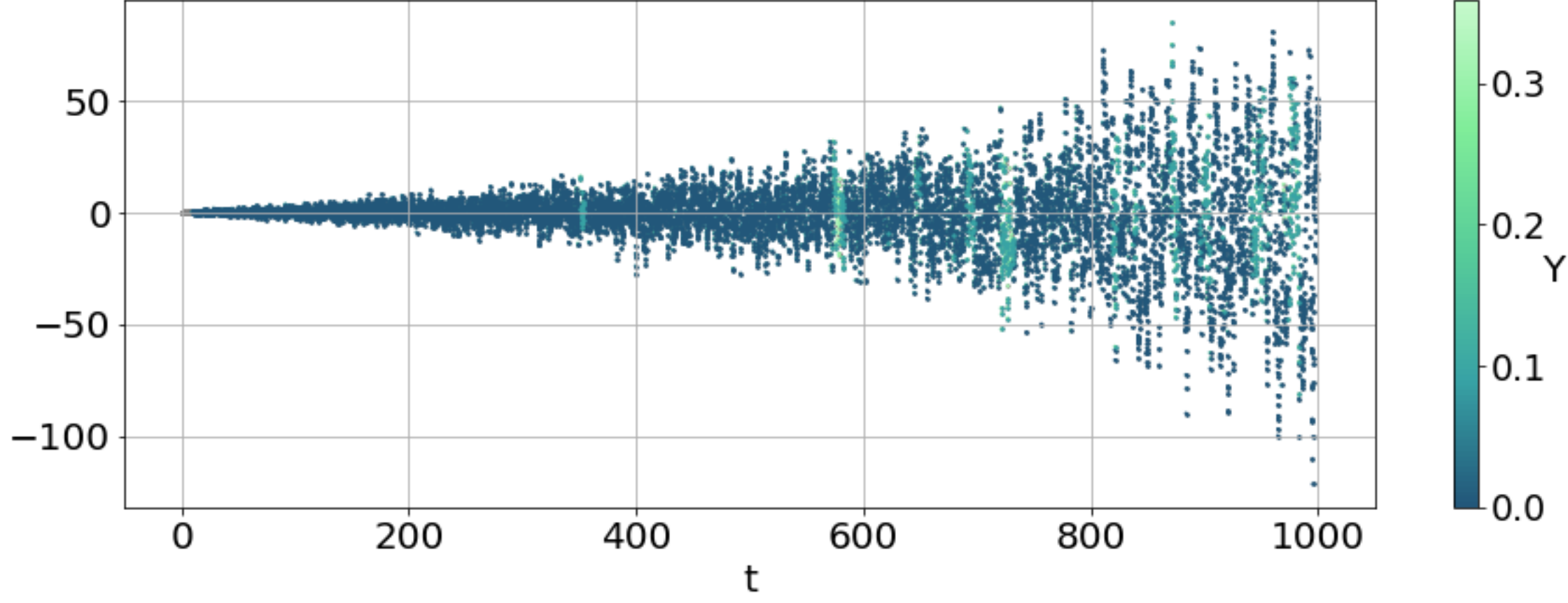


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

