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Why is Mn charge ordered and AFM coupled in YBaMn₂O₅?



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ABSTRACT

Synthesis conditions to obtain single-phase YBaMn₂O₅ are reported and its crystal structure refined from neutron powder diffraction data collected between 25 and 600 °C in the paramagnetic state. The structure parameters are used to reason why the phase maintains ordered ionic charges of Mn instead of mixing them at an elevated temperature, why this Mn^{2+} and Mn^{3+} ordering is of the checkerboard type in all three dimensions, and why these ions would couple antiferromagnetic upon cooling.

1. Introduction

The three solid phases of the A-site ordered double perovskite along the YBaMn₂O_{5-5.5-6} line convert one to another via reversible oxygen absorption at moderate temperatures, and this wide composition span opens for applications as oxygen-scavenger or -storage materials, solidoxide fuel-cell cathodes, or as catalytic oxidants [1]. The YBaMn^{II}Mn^{III}O₅ end phase [2] of Mn²⁺ and Mn³⁺ configurations d^5 and d^4 is one-step from d^4 and d^3 of LaCaMn^{III}Mn^{IV}O₆ [3] and other colossal-magnetoresistance perovskite manganites [4,5]. It is also one-step from d^6 and d^5 that leads to valence mixing in YBaFe^{II}Fe^{III}O₅ [6] above the Verwey-type [7] transition temperature.

There are notable differences between the Fe and Mn variants. While the d_{xx} orbital order in YBaFe₂O₅ stabilizes charge ordering that defies the electrostatic minimum for *point* charges by alternating chains of Fe²⁺ coordinations and chains of Fe³⁺ coordinations [6], the Mn²⁺ and Mn³⁺ in YBaMn₂O₅ comply with it and alternate in all three directions [8]. While iron atoms in YBaFe^{II}Fe^{III}O₅ become equivalent above 335 K [6], YBaMn^{III}Mn^{III}O₅ remains charge ordered even at 873 K, as observed in the neutron diffraction study [9] of TbBaMn₂O₅ in an equilibrium 1:1 mixture with TbBaMn₂O_{5.5}. Only the magnetic cooperative order is the same, as charge-ordered YBaFe₂O₅ [6], and somewhat surprisingly YBaMn₂O₅ at low temperatures [8], both adopt the G-type spin arrangement of Wollan–Koehler [3]. YBaFe₂O₅ does so with moments along *b*, YBaMn₂O₅ along *c*.

In this study, synthesis conditions to obtain phase-pure YBaMn^{II}Mn^{III}O₅ of non-magnetic and redox-stable yttrium were established, and the crystal structure investigated by high-intensity neutron powder diffraction up to 873 K for signs of approaching valence mixing. The refined structure parameters are used to calculate molecular-orbital energies, from which it is argued how the magnetic ordering in YBaMn₂O₅ complies with the Goodenough's model [10] for semicovalent exchange and what it might take to mix the Mn valences.

2. Experimental

Syntheses and characterizations. Polycrystalline YBaMn₂O_{5+w} was sintered from precursors obtained by liquid mixing in melted citric acid monohydrate [11]. Annealed Y2O3 (99.99%, Fluka) was dry-mixed with the citric acid (BDH AnalaR reagent, <0.01% sulfated ash, <0.002% Ca), in amount of 2 mol per mole of cation charge of each starting component, and dissolved upon melting assisted by water added just to cover the bottom of the beaker. Manganese lumps (99.9%, Ventron) were cleaned of oxide in a 1:1 diluted hydrochloric acid, washed, and dried. The obtained 8.3786 g of Mn was dissolved in 50 mL water upon gradual additions of in total 80 mL of 65% HNO3 and then added dropwise into the above melt to a complete release of nitrous gasses. After cooling below 100 °C, redistilled water was added on top of the viscous melt, and barium carbonate (Fluka, <0.2% Sr) was dissolved upon warming. The resulting clear melt of brown color was decomposed in a drying oven at 180 °C into an organic-based solid xerogel, which spontaneously incinerated overnight (unexpected, see typical synthesis details in Ref. [12]). The product was milled, pressed into cylindrical pellets of 15 mm in diameter and calcined for 24 h at 860 °C in a flowing atmosphere of 10.01(2) mole % H₂ in Ar (AGA, certified analysis) wetted in a saturated solution of KBr at 18 °C. The water-formation equilibrium then fixes the oxygen partial pressure in the reaction atmosphere to $\log(p_{02}/\text{bar}) \approx$ -18.5. Because small amounts of Y2O3 were detected in the calcination product, a second calcination was performed on the rehomogenized sample under flowing atmosphere of the above-used 10% H₂/Ar gas mixed with Ar (AGA, 6 N purity) to an Ar/H2 molar ratio of 78(3) and wetted at 19 °C to stabilize $\log(p_{02}/\text{bar}) \approx -16.6$. This yielded a pure mixture of the YBaMn₂O₅-type and YBaMn₂O_{5.5}-type phases of brown color. After testing the optimal sintering conditions (Table 1), pellets of 8 mm in diameter were pressed at 150 bar and sintered for 24 h at 1200 °C into a silvery YBaMn₂O_{5.006} under flowing atmosphere of 10% H₂/Ar gas mixed with Ar in two steps and wetted in the saturated solution of KBr at

