



Constraints on off-shell Higgs boson production and the Higgs boson total width in $ZZ \rightarrow 4\ell$ and $ZZ \rightarrow 2\ell 2\nu$ final states with the ATLAS detector



The ATLAS Collaboration*

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ABSTRACT

A measurement of off-shell Higgs boson production in the $ZZ \rightarrow 4\ell$ and $ZZ \rightarrow 2\ell 2\nu$ decay channels, where ℓ stands for either an electron or a muon, is performed using data from proton–proton collisions at a centre-of-mass energy of $\sqrt{s} = 13$ TeV. The data were collected by the ATLAS experiment in 2015 and 2016 at the Large Hadron Collider, and they correspond to an integrated luminosity of 36.1 fb^{-1} . An observed (expected) upper limit on the off-shell Higgs signal strength, defined as the event yield normalised to the Standard Model prediction, of 3.8 (3.4) is obtained at 95% confidence level (CL). Assuming the ratio of the Higgs boson couplings to the Standard Model predictions is independent of the momentum transfer of the Higgs production mechanism considered in the analysis, a combination with the on-shell signal-strength measurements yields an observed (expected) 95% CL upper limit on the Higgs boson total width of 14.4 (15.2) MeV.

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1. Introduction

The observation of the Higgs boson by the ATLAS and CMS experiments [1,2] at the Large Hadron Collider (LHC) marks a milestone towards the understanding of the mechanism of electroweak (EW) symmetry breaking [3–5]. Further studies of the spin, parity and couplings of the new particle have shown no significant deviation from the predictions for the Standard Model (SM) Higgs boson [6–10]. Efforts to measure the properties of the Higgs boson are primarily focused on on-shell production. For a Higgs boson at a mass of 125 GeV [10,11], the expected natural width of the SM Higgs boson is $\Gamma_H^{\text{SM}} \sim 4.1$ MeV [12]. However, above 125 GeV off-shell production of the Higgs boson has a substantial cross section at the LHC [13–16], due to the increased phase space as the vector bosons ($V = W, Z$) and top quark decay products become on-shell with the increasing energy scale. This provides an opportunity to study the Higgs boson properties at higher energy scales. Off-shell production can provide sensitivity to new physics that alters the interactions between the Higgs boson and other fundamental particles in the high-mass region [17–24].

The measured off-shell event yield from gluon–gluon fusion (ggF) production normalised to the SM prediction, where this

ratio is referred to as the signal strength $\mu_{\text{off-shell}}$, can be expressed as

$$\mu_{\text{off-shell}} = \frac{\sigma_{\text{off-shell}}^{gg \rightarrow H^* \rightarrow ZZ}}{\sigma_{\text{off-shell,SM}}^{gg \rightarrow H^* \rightarrow ZZ}} = \kappa_{g,\text{off-shell}}^2 \cdot \kappa_{Z,\text{off-shell}}^2,$$

where $\sigma_{\text{off-shell}}^{gg \rightarrow H^* \rightarrow ZZ}$ is the cross section of the off-shell Higgs boson production via ggF with subsequent decay into a ZZ pair, and $\kappa_{g,\text{off-shell}}$ and $\kappa_{Z,\text{off-shell}}$ are the off-shell coupling modifiers relative to the SM predictions associated with the $gg \rightarrow H^*$ production and the $H^* \rightarrow ZZ$ decay, respectively. The off-shell Higgs boson signal cannot be treated independently of the $gg \rightarrow ZZ$ background, as sizeable negative interference effects appear [13]. The interference term is assumed to be proportional to $\sqrt{\mu_{\text{off-shell}}} = \kappa_{g,\text{off-shell}} \cdot \kappa_{Z,\text{off-shell}}$. Similarly, $\mu_{\text{on-shell}}$ for the on-shell Higgs boson production via ggF is given by:

$$\mu_{\text{on-shell}} = \frac{\sigma_{\text{on-shell}}^{gg \rightarrow H \rightarrow ZZ^*}}{\sigma_{\text{on-shell,SM}}^{gg \rightarrow H \rightarrow ZZ^*}} = \frac{\kappa_{g,\text{on-shell}}^2 \cdot \kappa_{Z,\text{on-shell}}^2}{\Gamma_H / \Gamma_H^{\text{SM}}},$$

which depends on the Higgs boson total width Γ_H . A measurement of the relative off-shell and on-shell event yields, $\mu_{\text{off-shell}} / \mu_{\text{on-shell}}$, provides direct information about Γ_H , if one assumes identical on-shell and off-shell Higgs boson coupling modi-

* E-mail address: atlas.publications@cern.ch.

fiers [15,25]. The above formalism describing the ratio of off-shell to on-shell cross sections also applies to the vector-boson fusion (VBF) production mode. As in the previous measurement [26], for a measurement of Γ_H it is necessary to assume that the on-shell and off-shell coupling modifiers are the same, and for an upper limit that the on-shell coupling modifiers are not larger than the off-shell couplings. It is also assumed that any new physics which modifies the off-shell signal strength and the off-shell couplings does not modify the relative phase of the interfering signal and background processes. Further, it is assumed that there are neither sizeable kinematic modifications to the off-shell signal nor new sizeable signals in the search region of this analysis unrelated to an enhanced off-shell signal strength.

The ATLAS and CMS experiments have presented studies of the off-shell production of the Higgs boson using Run-1 proton–proton (pp) collisions data [26–29]. ATLAS obtained an observed (expected) upper limit on the off-shell Higgs boson signal strength ($\mu_{\text{off-shell}}$) in the range of 5.1–8.6 (6.7–11.0) [26], using the ZZ and WW channels. This range is determined by the assumption that the $gg \rightarrow ZZ$ and $gg \rightarrow WW$ background K -factors, corresponding to the ratio of the next-to-leading-order (NLO) QCD predictions to the leading-order (LO) predictions, lie between one-half and twice the value of the $gg \rightarrow H^* \rightarrow ZZ(WW)$ signal K -factor. An observed (expected) 95% confidence level (CL) upper limit of $\Gamma_H < 23(33)$ MeV was obtained, assuming the $gg \rightarrow ZZ(WW)$ background K -factor is equal to the $gg \rightarrow H^* \rightarrow ZZ(WW)$ signal K -factor. CMS presented a similar study in the ZZ and WW channels, with observed (expected) 95% CL upper limit of $\Gamma_H < 13(26)$ MeV [29]. By comparison, the precision of Γ_H from direct on-shell Higgs boson mass measurements alone is approximately 1 GeV [9,30,31], limited by measurement resolution.

This Letter presents an analysis of off-shell Higgs boson production in the $ZZ \rightarrow 4\ell$ and $ZZ \rightarrow 2\ell 2\nu$ final states ($\ell = e, \mu$), using 36.1 fb^{-1} of data collected by the ATLAS detector in pp collisions at $\sqrt{s} = 13$ TeV. The off-shell region is defined by requiring the invariant mass of the ZZ system (m_{ZZ}) to be above the on-shell ZZ production threshold, hence well above the Higgs boson mass, and the on-shell region is defined by a mass window around the 125 GeV resonance. This analysis adopts the same methodology used in the Run-1 analysis reported in Ref. [26]. The analysis for the $ZZ \rightarrow 4\ell$ final state closely follows the Higgs boson measurements and high-mass search in the same final state described in Refs. [32,33]. The off-shell Higgs signal strength is extracted using a matrix-element discriminant, defined in Section 4, in a mass region $220 \text{ GeV} < m_{4\ell} < 2000 \text{ GeV}$. The on-shell signal strength was measured in the $118 \text{ GeV} < m_{4\ell} < 129 \text{ GeV}$ region in Ref. [32]. The analysis of the $ZZ \rightarrow 2\ell 2\nu$ channel, described in Section 5, follows a strategy similar to that used in the search for heavy ZZ resonances described in Ref. [33]. For this channel, the signal strength is extracted from the transverse mass distribution in the 250 to 2000 GeV range. For off-shell production of the Higgs boson, the dominant processes of ggF and VBF are considered. Next-to-next-to-leading-order (NNLO) QCD and NLO EW corrections are known for the off-shell signal process $gg \rightarrow H^* \rightarrow ZZ$ [25]. More recently, NLO QCD corrections have also become available for the $gg \rightarrow ZZ$ background and for the signal–background interference [34,35], for which additional details are given in Section 3. Given that the QCD corrections for the off-shell signal processes have only been calculated inclusively in the jet multiplicity, the analysis is performed inclusively in jet observables and the event selection is designed to minimise the dependence on the momentum of the ZZ system, which is sensitive to the jet multiplicity.