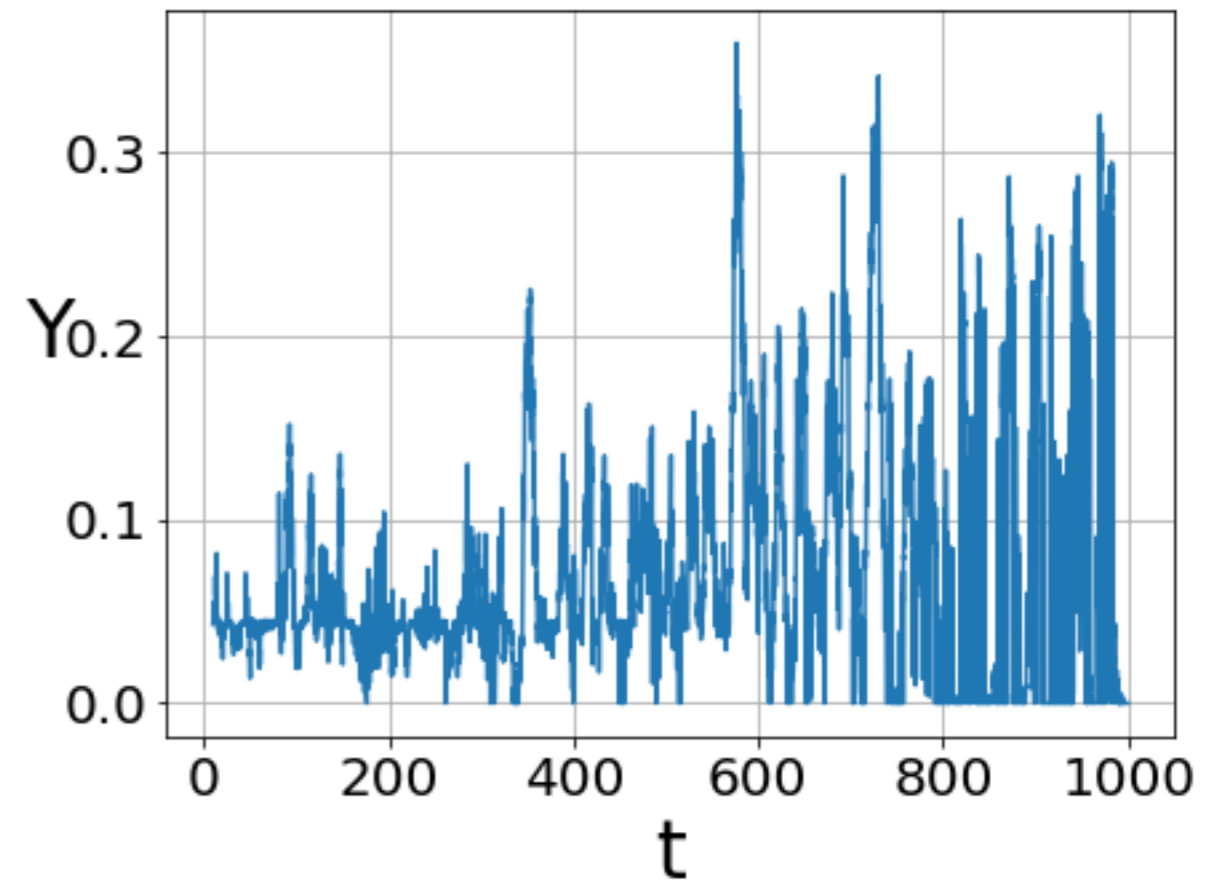
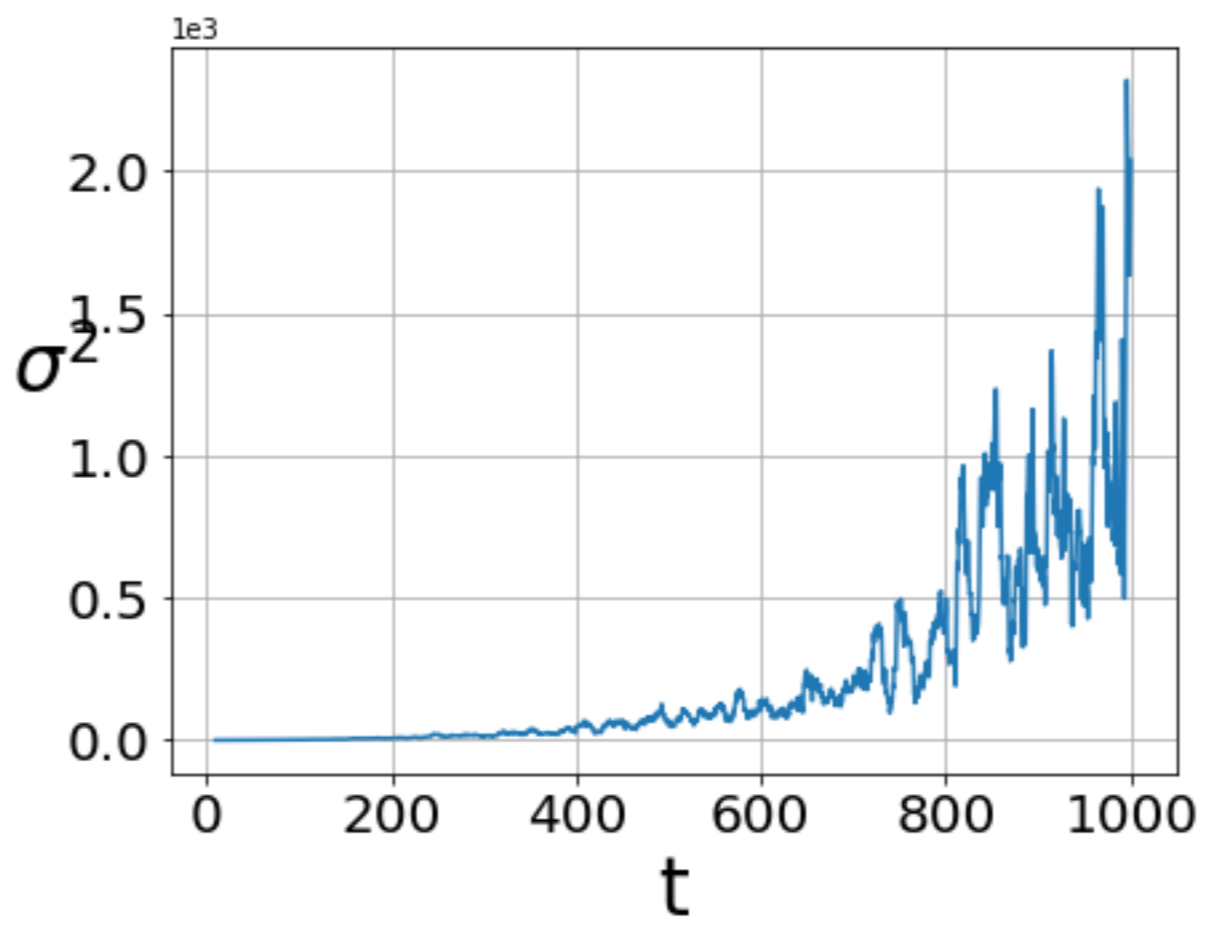
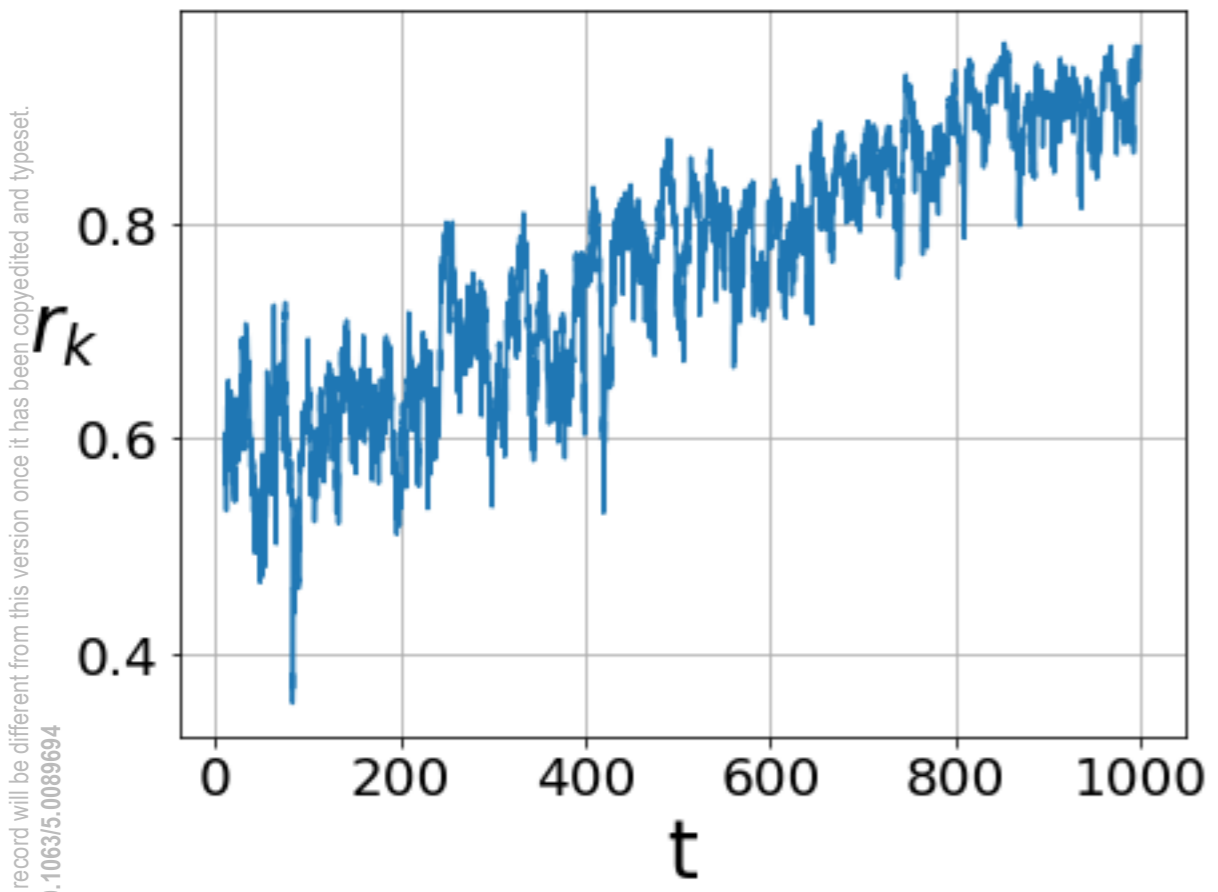