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Table 1

Sintering conditions that led to single-phase YBaMn₂O_{5+w} samples.

_						
w *		$t/^{\circ}C$ Ar/H ₂ **		$\log(p_{O2}/\text{bar})$	porosity	
	0.016	1000	1940(37)	-10.98(2)	60%	
	0.010	1100	2246(118)	-9.36(6)	37%	
	-0.003	1200	2277(155)	-8.09(5)	5.7%	
	0.006	1200	2551(48)	-7.94(2)	5.6%	

* Standard deviations are less than ± 1 unit of the last decimal.

** Volume ratio; standard deviations refer to gas-mixing variations over time.

21.8 °C (details in Table 1). A phase-pure product was obtained upon free cooling inside the furnace in the flowing atmosphere, the wetting of which was switched off below 1000 °C in order to exclude oxygen that would be needed to oxidize this large and highly sintered sample. Powder-diffraction checks were performed with a D5000 diffractometer and CuK_{α1} radiation.

Wet chemical analyses of oxygen content. The oxygen content in YBaMn₂O_{5+w} was determined iodometrically. A finely powdered sample of about 0.15 g and 0.8 g KI were weighed into a 10 mL glass ampoule, 5 mL H₂O was added and the ampoule was flushed by Ar. After adding 4 mL of concentrated hydrochloric acid, the ampoule was immediately sealed. The digestion of the samples lasted less than a minute when aided by heating at the bottom of the ampoule. The formed triiodide solutions were titrated under argon atmosphere with a 0.1 M thiosulfate solution of concentration standardized on KIO₃ of standard-substance quality (Merck). Close to the equivalence point, soluble starch was added as the indicator. Blank analyses were performed in order to correct for the small amount of iodine produced in the acidic environment employed for the digestion process.

Neutron powder diffraction (NPD). Data were collected at 25, 100, 200, 300, 400, 500 and 600 °C in the cryofurnace of the D2B instrument at the Institut Laue–Langevin, Grenoble, France. The two-axis diffractometer was equipped with 128 position-sensitive ³He counters, and the Ge monochromator was in orientation *hkl* 335. The wavelength $\lambda = 1.59421$ Å was standardized on CeO₂. Coarsely powdered YBaMn₂O_{5.006} (~10 g) was sealed in a vanadium container under He atmosphere. The primary collimation was removed from the setup in order to achieve higher intensity at the expense of resolution. The slit was set to 100/100 so that the beam size at the sample was 2×5 cm². The counting limit was 1.6×10^5 counts per step of 0.05° of the angle 2 θ . The angular range was $5-160^\circ$ (2θ).

Rietveld refinements. Crystal-structure refinements in the GSAS software suite [13] used the Chebyshev function of 18 parameters to model the background. The NPD peak profiles were fitted by Simpson integration of the Thompson–Cox–Hastings pseudo-Voigt profile with two Gaussian and one Lorentzian profile parameters refined together with the zero point, transparency and asymmetry.

Orbital energy calculations. The molecular-orbital energy levels for isolated manganese–oxygen coordinations were calculated from structural data with the CAESAR 2.0 (Crystal And Electronic Structure AnalyzeR) software package [14]. The basis-set coefficients for this Extended-Hückel tight-binding method were listed previously [6].

3. Results and discussion

The synthesis conditions of the highly sintered YBaMn₂O_{5.006(0)} NPD sample of ~10 g are listed in Table 1. The iodometric titration is well reproducible, and standard deviations of the oxygen content at the fourth decimal are not unusual. The non-stoichiometry span of the synthesis tests in Table 1 is close to the formation limits of this YBaMn₂O₅-type structure as a single phase under given synthesis conditions, also in agreement with results in Ref. [15]. The window in terms of p_{O2} in the reaction atmosphere is narrow. As an example, at the highest temperature tested, 1200 °C, increasing the Ar/H₂ mixing ratio to 9600 in the synthesis gas wetted at 22.8 °C yielded a solid solution YBaMn₂O_{5.67}, the



Fig. 1. P4/nmm unit cell of the paramagnetic YBaMn₂O₅ at 298 K.

lower composition limit of which is the $YBaMn_2O_{5.5}$ reported in Ref. [9] to coexist in equilibrium with $YBaMn_2O_5$.