2. ATLAS detector

The ATLAS experiment is described in Ref. [36]. ATLAS is a multipurpose detector with a forward–backward symmetric cylindrical geometry and a solid-angle¹ coverage of nearly 4π . The inner tracking detector, covering the region $|\eta| < 2.5$, consists of a silicon pixel detector, a silicon microstrip detector and a straw-tube transition-radiation tracker. The innermost layer of the pixel detector, the insertable B-layer [37], was installed between Run 1 and Run 2 of the LHC. The inner detector is surrounded by a thin superconducting solenoid providing a 2T magnetic field, and by a finely segmented lead/liquid-argon (LAr) electromagnetic calorimeter covering the region $|\eta| < 3.2$. A steel/scintillator-tile hadronic calorimeter provides coverage in the central region $|\eta| < 1.7$. The endcap and forward regions, covering the pseudorapidity range $1.5 < |\eta| < 4.9$, are instrumented with electromagnetic and hadronic LAr calorimeters, with copper or tungsten as the absorber material. A muon spectrometer system incorporating large superconducting toroidal air-core magnets surrounds the calorimeters. Three layers of precision wire chambers provide muon tracking in the range $|\eta| < 2.7$, while dedicated fast chambers are used for triggering in the region $|\eta| < 2.4$. The trigger system is composed of two stages [38]. The first stage, implemented with custom hardware, uses information from calorimeters and muon chambers to reduce the event rate from about 40 MHz to a maximum of 100 kHz. The second stage, called the high-level trigger, reduces the data acquisition rate to about 1 kHz on average. The high-level trigger is software-based and runs reconstruction algorithms similar to those used in the offline reconstruction.

3. Monte Carlo simulation and higher-order theory corrections

Monte Carlo (MC) samples of $gg \rightarrow (H^* \rightarrow)ZZ$ events, which include the SM Higgs boson signal, $gg \rightarrow H^* \rightarrow ZZ$, the continuum background, $gg \rightarrow ZZ$, and the signal–background interference contribution, were generated with the MC generator SHERPA-v2.2.2 + OPENLOOPS [39–42]. Matrix elements were calculated for zero jets and one jet at LO and merged with the SHERPA parton shower [43]. The NNPDF30NNLO [44] PDF set was used, and the QCD renormalisation and factorisation scales were set to $m_{ZZ}/2$.

The K -factor for the $gg \rightarrow H^* \rightarrow ZZ$ process is known up to NNLO in QCD as a function of m_{ZZ} [12,25]. More recently, a NLO QCD calculation which includes the $gg \rightarrow ZZ$ continuum process has become available [34,35] allowing m_{ZZ} differential K -factors to be calculated with an expansion in the inverse top mass ($1/m_t$) below $2m_t$, and assuming a massless-quark approximation above this threshold. This NLO QCD calculation was used to correct all three components with separate K -factors computed for the signal $gg \rightarrow H^* \rightarrow ZZ$ ($K^S(m_{ZZ})$), the background $gg \rightarrow ZZ$ ($K^B(m_{ZZ})$) and the interference ($K^I(m_{ZZ})$). Since the NNLO QCD correction is only known differentially in m_{ZZ} for the $gg \rightarrow H^* \rightarrow ZZ$ process and not for all three components in the off-shell region, an overall correction is applied by scaling the differential NLO QCD reweighted cross section by an additional factor of 1.2, which is assumed to be the same for the signal, background and interference. This additional constant scale factor is justified by the constant NNLO to

¹ The ATLAS experiment uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$.

NLO ratio of the QCD predictions over the data region considered in the analysis. Using these scaled NLO K -factors, the cross section for the $gg \rightarrow (H^* \rightarrow)ZZ$ process with any off-shell Higgs boson signal strength $\mu_{\text{off-shell}}$ can be obtained from a parameterisation of three SM MC samples: the $gg \rightarrow H^* \rightarrow ZZ$ signal ($\sigma_{gg \rightarrow H^* \rightarrow ZZ}^{\text{SM}}$), the $gg \rightarrow ZZ$ continuum background ($\sigma_{gg \rightarrow ZZ, \text{cont}}^{\text{SM}}$) and the full process with signal, background and interference $gg \rightarrow (H^* \rightarrow)ZZ$ ($\sigma_{gg \rightarrow (H^* \rightarrow)ZZ}^{\text{SM}}$), where the last sample is required to derive the interference sample:

$$\begin{aligned} \sigma_{gg \rightarrow (H^* \rightarrow)ZZ}(\mu_{\text{off-shell}}) &= \mu_{\text{off-shell}} \cdot 1.2 \cdot K^S(m_{ZZ}) \cdot \sigma_{gg \rightarrow H^* \rightarrow ZZ}^{\text{SM}} \\ &+ \sqrt{\mu_{\text{off-shell}}} \cdot 1.2 \cdot K^I(m_{ZZ}) \cdot \sigma_{gg \rightarrow ZZ, \text{Interference}}^{\text{SM}} \\ &+ 1.2 \cdot K^B(m_{ZZ}) \cdot \sigma_{gg \rightarrow ZZ, \text{cont}}^{\text{SM}} \end{aligned} \quad (1)$$

$$\begin{aligned} \sigma_{gg \rightarrow ZZ, \text{Interference}}^{\text{SM}} &= \sigma_{gg \rightarrow (H^* \rightarrow)ZZ}^{\text{SM}} - \sigma_{gg \rightarrow H^* \rightarrow ZZ}^{\text{SM}} - \sigma_{gg \rightarrow ZZ, \text{cont}}^{\text{SM}} \end{aligned} \quad (2)$$

The electroweak $pp \rightarrow VV + 2j$ processes containing both the VBF-like events and events from associated Higgs production with vector bosons (VH), which includes on-shell Higgs boson production, were simulated using MADGRAPH5_aMC@NLO [45] with matrix elements calculated at LO. The QCD renormalisation and factorisation scales were set to m_W following the recommendation in Ref. [46] and the NNPDF23LO PDF set [47] was used. PYTHIA 8.186 [48] was used for parton showering and hadronisation, with the A14 set of tuned parameters for the underlying event [49]. Due to the different Γ_H dependence, the on-shell and off-shell Higgs boson production processes are separated when weighting MC events as in Eqs. (1) by requiring that the generated Higgs boson mass satisfy $|m_H^{\text{gen}} - 125 \text{ GeV}| < 1 \text{ GeV}$. This requirement is fully efficient in selecting the on-shell VH process. The cross section $\sigma_{pp \rightarrow VV + 2j}(\mu_{\text{off-shell}})$ for the electroweak $pp \rightarrow VV + 2j$ process for any off-shell Higgs boson signal strength $\mu_{\text{off-shell}}$ is parameterised in the same way as for the $gg \rightarrow (H^* \rightarrow)ZZ$ process.

The $q\bar{q} \rightarrow ZZ$ background was simulated with SHERPA v2.2.2, using the NNPDF30NNLO PDF set for the hard-scattering process. NLO QCD accuracy is achieved in the matrix-element calculation for 0- and 1-jet final states and LO accuracy for 2- and 3-jet final states. The merging with the SHERPA parton shower was performed using the MePs@NLO prescription. NLO EW corrections are applied as a function of the particle-level m_{ZZ} [50,51].

The WW and WZ backgrounds were simulated at NLO in QCD using the POWHEG-Box v2 event generator [52] with the CT10NLO PDF set [53] and PYTHIA 8.186 for parton showering and hadronisation. The non-perturbative effects were modelled with the AZNLO set of tuned parameters [54]. The interference between the $q\bar{q} \rightarrow ZZ$ and $q\bar{q} \rightarrow WW$ processes for the $2\ell 2\nu$ final state is found to be negligible and thus is not considered.

Events containing a single Z boson with associated jets ($Z + \text{jets}$) were simulated using the SHERPA v2.2.1 event generator. Matrix elements were calculated for up to two partons at NLO and four partons at LO using the COMIX [55] and OPENLOOPS [41] matrix-element generators and merged with the SHERPA parton shower [43] using the MePs@NLO prescription. The NNPDF30NNLO PDF set was used in conjunction with dedicated parton-shower tuning developed by the SHERPA authors. The $Z + \text{jets}$ events are normalised using the NNLO cross sections [56].

The triboson backgrounds ZZZ , WZZ , and WWZ with fully leptonic decays and at least four prompt charged leptons were modelled using SHERPA v2.2.1. The contribution from triboson backgrounds with one W or Z boson decaying hadronically is not

included in the simulation, but the impact on the analysis is found to be negligible. For the fully leptonic $t\bar{t} + Z$ background, with four prompt charged leptons originating from the decays of the top quarks and Z boson, MADGRAPH5_aMC@NLO was used. The $t\bar{t}$ background, as well as the single-top and Wt production, were modelled using POWHEG-Box v2 interfaced to PYTHIA 6.428 [57] with the Perugia 2012 [58] set of tuned parameters for parton showering, hadronisation and the underlying event, and to EVTGEN v1.2.0 [59] for properties of the bottom and charm hadron decays.