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

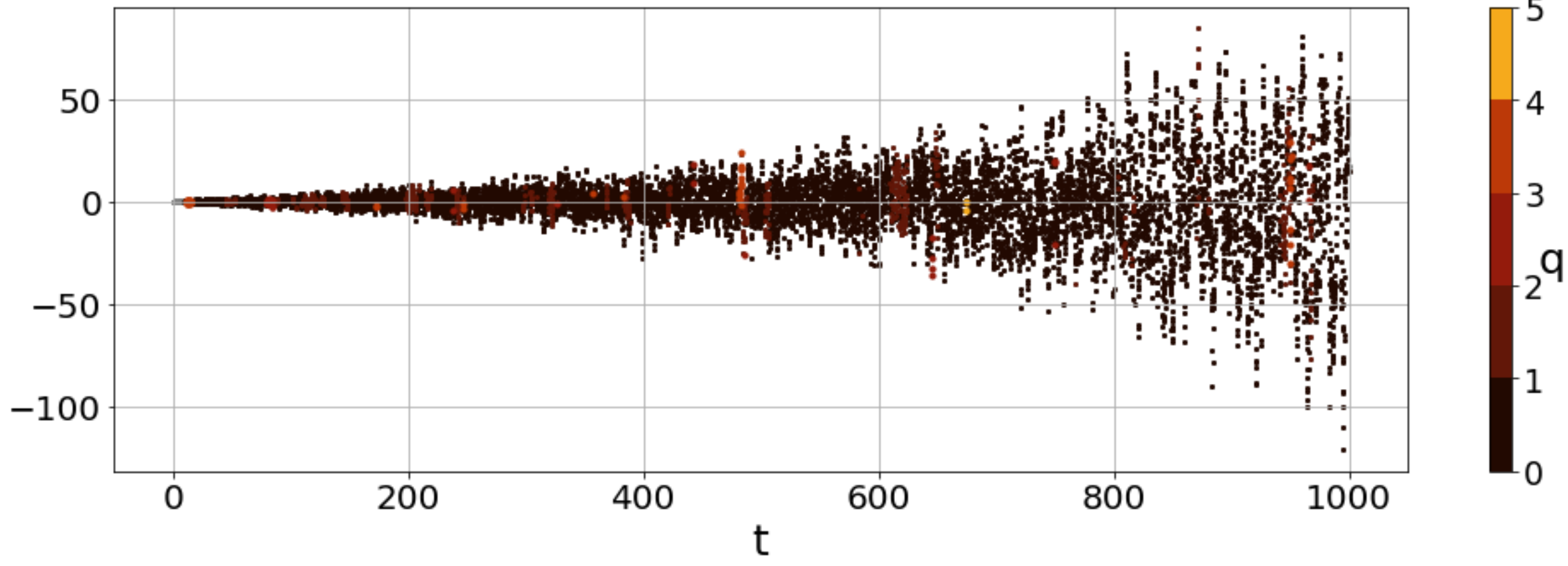


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

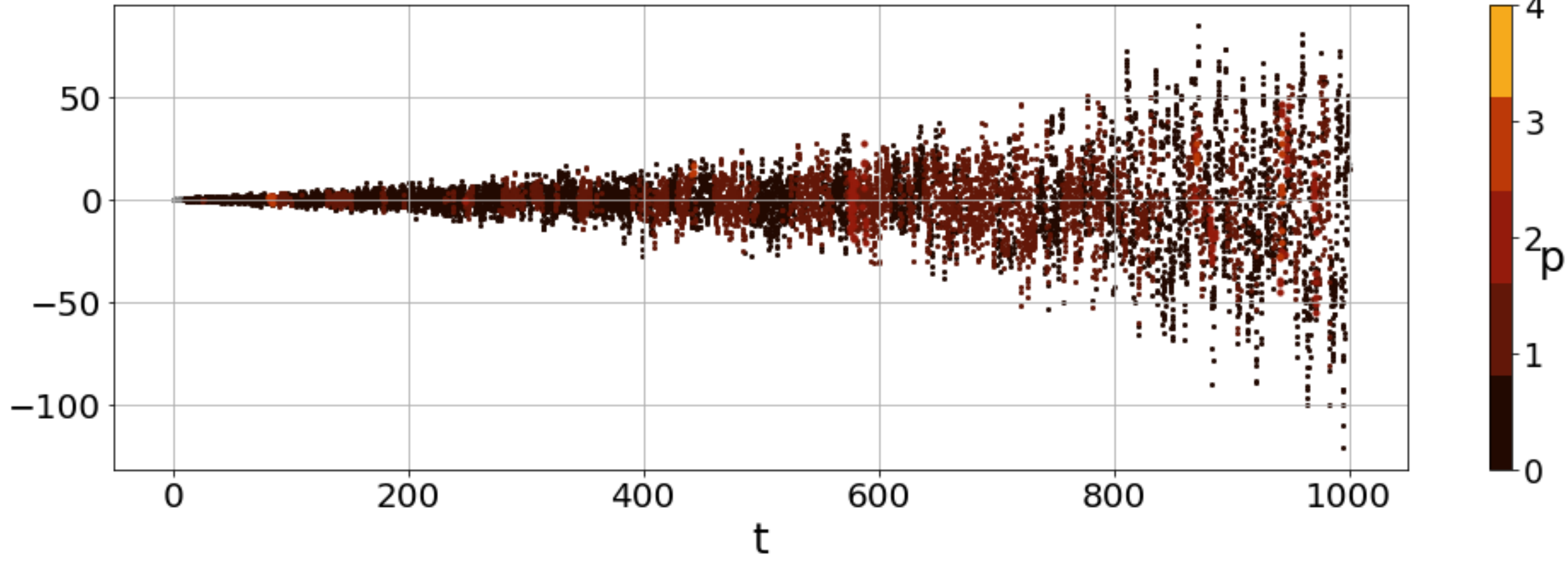
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