The crystal structure of YBaMn₂O_{5.006} was refined in the space-group P4/nmm, suggested in Ref. [8] for the three-dimensional checkerboard order of Mn^{2+} and Mn^{3+} (Fig. 1). The minute amount of non-stoichiometric oxygen is located in the yttrium layer between the bases of the square pyramids, and refinements of its occupancy yielded values between 0.008 and 0.014, in good agreement with the 0.006 oxygen per formula obtained by chemical analysis after the sample synthesis. The refinement shows that the crystal structure keeps the P4/nmm space-group symmetry within the temperature range investigated, as did TbBaMn₂O₅ in the NPD study of the two-phase sample in Ref. [9]. Numerical overview of the structural evolution, with refinement statistics and the refined unit cell and atomic parameters is in Table 2.

Fig. 2 illustrates the practically linear and continuous thermal expansion of this structure on unit-cell parameters per single-perovskite subcell. This suggests absence of structural or spin-state transitions in the investigated range of temperatures.

Fig. 3 shows thermal evolution of apical and basal Mn–O bond distances in the two different Mn coordination square pyramids. Whereas the four equal basal Mn–O distances are the shortest bonds in the small pyramid of d^4 Mn³⁺ (a d_{z^2} -stabilizing Jahn–Teller distortion), they are longest in the large pyramid of d^5 Mn²⁺ (a $d_{x^2-y^2}$ -stabilizing Jahn–Teller distortion), and make an alternating pattern of small and large squares on the *ab* plane. All basal bonds expand linearly with temperature. The two ionic charges clearly remain ordered at all temperatures investigated.

The thermal evolution of the apical Mn–O bond lengths is less straightforward. At low temperatures, the smaller Mn^{3+} pyramid has its apical bond expanded so much by the d^4 Jahn–Teller distortion stabilizing the highest-occupied molecular orbital (HOMO) d_{z^2} as less antibonding, that it is longer than the apical distance in the adjacent larger Mn^{2+} pyramid. Fig. 1 illustrates that this is achieved by having Mn^{3+} much closer to the base of its pyramid than Mn^{2+} is to its base. Upon increasing temperature, the length difference of these two apical bonds decreases (Fig. 3). This thermally induced decrease in the d^4 Jahn–Teller distortion was also observed in other Mn^{3+} compounds such as

Table 2

Temperature evolution of the	YBaMn ₂ O ₅ structure para	meters* and refinement fig	gures of merit.
1	201		

T (K)	298	373	473	573	673	773	873
R _{wp}	0.046	0.045	0.043	0.044	0.043	0.043	0.045
Rp	0.035	0.034	0.033	0.034	0.033	0.033	0.035
$\hat{R_F^2}$	0.033	0.038	0.041	0.049	0.054	0.058	0.070
χ^2	3.258	3.162	3.002	3.094	3.083	3.153	3.461
a (Å)	5.54891(8)	5.55306(8)	5.55801(8)	5.56386(8)	5.56928(4)	5.57506(9)	5.58125(9)
c (Å)	7.65836(13)	7.67326(13)	7.68955(12)	7.70642(13)	7.72131(13)	7.73706(13)	7.75332(13)
V (Å ³)	235.80(10)	236.62(10)	237.54(10)	238.56(10)	239.49(10)	240.48(10)	241.52(11)
$z \mathrm{Mn}^{2+}$	0.74624(54)	0.74571(54)	0.74429(55)	0.74384(58)	0.74445(61)	0.74590(61)	0.74641(68)
$z { m Mn}^{3+}$	0.27565(53)	0.27493(52)	0.27458(54)	0.27413(58)	0.27360(63)	0.27398(65)	0.27379(74)
<i>x</i> ₀₁	0.49138(18)	0.49135(18)	0.49163(19)	0.49174(20)	0.49227(23)	0.49235(25)	0.49287(31)
<i>z</i> ₀₁	0.31668(10)	0.31670(10)	0.31677(10)	0.31679(11)	0.31653(11)	0.31678(12)	0.31676(13)
z ₀₂	0.00724(58)	0.00833(55)	0.00950(53)	0.00988(57)	0.00983(62)	0.00910(67)	0.00841(79)
O3 occ.	0.008(4)	0.014(4)	0.010(4)	0.010(4)	0.009(4)	0.008(4)	0.012(4)

* Space group *P4/nmm* number 129 and origin choice 2, Wyckoff positions: Y 2*b*, Ba 2*a*, Mn²⁺ 2*c*, Mn³⁺ 2*c*, O1 8*j*, O2 2*c*, O3 2*c* of *z* fixed to ½ and occupancy (occ.) refined. Observed reflections: 166, refined variables: 39.