The particle-level events produced by each MC event generator were processed through the ATLAS detector simulation within the GEANT 4 framework [60,61] or the fast detector simulation package Atfast-II [61]. Additional pp interactions in the same and nearby bunch crossings (pile-up) are included in the simulation. The pile-up events were generated using PYTHIA 8 with the A2 set of tuned parameters [62] and the MSTW2008LO PDF set [63]. The simulation samples were weighted to reproduce the observed distribution of the mean number of interactions per bunch crossing in the data.

4. $ZZ \rightarrow 4\ell$ analysis

The analysis for the $ZZ \rightarrow 4\ell$ final state closely follows the on-shell Higgs boson measurements and high-mass search in the same final state described in Refs. [32,33], with the same event reconstruction, trigger and event selections, and background estimation methods. A matrix-element-based (ME-based) discriminant computed at LO is constructed to enhance the separation between the $gg \rightarrow H^* \rightarrow ZZ$ signal and the $gg \rightarrow ZZ$ and $q\bar{q} \rightarrow ZZ$ backgrounds, and this discriminant is subsequently used in a binned maximum-likelihood fit for the final result. To minimise the dependence of the $gg \rightarrow ZZ$ kinematics on higher-order QCD effects, the analysis is performed inclusively, ignoring the number of jets in the events.

The analysis is split into three channels (4μ , $2e2\mu$, $4e$). Each electron (muon) must have transverse momentum $p_T > 7$ (5) GeV and be measured in the pseudorapidity range $|\eta| < 2.47$ ($|\eta| < 2.7$). The highest- p_T lepton in the quadruplet must satisfy $p_T > 20$ GeV, and the second (third) lepton in p_T order is required to have $p_T > 15$ GeV ($p_T > 10$ GeV). Lepton pairs are formed from same-flavour opposite-charge leptons. For each channel, the quadruplet with a lepton pair whose mass is closest to the Z boson mass is kept. This pair is referred to as the leading dilepton pair and its invariant mass, m_{12} , is required to be between 50 GeV and 106 GeV. The second (subleading) pair is chosen from the remaining leptons as the pair closest in mass to the Z boson and in the range $50 \text{ GeV} < m_{34} < 115 \text{ GeV}$. The off-shell region is defined as the range $220 \text{ GeV} < m_{4\ell} < 2000 \text{ GeV}$, while the on-shell region is defined as $118 \text{ GeV} < m_{4\ell} < 129 \text{ GeV}$.

The dominant background in the $ZZ \rightarrow 4\ell$ channel arises from $q\bar{q} \rightarrow ZZ$ events. This is modelled using MC simulation, accurate to NLO QCD and NLO EW corrections as explained in Section 3. Other backgrounds, such as triboson production, $t\bar{t}V$, $Z + \text{jets}$, and top quark production, constitute less than 2% of the total background in the off-shell signal region, and are either taken from simulation or from dedicated data control regions.

Fig. 1(a) shows the observed and expected distributions of $m_{4\ell}$ combining all lepton channels in the off-shell region. The data are in agreement with the SM predictions, with two small excesses at $m_{4\ell}$ around 240 GeV and 700 GeV, each having a significance of about two standard deviations (2σ), as evaluated by the high-mass resonance search reported in Ref. [33]. Table 1 shows the expected and observed numbers of events in the signal region and additionally in the $400 \text{ GeV} < m_{4\ell} < 2000 \text{ GeV}$ mass range, which is

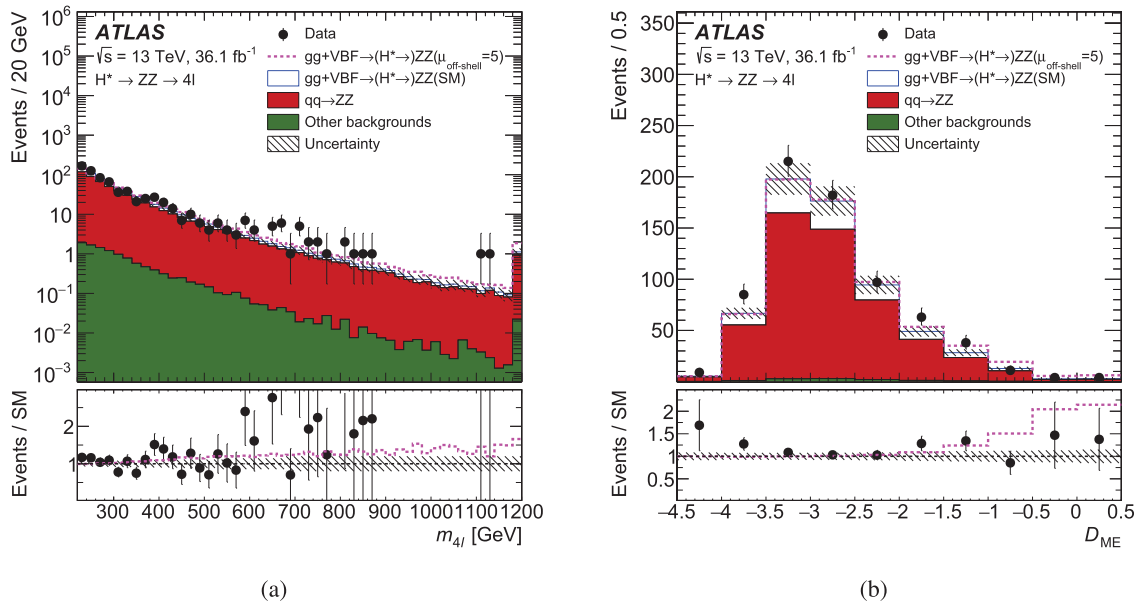


Fig. 1. Observed distributions in the range $220 \text{ GeV} < m_{4\ell} < 2000 \text{ GeV}$ for (a) the four-lepton invariant mass $m_{4\ell}$ and (b) the ME-based discriminant D_{ME} combining all lepton final states, compared to the expected contributions from the SM including the Higgs boson (stacked). Events with $m_{4\ell} > 1200 \text{ GeV}$ are included in the last bin of the $m_{4\ell}$ distribution. The hatched area shows the combined statistical and systematic uncertainties. The dashed line corresponds to the total expected event yield, including all backgrounds and the Higgs boson with $\mu_{\text{off-shell}} = 5$. The ratio plot shows the observed data yield divided by the SM prediction (black points) as well as the total expected event yield with $\mu_{\text{off-shell}} = 5$ divided by the SM prediction (dashed line) in each bin.

Table 1

The expected and observed numbers of events in the signal region for both final states. For the $ZZ \rightarrow 4\ell$ analysis, numbers are given for both the signal region and a signal-enriched region which covers the mass range $400 \text{ GeV} < m_{4\ell} < 2000 \text{ GeV}$. The other backgrounds in the $ZZ \rightarrow 4\ell$ final state include contributions from Z + jets and top quark processes, while in the $ZZ \rightarrow 2\ell 2\nu$ final state they include contributions from tri-boson production, the W + jets process, and top quark processes other than pair production. For the $ZZ \rightarrow 2\ell 2\nu$ analysis, the range $250 \text{ GeV} < m_{\tau\tau}^{ZZ} < 2000 \text{ GeV}$ is considered. The upper part of the table contains the expected events for the $gg \rightarrow (H^* \rightarrow)ZZ$ and VBF ($H^* \rightarrow)ZZ$ processes which include the Higgs boson signal, background and interference for the SM predictions. The SM estimates for the signal (S) and background (B) event yields without interference are given in parentheses. The lower part of the table contains the corresponding predictions for $\mu_{\text{off-shell}} = 5$. The uncertainties in the number of expected events include the statistical uncertainties from MC samples and systematic uncertainties, summed in quadrature. Empty entries correspond to contributions with event yields smaller than 0.1 events.