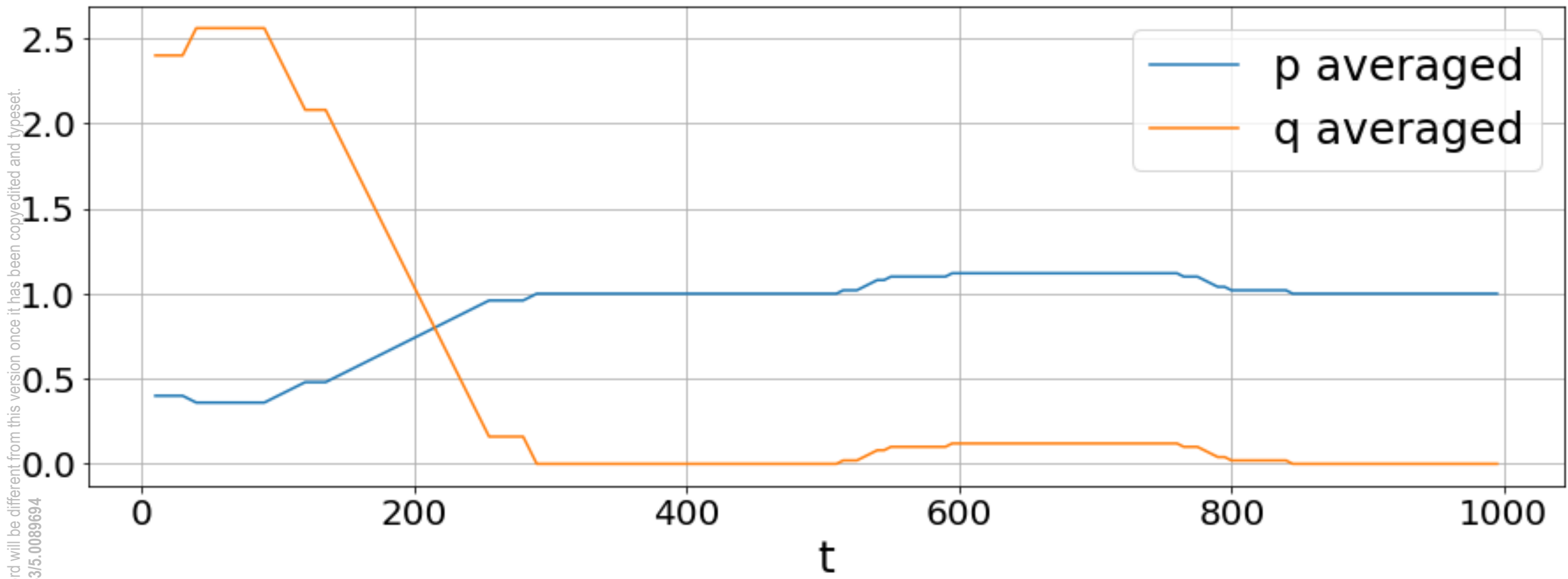


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



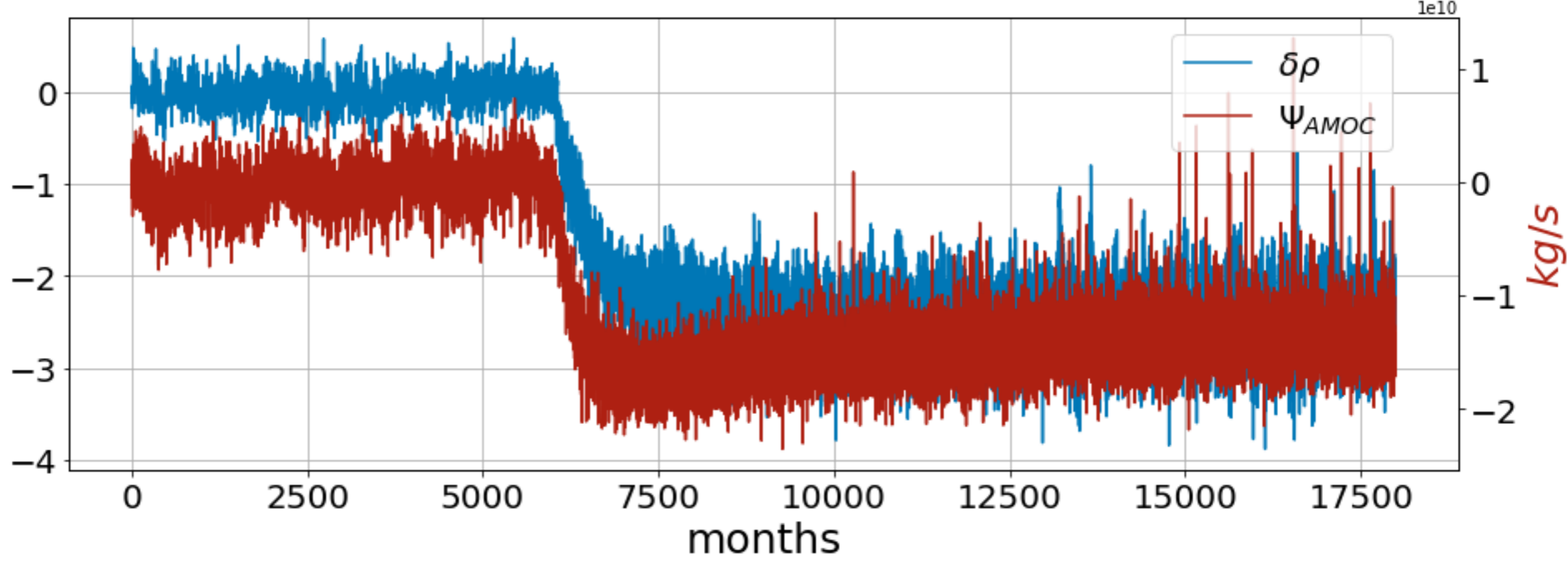
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694





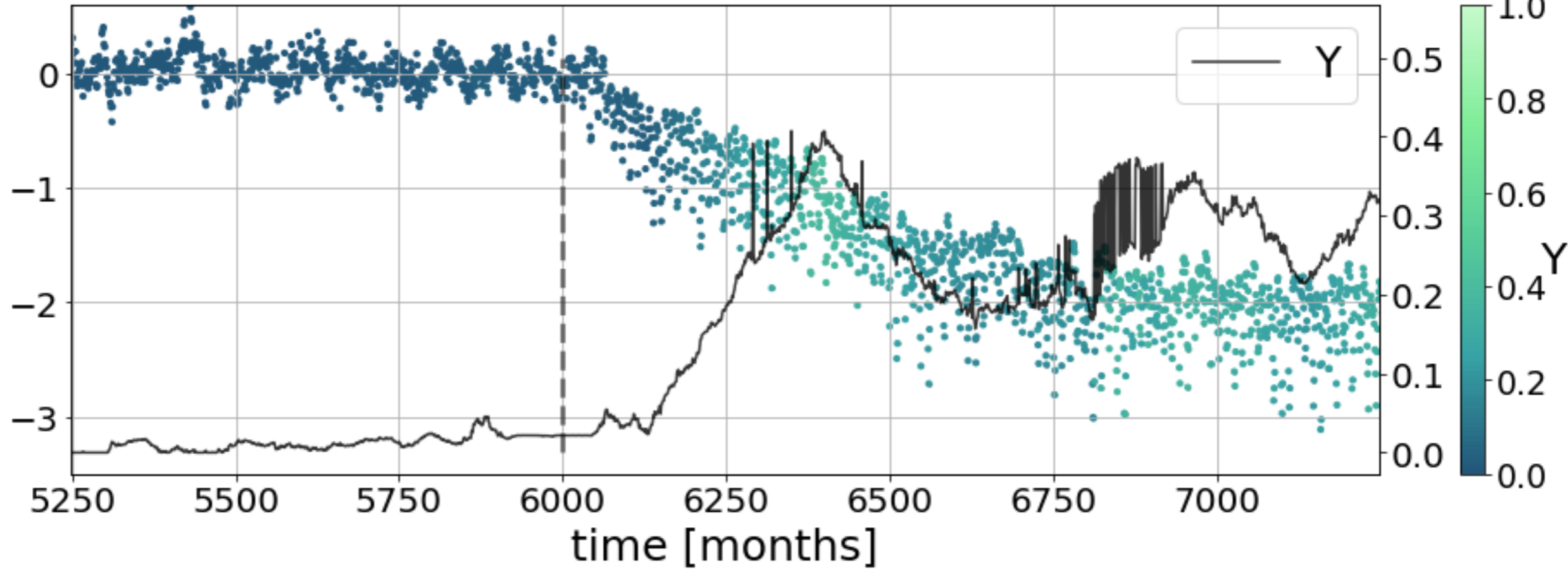
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