Fig. 2. Thermal evolution of the unit-cell parameters per single-perovskite subcell. The values for 1.5 K are from Ref. [8].



Fig. 3. Thermal evolutions of bond lengths $Mn^{2+}-O$ (magenta) and $Mn^{3+}-O$ (brown) to the square-pyramidal basis (\blacksquare) and to the apex (\blacktriangle). The 1.5 K values are from Ref. [8]. Lines are guides for eye.

stoichiometric LaMnO₃, where it disappears at about 750 K [16,17]. In YBaMn₂O₅, the apical and basal Mn–O distances in each pyramid become closest at about 600 K (Fig. 3) due to this thermal alleviation of the Jahn–Teller distortion. From that point up, a slow mutual approach of the

two Mn valences starts, demonstrated by resumed contraction of the apical distance in the larger ${\rm Mn}^{2+}$ pyramid versus expansion in the smaller ${\rm Mn}^{3+}$ one, both upon the overall thermal expansion.

Fig. 4 shows the evolution of the bond-valence sums (BVS) of the two



Fig. 4. Thermal evolution of the bond-valence sums (BVS) at the two crystallographically different Mn atoms. The BVS for 1.5 K is calculated from data in Ref. [8].



Fig. 5. Evolution of energies of essentially nonbonding Mn orbitals calculated for $(Mn^{3+}O_5)^{7-}$ (brown) and $(Mn^{2+}O_5)^{8-}$ (magenta). Orbital symbols refer to local coordinate systems centered at Mn and oriented in the direction of the O atoms. The energies for 1.5 K are calculated from atomic coordinates in Ref. [8].

manganese atoms, calculated with the average Mn–O single-bond length 1.75 Å according to Ref. [18]. As expected, the BVS values are close to +2 and +3, and they do start a mutual approach around 600 K. Extrapolation of the four highest-temperature points in Fig. 4 gives an upper estimate for how far YBaMn₂O₅ may be from valence mixing.

The energies of the essentially nonbonding orbitals of the YBaMn₂O₅ high-spin [2,8] cations Mn²⁺ and Mn³⁺ in their ligand field were calculated with CAESAR 2.0 and are plotted in Fig. 5 as a function of temperature. In the small Mn³⁺ coordination pyramid, the Jahn–Teller expansion of the apical bond upon contraction of basal bonds destabilizes the empty $d_{x^2-y^2}$ orbital by as much as 4.9 eV from d_{z^2} at 300 K, and that makes the $Mn^{3+} d^4$ configuration behave as a *half*-filled orbital set. The apical oxide anion then mediates an antiferromagnetic semicovalent superexchange [10] with the truly half-filled spin-polarized orbital set of $Mn^{2+} d^5$ (Fig. 6). This is what yields below $T_N = 167$ K [8,19] the observed [8] G-type [3] coupling of all manganese spins. The basal oxygens of the pyramids mediate likewise. Their buckling angle Mn²⁺–O–Mn³⁺ in YBaMn₂O₅ (157.20(10)° at 298 K versus 156.75(14) at 873 K) is the same as in the charge-ordered phases YBaFe₂O₅ of 158.59(31)° at 20 K [6] or YBaCo₂O₅ of 156.30(88)° at 50 K [20], where the G-type coupling was identified by NPD in their charge-ordered state (on HoBaFe₂O₅ [21] and HoBaCo₂O₅ [22]). Also DFT calculations [23,

24] confirm the YBaMn₂O₅ G-type spin order. Only YBaMn₂O₅ turns into a ferrimagnet upon cooling, though, since its checkerboard charge order coincides with the G-type spin order, which makes all Mn^{2+} moments to point up (along *c* [8]) and all Mn^{3+} down.

Valence mixing via sharing the excess electron of Mn^{2+} by both manganese atoms that face each other in Fig. 1 would have to be accomplished via their d_{z^2} orbitals, as in YBaFe₂O₅ [25]. Fig. 6 illustrates how much the energies of the d_{z^2} orbital (in blue) for the electron of fixed spin orientation differ between Mn^{2+} and Mn^{3+} . In contrast, in the charge-ordered YBaFe₂O₅ of Fig. 7, the minority-spin Fe²⁺ electron has only 0.8 eV to the Fe²⁺ d_{z^2} orbital that is 0.2 eV below the d_{z^2} of Fe³⁺ to mix with and thus couple these two Fe ferromagnetic (blue arrow) by this electron sharing across the Y layer [6].