Process	$ZZ \rightarrow 4\ell$		$ZZ \rightarrow 2\ell 2\nu$
	$m_{4\ell} > 220 \text{ GeV}$	$m_{4\ell} > 400 \text{ GeV}$	$m_{\tau\tau}^{ZZ} > 250 \text{ GeV}$
$gg \rightarrow (H^* \rightarrow)ZZ$	96 ± 15	10.6 ± 2.0	22 ± 4
($gg \rightarrow H^* \rightarrow ZZ$ (S))	9.8 ± 1.5	5.9 ± 1.0	20.1 ± 3.3
($gg \rightarrow ZZ$ (B))	101 ± 16	11.8 ± 2.2	28 ± 6
VBF ($H^* \rightarrow)ZZ$	8.29 ± 0.34	3.07 ± 0.13	2.83 ± 0.14
(VBF $H^* \rightarrow ZZ$ (S))	1.67 ± 0.08	1.14 ± 0.04	5.45 ± 0.30
(VBF ZZ (B))	9.9 ± 0.4	4.17 ± 0.18	6.92 ± 0.35
$q\bar{q} \rightarrow ZZ$	520 ± 42	77 ± 8	132 ± 15
$q\bar{q} \rightarrow WZ$	–	–	68 ± 4
$WW/t\bar{t}/Wt/Z \rightarrow \tau\tau$	–	–	2.6 ± 1.0
Z + jets	–	–	6.0 ± 2.8
Other backgrounds	14.6 ± 0.7	2.15 ± 0.15	1.14 ± 0.08
Total Expected (SM)	639 ± 60	93 ± 10	234 ± 16
Observed	704	114	261
Other signal hypothesis			
$gg \rightarrow (H^* \rightarrow)ZZ$ ($\mu_{\text{off-shell}} = 5$)	117 ± 18	26 ± 5	61 ± 12
VBF ($H^* \rightarrow)ZZ$ ($\mu_{\text{off-shell}} = 5$)	11.0 ± 0.5	4.85 ± 0.22	8.8 ± 0.4

enriched in signal. The latter mass region was chosen for this table since it is optimal for a counting experiment.

The matrix-element kinematic discriminant fully exploits the event kinematics in the centre-of-mass frame of the 4ℓ system. It is computed from eight kinematic observables: the three masses $m_{4\ell}$, m_{12} and m_{34} , and the leading Z boson production angle and four decay angles defined in Ref. [64]. These observables are

used to calculate the matrix elements for the different processes with the MCFM program [15] at LO. The following matrix elements are calculated for each event in the mass range $220 \text{ GeV} < m_{4\ell} < 2000 \text{ GeV}$:

- $P_{q\bar{q}}$: the matrix element squared for the $q\bar{q} \rightarrow ZZ \rightarrow 4\ell$ process,

- P_{gg} : the matrix element squared for the $gg \rightarrow (H^* \rightarrow) ZZ \rightarrow 4\ell$ process, which includes the Higgs boson with SM couplings, the continuum background and their interference,
- P_H : the matrix element squared for the $gg \rightarrow H^* \rightarrow ZZ \rightarrow 4\ell$ process without continuum background or interference.

The ME-based discriminant is defined as in Ref. [15]:

$$D_{\text{ME}} = \log_{10} \left(\frac{P_H}{P_{gg} + c \cdot P_{q\bar{q}}} \right),$$

where $c = 0.1$ is a constant whose value is chosen to balance the overall cross sections of the $q\bar{q} \rightarrow ZZ$ and $gg \rightarrow (H^* \rightarrow) ZZ$ processes. The value of c has a small effect on the analysis sensitivity. Fig. 1(b) shows the observed and expected distributions of D_{ME} . Events with a D_{ME} value between -4.5 and 0.5 are used for the final result.

5. $ZZ \rightarrow 2\ell 2\nu$ analysis

The analysis in the $ZZ \rightarrow 2\ell 2\nu$ final state closely follows the one performed to search for ZZ resonances [33]. The reconstruction, identification and selection of electrons, muons, jets, b -jets and missing transverse momentum are identical while the event selection is optimised for the current analysis.

To discriminate the signal from the background and enhance the sensitivity to off-shell Higgs boson production, the transverse mass of the ZZ system (m_{T}^{ZZ}) is used, defined as:

$$m_{\text{T}}^{ZZ} \equiv \sqrt{\left[\sqrt{m_Z^2 + (p_{\text{T}}^{\ell\ell})^2} + \sqrt{m_Z^2 + (E_{\text{T}}^{\text{miss}})^2} \right]^2 - \left| \vec{p}_{\text{T}}^{\ell\ell} + \vec{E}_{\text{T}}^{\text{miss}} \right|^2},$$

where $p_{\text{T}}^{\ell\ell}$ is the transverse momentum of the dilepton system, m_Z is the mass of the Z boson fixed to $m_Z = 91.187$ GeV [65] and $E_{\text{T}}^{\text{miss}}$ is the magnitude of the missing transverse momentum $\vec{E}_{\text{T}}^{\text{miss}}$. The latter is computed as the negative sum of transverse momenta of all the leptons and jets, as well as the tracks originating from the primary vertex but not associated with any of the leptons or jets, the so-called soft term.

The event selection is designed to minimise the dependence on the p_{T} of the ZZ system, and thus is performed inclusively in number of jets. First, events with two opposite-charge leptons of the same flavour are selected with the requirement of $p_{\text{T}} > 30$ (20) GeV for the leading (sub-leading) lepton. The dilepton invariant mass $m_{\ell\ell}$ is required to be in the range 76 GeV $< m_{\ell\ell} < 106$ GeV. Events with additional, loosely identified leptons with $p_{\text{T}} > 7$ GeV are rejected to reduce the amount of WZ background. The two Z bosons originating from the decay of an off-shell Higgs boson are boosted and tend to be back-to-back in the transverse plane. A series of selection requirements are applied to reduce the $Z + \text{jets}$ background: the two leptons are required to be produced with an angular separation of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} < 1.8$; $E_{\text{T}}^{\text{miss}}$ is required to be larger than 175 GeV; the azimuthal angle between the transverse momentum of the dilepton system and the missing transverse momentum is required to be large, $\Delta\phi(\vec{p}_{\text{T}}^{\ell\ell}, \vec{E}_{\text{T}}^{\text{miss}}) > 2.7$; $p_{\text{T}}^{\ell\ell}$ is required to be balanced by the $E_{\text{T}}^{\text{miss}}$ and jets, $|p_{\text{T}}^{\text{miss,jet}} - p_{\text{T}}^{\ell\ell}|/p_{\text{T}}^{\ell\ell} < 0.2$, where $p_{\text{T}}^{\text{miss,jet}} = |\vec{E}_{\text{T}}^{\text{miss}} + \sum_{\text{jet}} \vec{p}_{\text{T}}^{\text{jet}}|$; and $E_{\text{T}}^{\text{miss}}/H_{\text{T}} > 0.33$, where H_{T} is the scalar sum of lepton and jet transverse momenta. Finally, events with a b -jet with $p_{\text{T}} > 20$ GeV and $|\eta| < 2.5$, identified by the MV2c10 algorithm [66,67] with 85% tagging efficiency, are vetoed to suppress the top quark background.

The dominant backgrounds in the $ZZ \rightarrow 2\ell 2\nu$ channel consist of $q\bar{q} \rightarrow ZZ$ events, followed by WZ events. Other background processes with two genuine leptons not directly originating from

a Z boson decay include WW , $t\bar{t}$, Wt and $Z \rightarrow \tau\tau$. The remaining background comes from $Z + \text{jets}$ with poorly reconstructed $E_{\text{T}}^{\text{miss}}$, $W + \text{jets}$ events with at least one misidentified electron or muon, semileptonic top decays, and multi-jet events.

The $q\bar{q} \rightarrow ZZ$ background is modelled in the same manner as for the $ZZ \rightarrow 4\ell$ channel. The WZ background is estimated with simulation using a normalisation correction factor extracted from a dedicated control region (CR). This WZ -enriched CR is defined by selecting $Z \rightarrow \ell\ell$ candidates with an additional electron or muon with $p_{\text{T}} > 20$ GeV. Events with a b -jet are rejected to suppress leptonic $t\bar{t}$ decays and a $m_{\text{T}}(W) > 60$ GeV requirement is applied to reduce the $Z + \text{jets}$ contamination. The correction factor is then calculated in the CR as the number of data events, after subtracting the non- WZ contributions, divided by the predicted WZ yield, and is found to be 1.29. The statistical uncertainty of the WZ estimate is about 2%, while the systematic uncertainty is estimated to be 5% from theoretical and experimental uncertainties in the simulation-based transfer factor between the three-lepton control region and the two-lepton signal region.

The non-resonant- $\ell\ell$ background, including WW , $t\bar{t}$, Wt and $Z \rightarrow \tau\tau$ processes, is estimated from a control sample of $e\mu$ events, in the same manner as in Ref. [33], except that the $e\mu$ CR is defined by requiring $E_{\text{T}}^{\text{miss}} > 120$ GeV. The background estimation is performed by extrapolating the result obtained with the relaxed $E_{\text{T}}^{\text{miss}}$ requirement to the SR, extracting the efficiency of the $E_{\text{T}}^{\text{miss}}$ selection criteria from MC simulation of the non-resonant- $\ell\ell$ background. The m_{T}^{ZZ} distributions for the non-resonant- $\ell\ell$ background are derived from the data CR and extrapolated to the SR. The total uncertainty in the non-resonant- $\ell\ell$ estimate is about 40%, including the statistical uncertainty of the data in the control region, the extrapolation and the method bias estimated from simulation. The m_{T}^{ZZ} distribution differences between data and simulation are taken as a shape uncertainty ($\sim 10\%$).