kg/m^3



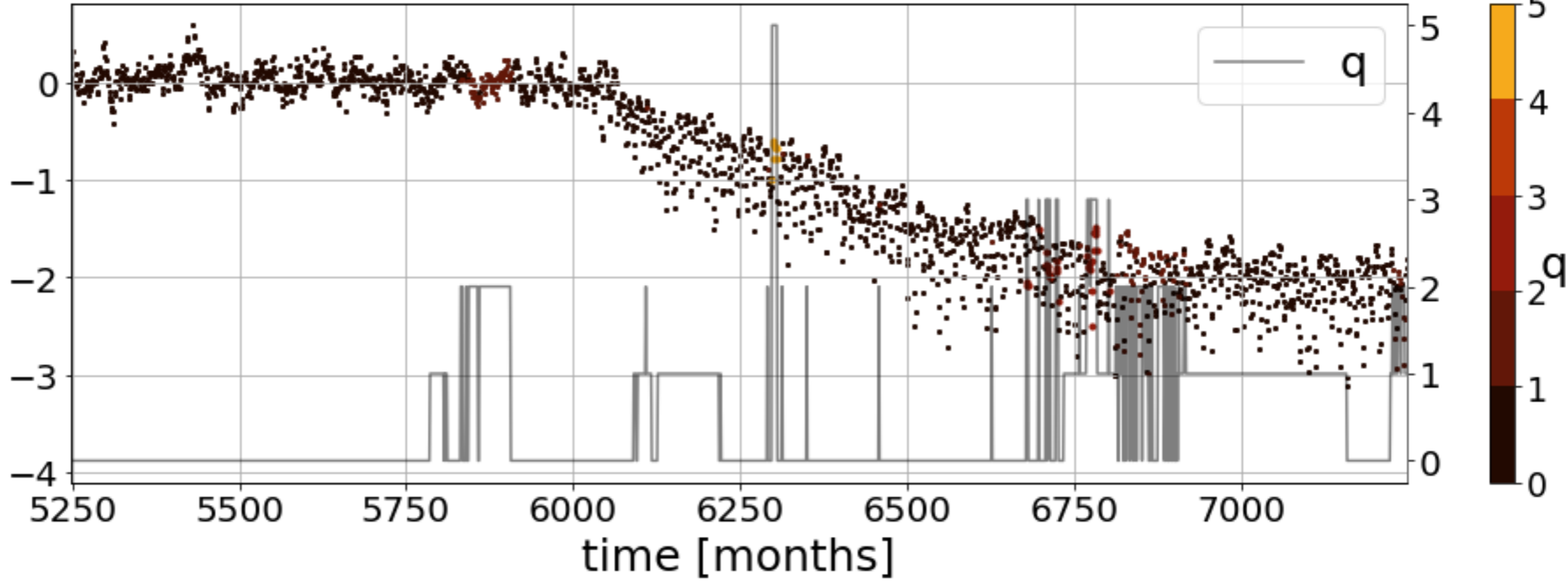
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

$\delta\rho$ [kg/m³]

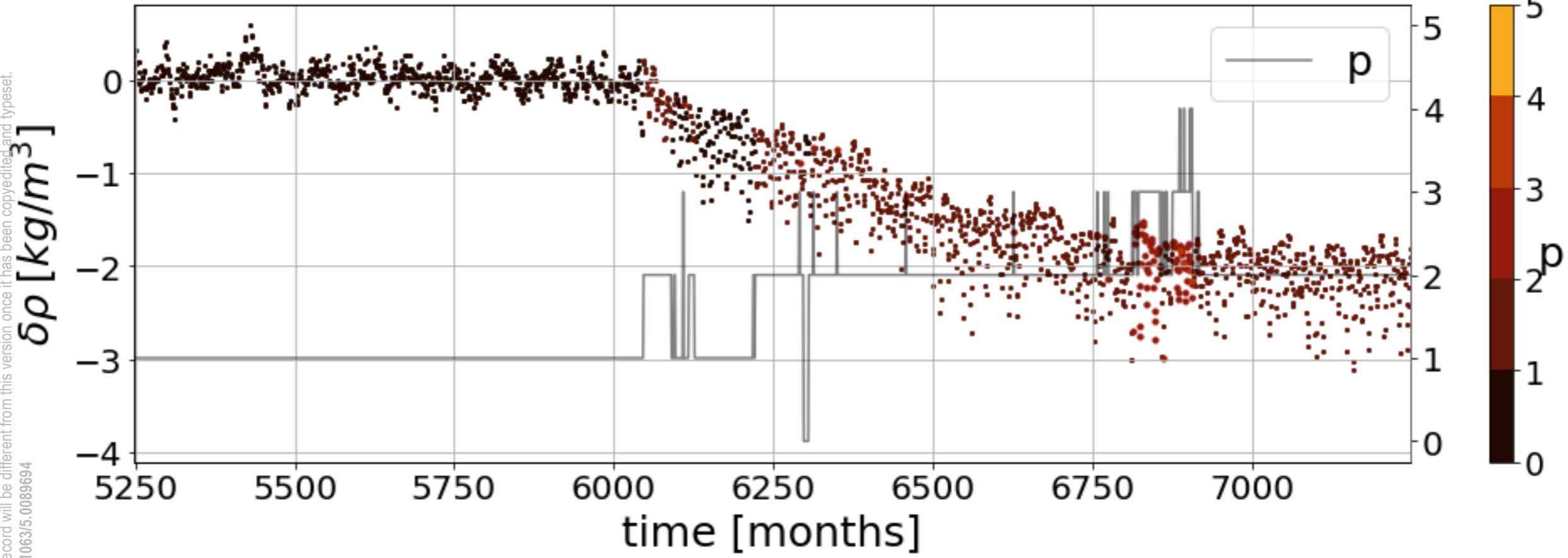


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

$\delta\rho$ [kg/m³]