4. Conclusions

Synthesis conditions to obtain single-phase $YBaMn^{II}Mn^{II}O_5$ by sintering at high-temperatures make a narrow window of oxygen partial pressures in the reaction atmosphere (Table 1). Likewise narrow is the range in temperatures. Whereas sintering at 1000 °C yields porosity of 60%, the silvery product at 1200 °C has 6%. The crystal structure in the paramagnetic state, refined from neutron powder diffraction data



Fig. 6. Charge-ordered YBaMn₂O₅. Semicovalent linear superexchange interaction of effectively half-filled d^4 Mn³⁺ and half-filled d^5 Mn²⁺ cations. The spin-pairing energy $U_{\rm C}$ is drawn to permit the observed [8] Mn²⁺ high-spin state. The two d_{z^2} orbitals to eventually share the valence-mixing electron at high temperature are marked in blue.



Fig. 7. Charge-ordered YBaFe₂O₅ on the scale of Fig. 6. Semicovalent linear superexchange interaction of half-filled $d^5 \text{ Fe}^{3+}$ and of $d^6 \text{ Fe}^{2+}$ cations. The spinpairing energy $U_{\rm C}$ is drawn to permit the observed [6] Fe²⁺ high-spin state. In blue: Upon demise of the charge ordering, the minority-spin Fe²⁺ electron in the ordered d_{xz} orbital will mix via d_{z^2} orbitals (blue shapes) of the two Fe atoms facing each other across their pyramidal bases and couple them ferromagnetic (blue double arrow) in the still antiferromagnetic crystal [6].

collected between 25 and 600 °C, maintains Mn^{2+} and Mn^{3+} ordered in a three-dimensional checkerboard pattern of electrostatic minimization; a pattern that in YBaFe₂O₅ is modified by orbital ordering [6].

At ambient temperature, the presence of one electron in the Mn^{3+} HOMO d_{z^2} expands the apical Mn–O distance of this small coordination pyramid due to Jahn–Teller distortion. The apical distance of the larger Mn^{2+} square pyramid is contracted due to corresponding expansion of its basal Mn–O distances that stabilizes the Mn^{2+} HOMO $d_{x^2-y^2}$. Upon warming to ~600 K, the apical and basal Mn–O distances become closest both in the larger Mn^{2+} and in the smaller Mn^{3+} pyramid (Fig. 3), due to thermal alleviation of the Jahn–Teller distortion. Above this temperature, the Mn^{2+} –O apical distance begins to contract and Mn^{3+} –O expand upon the overall thermal expansion. From this point, the BVS at Mn^{2+} and Mn^{3+} start slowly approaching one another as a first sign of a possible valence mixing at a higher temperature.

Extended-Hückel calculations of the orbital energies for the isolated square-pyramidal coordinations around Mn^{3+} and Mn^{2+} bring further details. The very high energy of the Mn^{3+} LUMO $d_{x^2-y^2}$ makes this electron configuration effectively half-filled. The semicovalent super-exchange interaction [10] of such d^4 Mn^{3+} and d^5 Mn^{2+} half-filled

electron configurations via the intervening oxygen then yields antiferromagnetic coupling of the two unequal-spin manganese atoms. Furthermore, the calculations show that the energies of the d_{z^2} electron of the same spin orientation differ profoundly between Mn^{2+} and Mn^{3+} , and this is what makes it hard to mix thermally the Mn^{2+} excess electron with Mn^{3+} in YBaMn₂O₅ when compared with the easy mixing in YBaFe₂O₅.

Interestingly, combining Mn and Co in the prototype double-cell perovskite structure does not simulate YBaFe₂O₅. There is no charge ordering, no valence mixing, just a Co²⁺ and Mn³⁺ solid solution of d^7-d^4 configurations [26]. The reason is redox chemistry. If the product were YBaCo^{III}Mn^{II}O₅, the Co³⁺ and Mn²⁺ d^6-d^5 configuration would have behaved as it does in YBaFe^{II}Fe^{III}O₅.

CRediT authorship contribution statement

Pavel Karen: Conceptualization, Data curation, Formal analysis, Investigation, Methodology, Project administration, Resources, Software, Validation, Visualization, Writing – original draft, Writing – review & editing.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

Data will be made available on request.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at https://do i.org/10.1016/j.jssc.2022.123469.

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