The $Z + \text{jets}$ background, expected to be $\sim 2\%$ of the total background, is estimated from a combination of MC and data-driven techniques. A $Z + \text{jets}$ enriched CR is defined by reversing the $E_{\text{T}}^{\text{miss}}/H_{\text{T}}$ selection. Additionally, the b -jet veto and the requirement on $\Delta\phi(\vec{p}_{\text{T}}^{\ell\ell}, \vec{E}_{\text{T}}^{\text{miss}})$ are removed to allow more data events. The estimation is performed by extrapolating the number of events observed in the CR, after subtracting non- Z -boson backgrounds, to the SR with a correction factor based on simulation. The m_{T}^{ZZ} distribution for the $Z + \text{jets}$ background is derived from simulation. The total uncertainty in the $Z + \text{jets}$ estimate is about 50% (80%) for the ee ($\mu\mu$) channel, including the statistical uncertainty of the data in the control region and the extrapolation factor. The shape difference in m_{T}^{ZZ} between $Z + \text{jets}$ MC events in the SR and those in the CR is taken into account as a systematic uncertainty.

Other backgrounds, such as triboson production, ttV , $W + \text{jets}$, and top quark processes other than pair production, constitute only a tiny fraction of the total background in the off-shell signal region, $< 1\%$, and are taken from simulation. The contribution from the on-shell Higgs production is negligible in the off-shell signal region.

The expected and observed numbers of events in the signal region for the $ZZ \rightarrow 2\ell 2\nu$ analysis are summarised in Table 1. Fig. 2 shows the observed and expected distributions of m_{T}^{ZZ} in both the ee and $\mu\mu$ channels in the off-shell region.

6. Systematic uncertainties

Systematic uncertainty sources impacting the analysis of both channels can be divided into two categories: uncertainties in the theoretical description of the signal and background processes and experimental uncertainties related to the detector or to the reconstruction algorithms. The largest systematic uncertainties arise

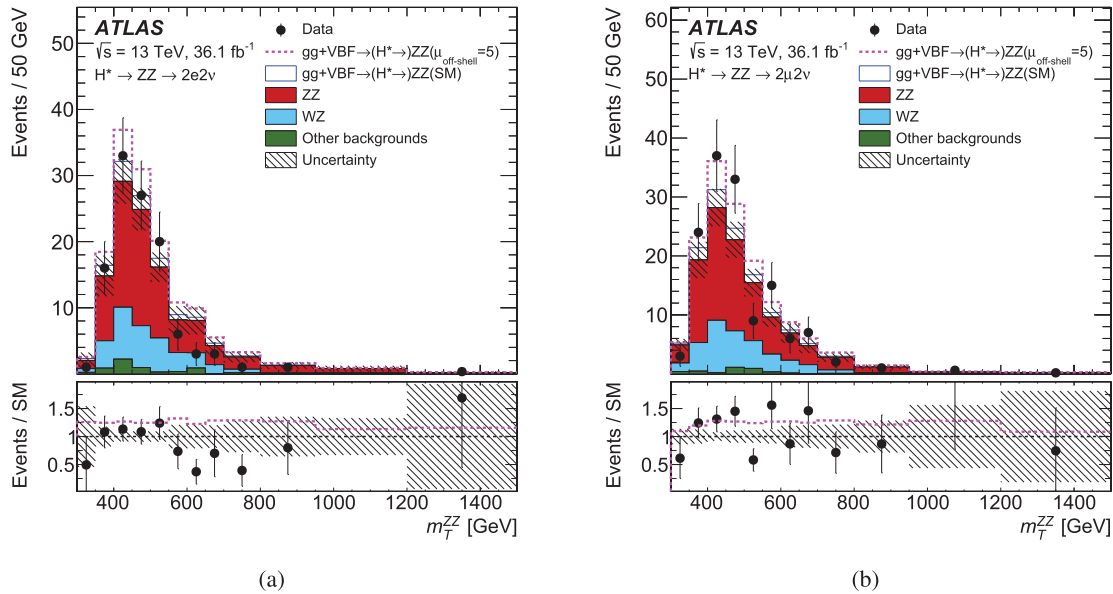


Fig. 2. Observed transverse mass m_T^{ZZ} distribution in the (a) ee channel and (b) $\mu\mu$ channel of the $ZZ \rightarrow 2\ell 2\nu$ off-shell region, compared to the expected contributions from the SM including the Higgs boson (stacked). Events with $1500 < m_T^{ZZ} < 2000$ GeV are included in the last bin of the distribution. The hatched area shows the combined statistical and systematic uncertainties. The dashed line corresponds to the total expected event yield, including all backgrounds and the Higgs boson with $\mu_{\text{off-shell}} = 5$. The ratio plot shows the observed data yield divided by the SM prediction (black points) as well as the total expected event yield with $\mu_{\text{off-shell}} = 5$ divided by the SM prediction (dashed line) in each bin.

from the theoretical uncertainties in the gg -initiated ZZ processes and the $q\bar{q} \rightarrow ZZ$ background process. The uncertainties from experimental measurements are generally small compared to the theoretical uncertainties in this analysis.

The theoretical uncertainties originate from the PDF choice, the missing higher-order corrections, and the parton-shower modelling. The PDF uncertainty corresponds to the 68% CL variations of the nominal PDF set NNPDF30NNLO for both $q\bar{q} \rightarrow ZZ$ and $gg \rightarrow (H^* \rightarrow)ZZ$, as well as the difference from alternative PDF sets. The alternative PDF sets used are CT10NNLO [68] and MMHT2014NNLO [69]. The uncertainty due to PDF is found to be about 3% in the high-mass region considered. The uncertainty due to higher-order QCD corrections (QCD scale uncertainty) is estimated by varying the renormalisation and factorisation scales independently, ranging from a factor of one-half to two. The uncertainty in the K -factors due to the NLO QCD scale uncertainty is 10–20% as a function of m_{ZZ} for the gg -initiated ZZ processes in the probed high-mass region, and ranges from 5% to 10% as a function of m_{ZZ} for the $q\bar{q} \rightarrow ZZ$ background. The QCD scale uncertainties are treated as correlated among the gg -initiated ZZ processes, and uncorrelated with the $q\bar{q}$ -initiated ZZ process. There are a few additional normalisation uncertainties associated with the NLO K -factors discussed in section 3. In the region below $2m_t$, the higher-order corrections are computed with a maximum jet transverse momentum of 150 GeV to ensure a good description by the $1/m_t$ expansion. The default scale uncertainty is therefore doubled for events which have a jet with $p_T > 150$ GeV, corresponding to about 8% of the events in this region. The scale uncertainty is also increased by 50% around the $2m_t$ threshold, with a Gaussian-smoothed transition decreasing to the default uncertainty within 50 GeV of the threshold. This is intended to allow for possible effects on the K -factor which have not been estimated as the top quark moves on-shell. It is assumed that the 10–20% NLO QCD scale uncertainty for the gg -initiated ZZ processes covers the assumption of massless loops above the $2m_t$ threshold, and as well the uncertainties in the 1.2 scale factor estimated only for the NNLO/NLO signal correction but also applied to the background and interference components. These NLO QCD scale uncertainties

are larger than those associated with the NNLO QCD signal uncertainties. The EW correction uncertainty for $q\bar{q} \rightarrow ZZ$ is evaluated using the same method as in Ref. [26] and its impact is estimated to be about 1%. The parton-shower uncertainty is evaluated by varying parameters in the parton-shower tunes according to Refs. [49,54] and found to be 2–3% in normalisation.

The theoretical uncertainties due to the missing higher-order corrections and PDF variations are small for VH -like and VBF-like processes $pp \rightarrow ZZ + 2j$; therefore, they are not included in the analysis.

For the $ZZ \rightarrow 4\ell$ analysis, the same sources of experimental uncertainty as in Ref. [32] are evaluated. The leading experimental systematic uncertainties are due to the electron and muon reconstruction and selection efficiency uncertainties, which are smaller than the uncertainties associated with the theoretical predictions.

Similarly, for the $ZZ \rightarrow 2\ell 2\nu$ channel, the same sources of experimental uncertainty as in Ref. [70] are evaluated. These experimental uncertainties affect the sensitivity of the $\mu_{\text{off-shell}}$ measurement only at the percent level.