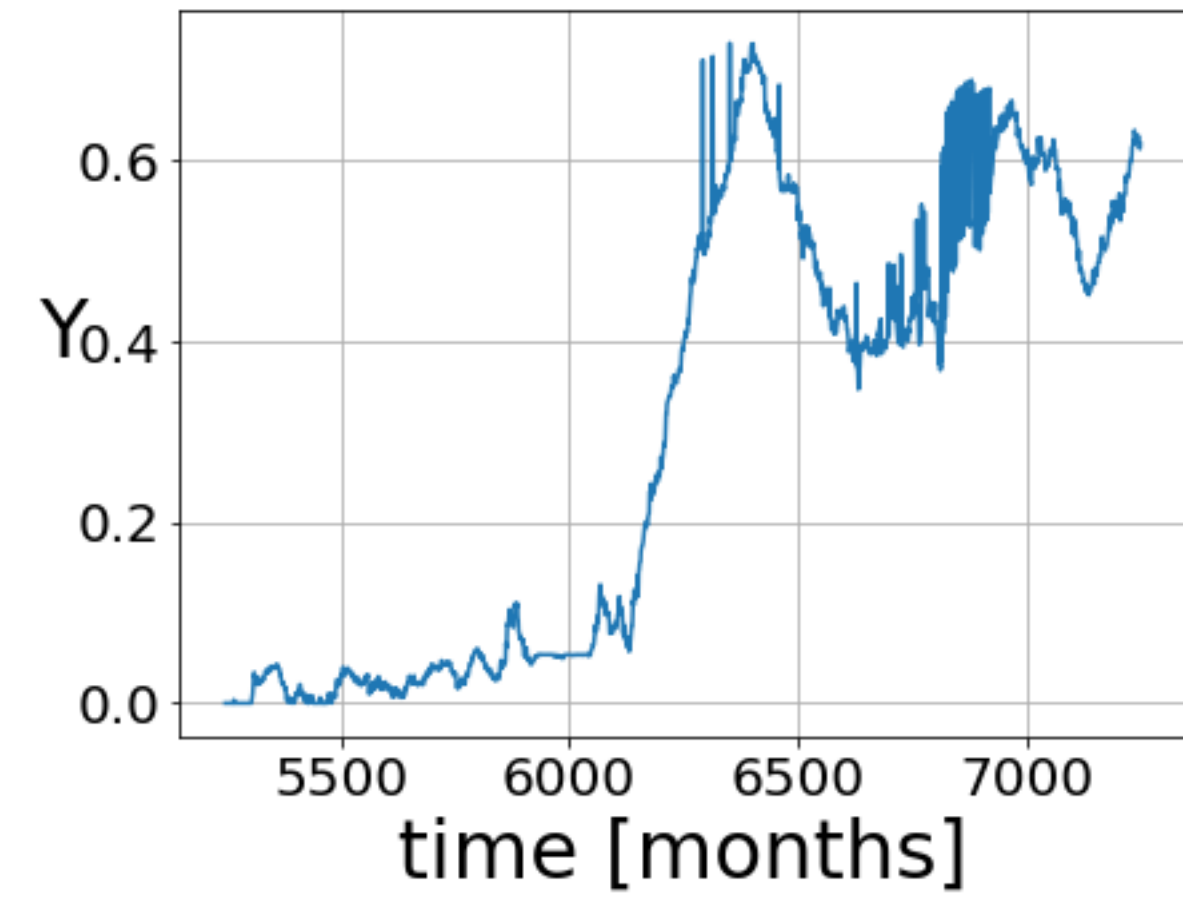
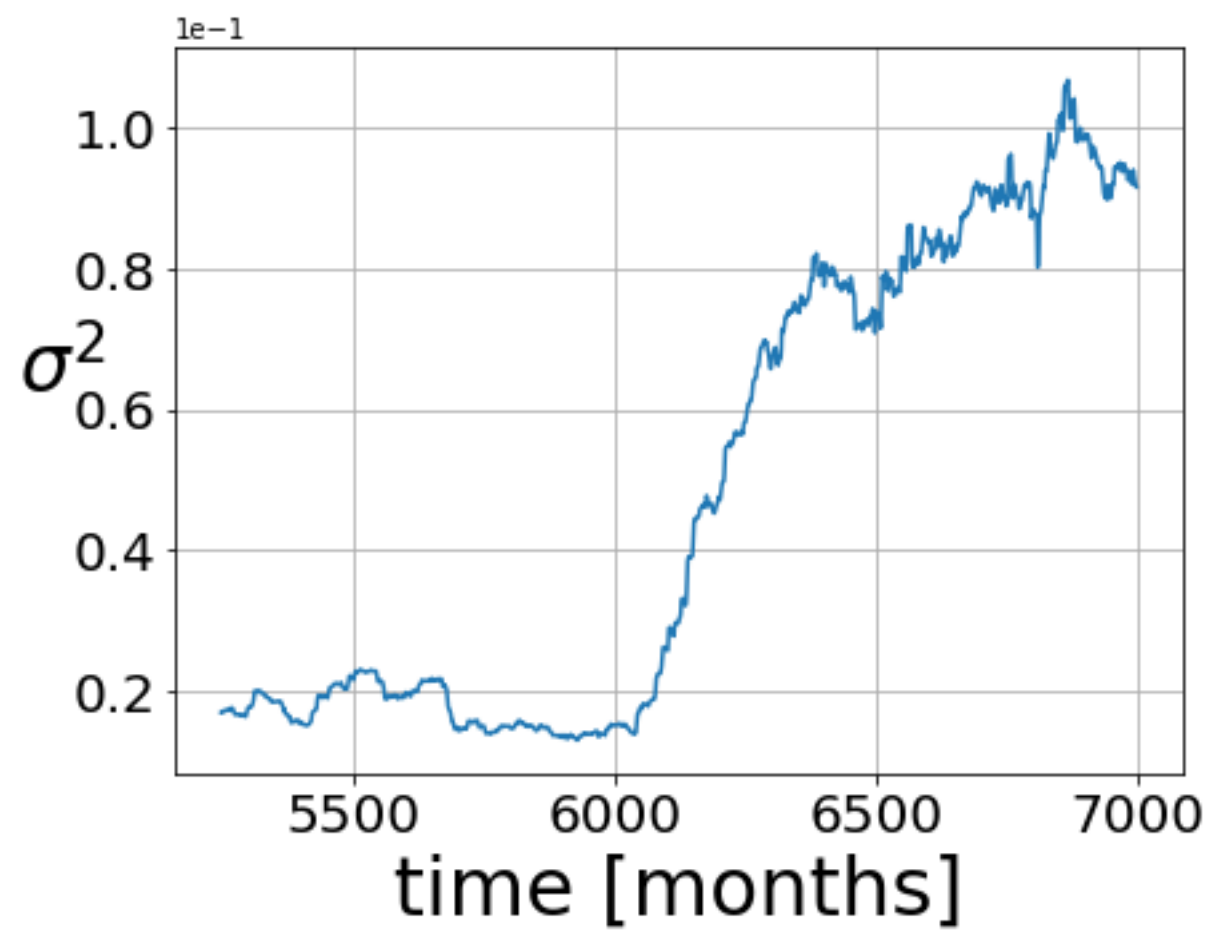
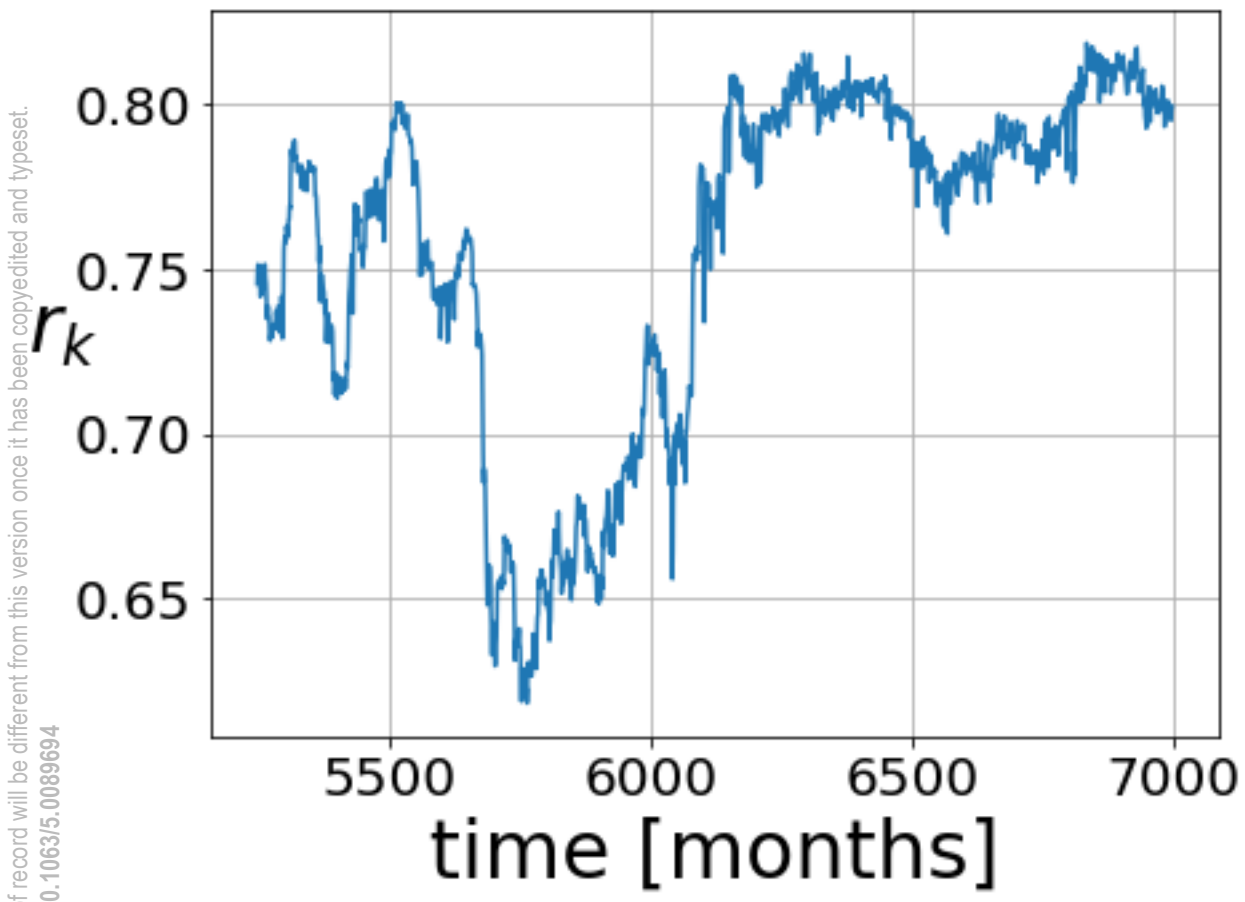


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



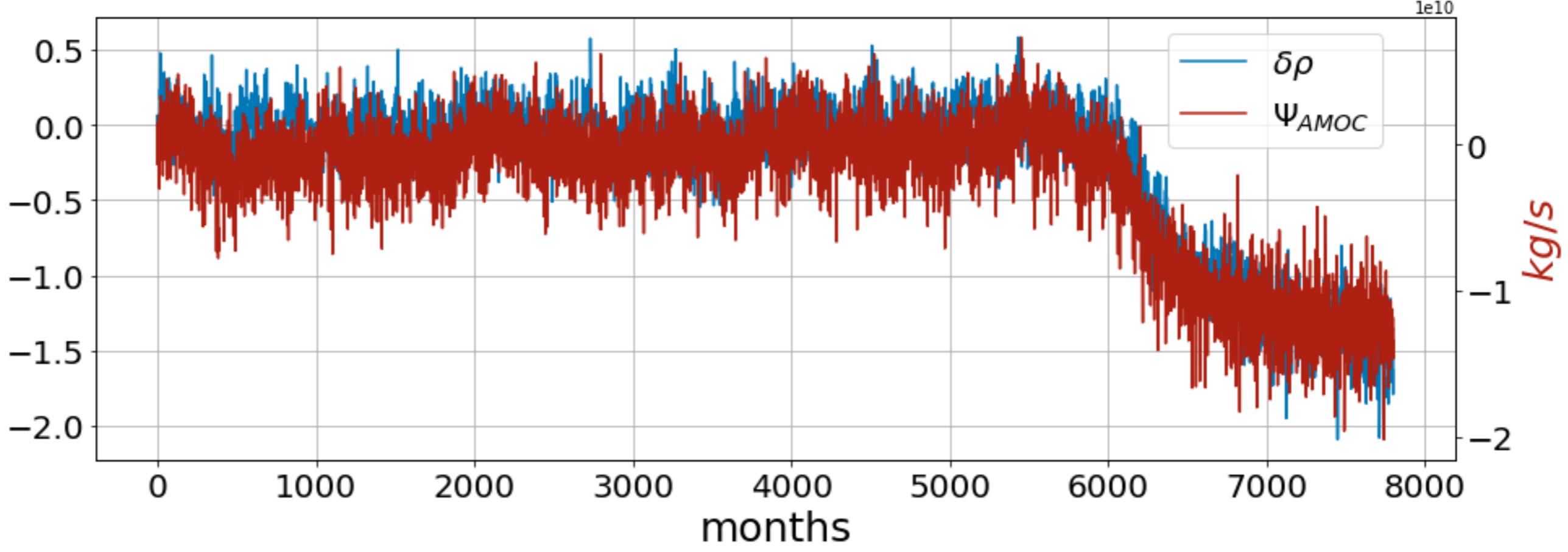
This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