The uncertainty in the combined 2015 and 2016 integrated luminosity is 2.1%, derived following a methodology similar to that detailed in Ref. [71], from a preliminary calibration of the luminosity scale using x - y beam-separation scans. This uncertainty is applied to the normalisation of the signal and also to background contributions whose normalisations are derived from MC simulations. A variation in the pile-up reweighting of MC events is included to cover the uncertainty in the ratio of the predicted and measured inelastic cross sections in Ref. [72].

7. Results

The results for the $ZZ \rightarrow 4\ell$ and $ZZ \rightarrow 2\ell 2\nu$ analyses are first translated into limits on the off-shell signal strength $\mu_{\text{off-shell}}$. A single off-shell signal-strength parameter is applied for all production modes, assuming that the ratio of the off-shell production rates via the ggF process to those via the VBF process are as predicted in the SM, namely $\mu_{\text{off-shell}}^{\text{ggF}}/\mu_{\text{off-shell}}^{\text{VBF}} = 1$. In a second step, the off-shell analyses are combined with the on-shell $ZZ^* \rightarrow$

Table 2

The 95% CL upper limits on $\mu_{\text{off-shell}}$, $\Gamma_H/\Gamma_H^{\text{SM}}$ and R_{gg} . Both the observed and expected limits are given. The 1σ (2σ) uncertainties represent 68% (95%) confidence intervals for the expected limit. The upper limits are evaluated using the CL_s method, with the SM values as the alternative hypothesis for each interpretation.

		Observed	Expected		
			Median	$\pm 1\sigma$	$\pm 2\sigma$
$\mu_{\text{off-shell}}$	$ZZ \rightarrow 4\ell$ analysis	4.5	4.3	[3.3, 5.4]	[2.7, 7.1]
	$ZZ \rightarrow 2\ell 2\nu$ analysis	5.3	4.4	[3.4, 5.5]	[2.8, 7.0]
	Combined	3.8	3.4	[2.7, 4.2]	[2.3, 5.3]
$\Gamma_H/\Gamma_H^{\text{SM}}$	Combined	3.5	3.7	[2.9, 4.8]	[2.4, 6.5]
R_{gg}	Combined	4.3	4.1	[3.3, 5.6]	[2.7, 8.2]

4ℓ [73] analysis, where the on-shell Higgs signal strength is measured to be $\mu_{\text{on-shell}} = 1.28_{-0.19}^{+0.21}$. The combination with the on-shell analysis is performed with two assumptions that correspond to different interpretations of the results. In the first combination, the parameter of interest is the ratio of off-shell to on-shell signal strengths, which can be interpreted as the Higgs boson width normalised to its SM prediction: $\mu_{\text{off-shell}}/\mu_{\text{on-shell}} = \Gamma_H/\Gamma_H^{\text{SM}}$. This interpretation assumes that the off- and on-shell coupling modifiers are the same for both ggF and VBF production modes (i.e., $\kappa_{g,\text{on-shell}} = \kappa_{g,\text{off-shell}} = \kappa_{V,\text{on-shell}} = \kappa_{V,\text{off-shell}}$). In the second combination, the parameter of interest is the ratio of off-shell to on-shell signal strengths for the ggF production only, $R_{gg} = \mu_{\text{off-shell}}^{\text{ggF}}/\mu_{\text{on-shell}}^{\text{ggF}}$, which can be interpreted as the ratio of off-shell to on-shell gluon couplings: $R_{gg} = \kappa_{g,\text{off-shell}}^2/\kappa_{g,\text{on-shell}}^2$. In this case the coupling scale factors $\kappa_V = \kappa_{V,\text{on-shell}} = \kappa_{V,\text{off-shell}}$ associated with on- and off-shell VBF production and the $H^{(*)} \rightarrow ZZ$ decay are assumed to be the same and fitted to the data (profiled). This also assumes that the total width is equal to the SM prediction.

The statistical analysis is based on the framework described in Refs. [74–76]. A binned likelihood function is constructed as a product of Poisson probability terms over all bins of the fit templates considered. This function depends on the parameter of interest μ , corresponding to one of the different interpretations discussed above ($\mu_{\text{off-shell}}$, $\Gamma_H/\Gamma_H^{\text{SM}}$ and R_{gg}), and θ , a set of nuisance parameters that encode the effects of systematic uncertainties on the signal and expected backgrounds, as described in Section 6. The nuisance parameters are constrained using either Gaussian or log-normal terms.

In the $ZZ \rightarrow 4\ell$ channel, a binned maximum-likelihood fit to the D_{ME} distribution is performed to extract the limits on μ . The fit model accounts for signal and background processes, including $gg \rightarrow (H^* \rightarrow)ZZ$, VBF ($H^* \rightarrow)ZZ$, $q\bar{q} \rightarrow ZZ$ and other backgrounds. The probability density functions of the signal-related processes $gg \rightarrow (H^* \rightarrow)ZZ$ and VBF ($H^* \rightarrow)ZZ$ are parameterised as a function of the off-shell Higgs boson signal strength $\mu_{\text{off-shell}}$ as given in Eqs. (1) and (2). In the $ZZ \rightarrow 2\ell 2\nu$ channel, a similar maximum-likelihood fit to the $m_{\tau\tau}^{ZZ}$ distribution is performed. The modelling of the dominant signal and background processes is the same as in the $ZZ \rightarrow 4\ell$ channel. The likelihood function for the combination of the $ZZ \rightarrow 4\ell$ and $ZZ \rightarrow 2\ell 2\nu$ channels is the product of the Poisson likelihoods of these individual channels. The main common theoretical and experimental systematic uncertainties are treated as correlated within different channels.

The PDF uncertainties and uncertainties from higher-order QCD corrections applied to the $q\bar{q} \rightarrow ZZ$ process are considered correlated between the on-shell and off-shell measurements. Given the different theoretical computations, the corresponding uncertainties are considered uncorrelated for the gg-initiated ZZ processes between the on-shell and off-shell measurements, and the impact of such a correlation effect is found to be small. In addition to the main theoretical uncertainties, the common experimental system-

atic uncertainties are treated as correlated between the on-shell and off-shell measurements.

Hypothesis testing and confidence intervals are based on the profile likelihood ratio [77]. The parameters of interest are different in the various tests, while the remaining parameters are profiled. All 95% CL upper limits are derived using the CL_s method [78], based on the ratio of one-sided p -values: $R_{\text{CL}_s}(\mu) = p_\mu/(1 - p_1)$ where p_μ is the p -value for testing a given $\mu = \mu_{\text{off-shell}}$ or $\mu = \Gamma_H/\Gamma_H^{\text{SM}}$ (the non-SM hypothesis) and p_1 is the p -value derived from the same test statistic under the SM hypothesis of $\mu_{\text{off-shell}} = 1$ in the first case and $\Gamma_H/\Gamma_H^{\text{SM}} = 1$ in the second case.²

The negative log-likelihood, $-2\ln\lambda$, is scanned as a function of a single parameter of interest, chosen to be $\mu_{\text{off-shell}}$, $\Gamma_H/\Gamma_H^{\text{SM}}$ or R_{gg} . The results are shown in Fig. 3 for observed and expected values. The results based on the CL_s method for the two individual analyses and their combination are reported in Table 2. As a result of the small data excess observed in the off-shell region, the observed limits on $\mu_{\text{off-shell}}$ are less stringent than the expected ones. The observed (expected) limit on $\Gamma_H/\Gamma_H^{\text{SM}}$ is 3.5 (3.7) at the 95% CL. Due to the fact that the measured on-shell signal strength $\mu_{\text{on-shell}}$ is larger than one [32], the observed limit on $\Gamma_H/\Gamma_H^{\text{SM}}$ is smaller than the expected limit. The limit on $\Gamma_H/\Gamma_H^{\text{SM}}$ can be translated into a limit on the total width of the Higgs boson, leading to an observed (expected) 95% CL upper limit on the Higgs boson total width of 14.4 (15.2) MeV.

These results are significantly improved compared to the Run-1 publication [26], the expected limit being about a factor two better.

If instead of constraining the $q\bar{q} \rightarrow ZZ$ background to the theoretical expectation, the normalisation is left as a free parameter in the profile likelihood fit, the upper limits on $\mu_{\text{on-shell}}$ are about 4% worse in the $ZZ \rightarrow 4\ell$ channel. If only the NLO K -factor are applied to the SM prediction of the gg-initiated ZZ processes, without the additional NNLO/NLO K -factor of 1.2 (Section 3), the upper limits on $\mu_{\text{off-shell}}$ and $\Gamma_H/\Gamma_H^{\text{SM}}$ are about 10% worse.

The impact of the various systematic uncertainties on the expected limit in the $\mu_{\text{off-shell}}$ fit are listed in Table 3. The values in this table were derived by fixing all the nuisance parameters associated with the systematic uncertainties to the values derived from the SM-conditional fit to the data, with the exception of the one under study. The uncertainties with the largest impact on the sensitivity of $\mu_{\text{off-shell}}$ are the theoretical uncertainties of the gg- and $q\bar{q}$ -initiated ZZ processes.