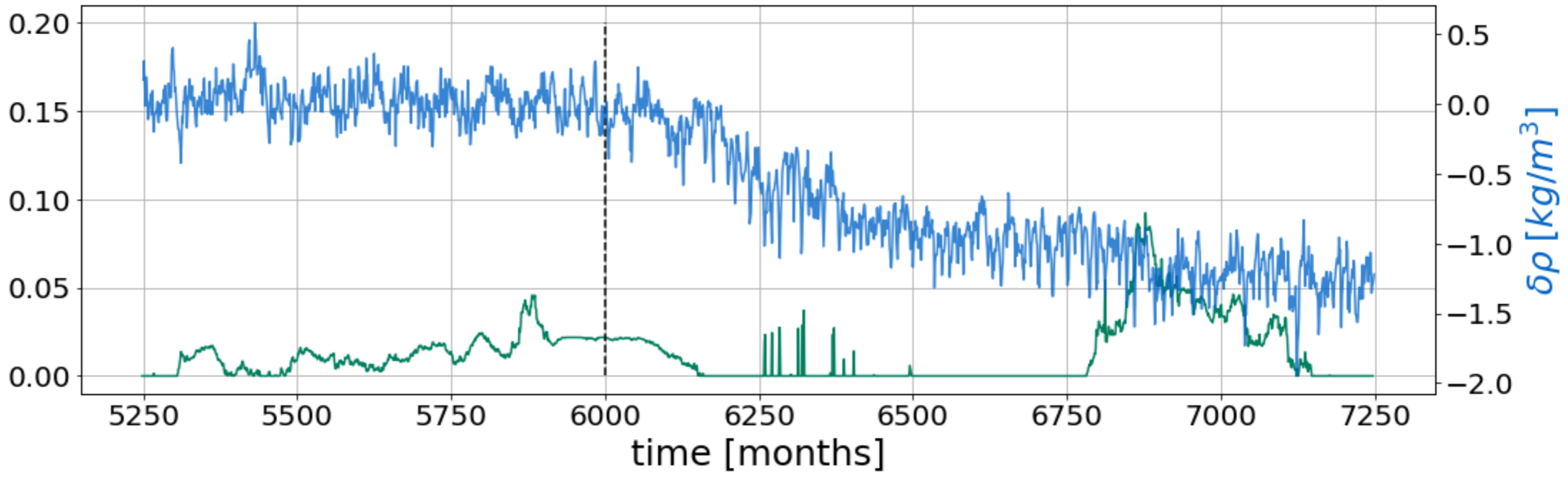


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

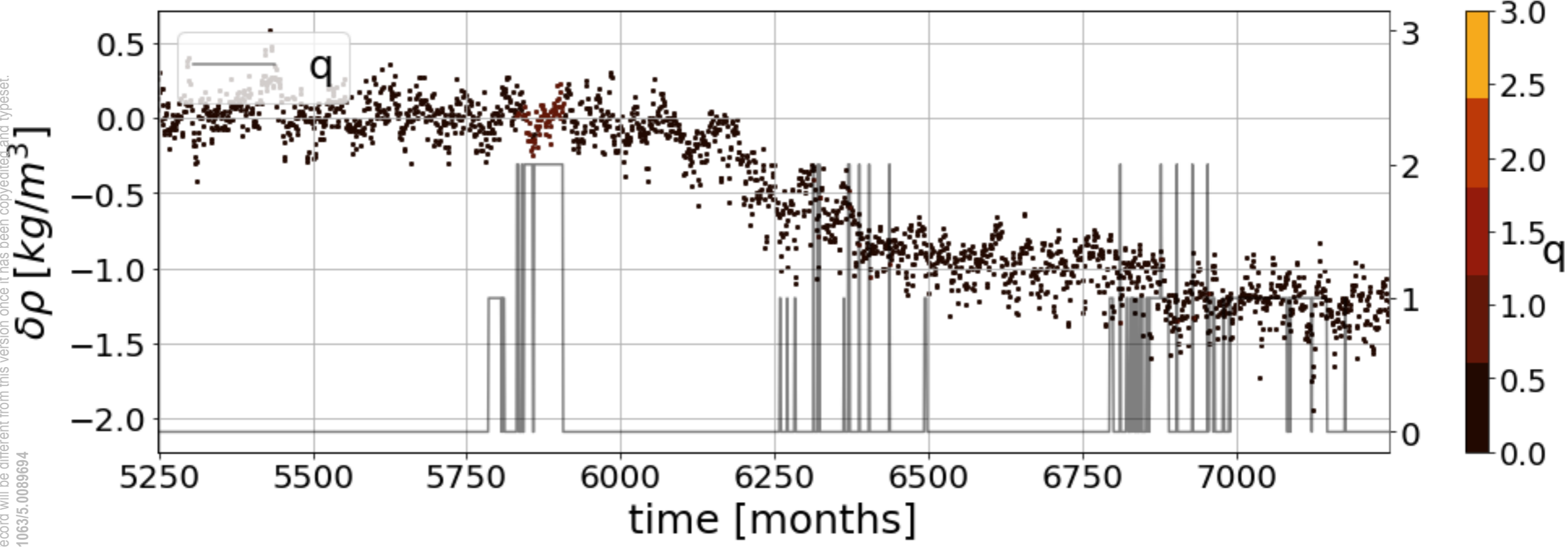
kg/m^3



This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/1.50089694

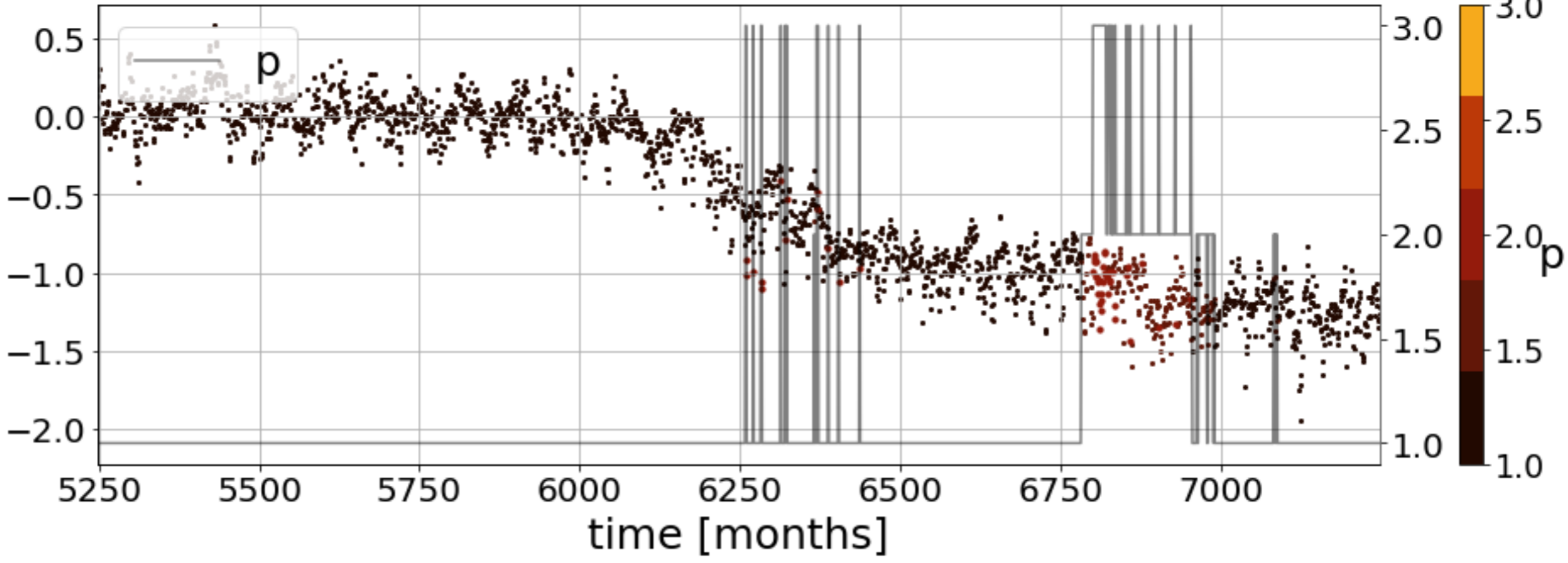


This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694



This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

$\delta\rho$ [kg/m³]



This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0089694