8. Conclusion

A determination of the off-shell Higgs boson signal strength in the $ZZ \rightarrow 4\ell$ and $ZZ \rightarrow 2\ell 2\nu$ final states and their combination is

² In the context of this analysis the alternative hypothesis is given by the SM value(s) for all relevant parameters of the fit model.

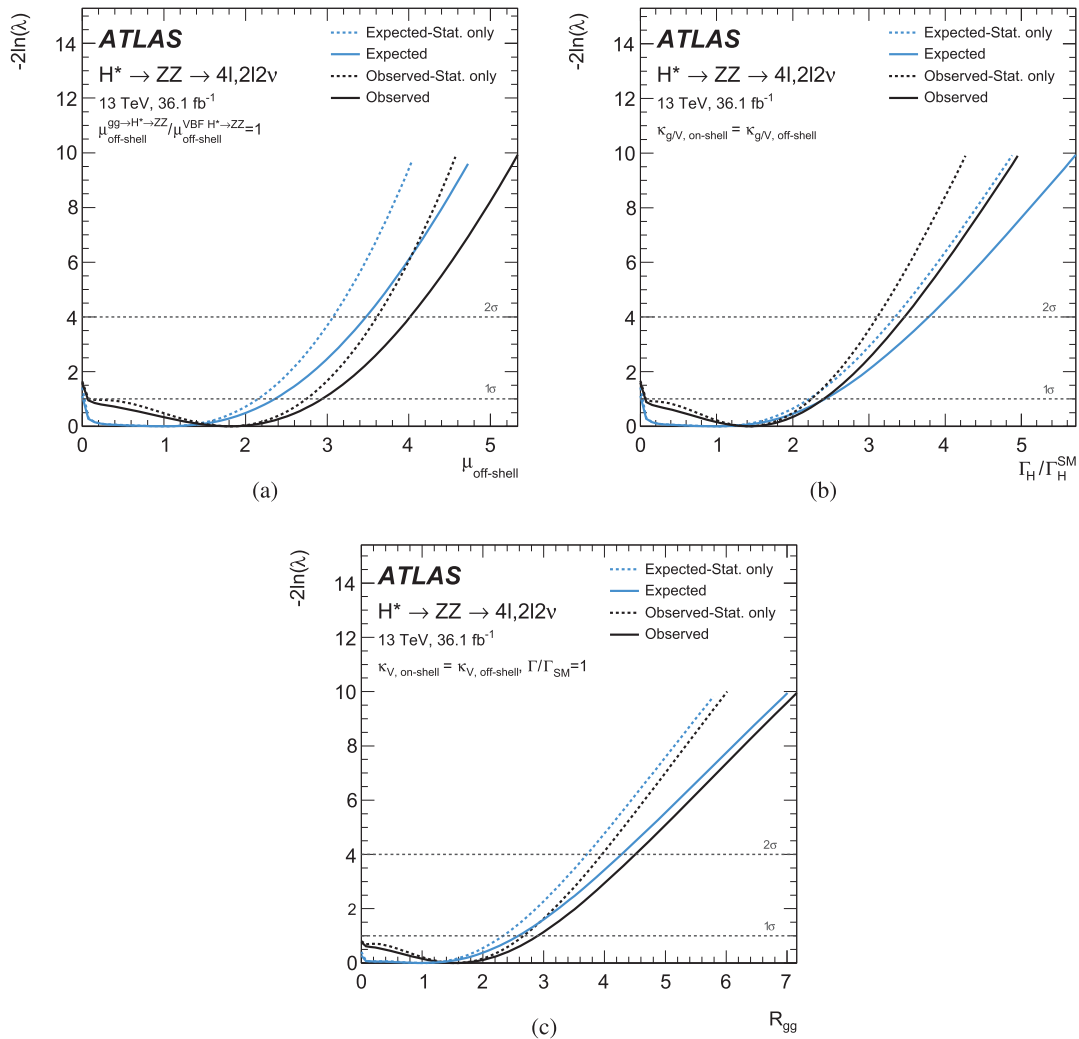


Fig. 3. Scan of the negative log-likelihood, $-2\ln\lambda$, for the (a) off-shell Higgs signal strength, $\mu_{\text{off-shell}}$ (b) $\Gamma_H/\Gamma_H^{\text{SM}}$ ratio (c) $R_{gg} = \kappa_{g,\text{off-shell}}^2/\kappa_{g,\text{on-shell}}^2$. The solid lower black (upper blue) line represents the observed (expected) value including all systematic uncertainties, while the dashed lower black (upper blue) line is for the observed (expected) value without systematic uncertainties (lower and upper refer here to the position of the lines in the legend). The double minimum structure of the scan when the parameter of interest approaches zero is the consequence of the parametrisation as shown in Eqs. (1). (For interpretation of the colours in the figure(s), the reader is referred to the web version of this article.)

Table 3

The expected 95% CL upper limit on $\mu_{\text{off-shell}}$ with a ranked listing of the impact of the leading systematic uncertainty individually, comparing with no systematic uncertainty or all systematic uncertainties. The upper limits are evaluated using the CL_s method.

Systematic uncertainty	95% CL upper limit on $\mu_{\text{off-shell}}$		
	$ZZ \rightarrow 4\ell$	$ZZ \rightarrow 2\ell 2\nu$	Combined
QCD scale $q\bar{q} \rightarrow ZZ$	4.2	3.9	3.2
QCD scale $gg \rightarrow (H^* \rightarrow)ZZ$	4.2	3.6	3.1
Luminosity	4.1	3.5	3.1
Remaining systematic uncertainties	4.1	3.5	3.0
All systematic uncertainties	4.3	4.4	3.4
No systematic uncertainties	4.0	3.4	3.0

presented. The result is based on pp collision data collected by the ATLAS experiment at the LHC, corresponding to an integrated luminosity of 36.1 fb^{-1} at a collision energy of $\sqrt{s} = 13 \text{ TeV}$. Using the CL_s method, the observed (expected) 95% confidence level (CL) upper limit on the off-shell signal strength is 3.8 (3.4). Assuming the ratio of the relevant Higgs boson couplings to the SM predictions are constant with energy from on-shell production to the

high-mass range considered in this analysis, a combination with the on-shell measurements yields an observed (expected) 95% CL upper limit on the Higgs boson total width of 14.4 (15.2) MeV.

Assuming that the total width of the Higgs boson is as expected in the SM, and the coupling scale factors associated with on- and off-shell VBF production and the $H^{(*)} \rightarrow ZZ$ decay are the same, the same combination can be interpreted as a limit on the ratio of the off-shell to the on-shell couplings to gluons $R_{gg} = \kappa_{g,\text{off-shell}}^2/\kappa_{g,\text{on-shell}}^2$. An observed (expected) limit of 4.3 (4.1) at 95% CL on R_{gg} is obtained.

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The ATLAS Collaboration

M. Aaboud^{34d}, G. Aad⁹⁹, B. Abbott¹²⁴, O. Abdinov^{13,*}, B. Abeloos¹²⁸, D.K. Abhayasinghe⁹¹, S.H. Abidi¹⁶⁴, O.S. AbouZeid³⁹, N.L. Abraham¹⁵³, H. Abramowicz¹⁵⁸, H. Abreu¹⁵⁷, Y. Abulaiti⁶, B.S. Acharya^{64a,64b,n}, S. Adachi¹⁶⁰, L. Adam⁹⁷, L. Adamczyk^{81a}, J. Adelman¹¹⁹, M. Adersberger¹¹², A. Adiguzel^{12c,ag}, T. Adye¹⁴¹, A.A. Affolder¹⁴³, Y. Afik¹⁵⁷, C. Agheorghiesei^{27c}, J.A. Aguilar-Saavedra^{136f,136a}, F. Ahmadov^{77,ae}, G. Aielli^{71a,71b}, S. Akatsuka⁸³, T.P.A. Åkesson⁹⁴, E. Akilli⁵², A.V. Akimov¹⁰⁸, G.L. Alberghi^{23b,23a}, J. Albert¹⁷³, P. Albicocco⁴⁹, M.J. Alconada Verzini⁸⁶, S. Alderweireldt¹¹⁷, M. Aleksa³⁵, I.N. Aleksandrov⁷⁷, C. Alexa^{27b}, T. Alexopoulos¹⁰, M. Alhroob¹²⁴, B. Ali¹³⁸, G. Alimonti^{66a}, J. Alison³⁶, S.P. Alkire¹⁴⁵, C. Allaire¹²⁸, B.M.M. Allbrooke¹⁵³, B.W. Allen¹²⁷, P.P. Allport²¹, A. Aloisio^{67a,67b}, A. Alonso³⁹, F. Alonso⁸⁶, C. Alpigiani¹⁴⁵, A.A. Alshehri⁵⁵, M.I. Alstaty⁹⁹, B. Alvarez Gonzalez³⁵, D. Álvarez Piqueras¹⁷¹, M.G. Alvigi^{67a,67b}, B.T. Amadio¹⁸, Y. Amaral Coutinho^{78b}, A. Ambler¹⁰¹, L. Ambroz¹³¹, C. Amelung²⁶, D. Amidei¹⁰³, S.P. Amor Dos Santos^{136a,136c}, S. Amoroso⁴⁴, C.S. Amrouche⁵², C. Anastopoulos¹⁴⁶, L.S. Ancu⁵², N. Andari¹⁴², T. Andeen¹¹, C.F. Anders^{59b}, J.K. Anders²⁰, K.J. Anderson³⁶, A. Andreazza^{66a,66b}, V. Andrei^{59a}, C.R. Anelli¹⁷³, S. Angelidakis³⁷, I. Angelozzi¹¹⁸, A. Angerami³⁸, A.V. Anisenkov^{120b,120a}, A. Annovi^{69a}, C. Antel^{59a}, M.T. Anthony¹⁴⁶, M. Antonelli⁴⁹, D.J.A. Antrim¹⁶⁸, F. Anulli^{70a}, M. Aoki⁷⁹, J.A. Aparisi Pozo¹⁷¹, L. Aperio Bella³⁵, G. Arabidze¹⁰⁴, J.P. Araque^{136a}, V. Araujo Ferraz^{78b}, R. Araujo Pereira^{78b}, A.T.H. Arce⁴⁷, R.E. Ardell⁹¹, F.A. Arduh⁸⁶, J.-F. Arguin¹⁰⁷, S. Argyropoulos⁷⁵, A.J. Armbruster³⁵, L.J. Armitage⁹⁰, A. Armstrong¹⁶⁸, O. Arnaez¹⁶⁴, H. Arnold¹¹⁸, M. Arratia³¹, O. Arslan²⁴, A. Artamonov^{109,*}, G. Artoni¹³¹, S. Artz⁹⁷, S. Asai¹⁶⁰, N. Asbah⁵⁷, E.M. Asimakopoulou¹⁶⁹, L. Asquith¹⁵³, K. Assamagan²⁹, R. Astalos^{28a}, R.J. Atkin^{32a}, M. Atkinson¹⁷⁰, N.B. Atlay¹⁴⁸, K. Augsten¹³⁸,

G. Avolio³⁵, R. Avramidou^{58a}, M.K. Ayoub^{15a}, G. Azuelos^{107,ar}, A.E. Baas^{59a}, M.J. Baca²¹,
H. Bachacou¹⁴², K. Bachas^{65a,65b}, M. Backes¹³¹, P. Bagnaia^{70a,70b}, M. Bahmani⁸², H. Bahrasemani¹⁴⁹,
A.J. Bailey¹⁷¹, J.T. Baines¹⁴¹, M. Bajic³⁹, C. Bakalis¹⁰, O.K. Baker¹⁸⁰, P.J. Bakker¹¹⁸, D. Bakshi Gupta⁹³,
S. Balaji¹⁵⁴, E.M. Baldin^{120b,120a}, P. Balek¹⁷⁷, F. Balli¹⁴², W.K. Balunas¹³³, J. Balz⁹⁷, E. Banas⁸²,
A. Bandyopadhyay²⁴, S. Banerjee^{178,j}, A.A.E. Bannoura¹⁷⁹, L. Barak¹⁵⁸, W.M. Barbe³⁷, E.L. Barberio¹⁰²,
D. Barberis^{53b,53a}, M. Barbero⁹⁹, T. Barillari¹¹³, M.-S. Barisits³⁵, J. Barkeloo¹²⁷, T. Barklow¹⁵⁰,
R. Barnea¹⁵⁷, S.L. Barnes^{58c}, B.M. Barnett¹⁴¹, R.M. Barnett¹⁸, Z. Barnovska-Blenessy^{58a}, A. Baroncelli^{72a},
G. Barone²⁶, A.J. Barr¹³¹, L. Barranco Navarro¹⁷¹, F. Barreiro⁹⁶, J. Barreiro Guimarães da Costa^{15a},
R. Bartoldus¹⁵⁰, A.E. Barton⁸⁷, P. Bartos^{28a}, A. Basalae¹³⁴, A. Bassalat¹²⁸, R.L. Bates⁵⁵, S.J. Batista¹⁶⁴,
S. Batlamous^{34e}, J.R. Batley³¹, M. Battaglia¹⁴³, M. Baucé^{70a,70b}, F. Bauer¹⁴², K.T. Bauer¹⁶⁸,
H.S. Bawa^{150,l}, J.B. Beacham¹²², T. Beau¹³², P.H. Beauchemin¹⁶⁷, P. Bechtel²⁴, H.C. Beck⁵¹, H.P. Beck^{20,q},
K. Becker⁵⁰, M. Becker⁹⁷, C. Becot⁴⁴, A. Beddall^{12d}, A.J. Beddall^{12a}, V.A. Bednyakov⁷⁷, M. Bedognetti¹¹⁸,
C.P. Bee¹⁵², T.A. Beermann³⁵, M. Begalli^{78b}, M. Begel²⁹, A. Behera¹⁵², J.K. Behr⁴⁴, A.S. Bell⁹²,
G. Bella¹⁵⁸, L. Bellagamba^{23b}, A. Bellerive³³, M. Bellomo¹⁵⁷, P. Bellos⁹, K. Belotskiy¹¹⁰, N.L. Belyaev¹¹⁰,
O. Benary^{158,*}, D. Benchekroun^{34a}, M. Bender¹¹², N. Benekos¹⁰, Y. Benhammou¹⁵⁸,
E. Benhar Noccioli¹⁸⁰, J. Benitez⁷⁵, D.P. Benjamin⁴⁷, M. Benoit⁵², J.R. Bensinger²⁶, S. Bentvelsen¹¹⁸,
L. Beresford¹³¹, M. Beretta⁴⁹, D. Berge⁴⁴, E. Bergeaas Kuutmann¹⁶⁹, N. Berger⁵, L.J. Bergsten²⁶,
J. Beringer¹⁸, S. Berlendis⁷, N.R. Bernard¹⁰⁰, G. Bernardi¹³², C. Bernius¹⁵⁰, F.U. Bernlochner²⁴,
T. Berry⁹¹, P. Berta⁹⁷, C. Bertella^{15a}, G. Bertoli^{43a,43b}, I.A. Bertram⁸⁷, G.J. Besjes³⁹,
O. Bessidskaia Bylund¹⁷⁹, M. Bessner⁴⁴, N. Besson¹⁴², A. Bethani⁹⁸, S. Bethke¹¹³, A. Betti²⁴,
A.J. Bevan⁹⁰, J. Beyer¹¹³, R.M. Bianchi¹³⁵, O. Biebel¹¹², D. Biedermann¹⁹, R. Bielski³⁵, K. Bierwagen⁹⁷,
N.V. Biesuz^{69a,69b}, M. Biglietti^{72a}, T.R.V. Billoud¹⁰⁷, M. Bindi⁵¹, A. Bingul^{12d}, C. Bini^{70a,70b},
S. Biondi^{23b,23a}, M. Birman¹⁷⁷, T. Bisanz⁵¹, J.P. Biswal¹⁵⁸, C. Bittrich⁴⁶, D.M. Bjergaard⁴⁷, J.E. Black¹⁵⁰,
K.M. Black²⁵, T. Blazek^{28a}, I. Bloch⁴⁴, C. Blocker²⁶, A. Blue⁵⁵, U. Blumenschein⁹⁰, Dr. Blunier^{144a},
G.J. Bobbink¹¹⁸, V.S. Bobrovnikov^{120b,120a}, S.S. Bocchetta⁹⁴, A. Bocci⁴⁷, D. Boerner¹⁷⁹, D. Bogavac¹¹²,
A.G. Bogdanchikov^{120b,120a}, C. Bohm^{43a}, V. Boisvert⁹¹, P. Bokan^{169,x}, T. Bold^{81a}, A.S. Boldyrev¹¹¹,
A.E. Bolz^{59b}, M. Bomben¹³², M. Bona⁹⁰, J.S. Bonilla¹²⁷, M. Boonekamp¹⁴², A. Borisov¹⁴⁰, G. Borissov⁸⁷,
J. Bortfeldt³⁵, D. Bortoletto¹³¹, V. Bortolotto^{71a,71b}, D. Boscherini^{23b}, M. Bosman¹⁴, J.D. Bossio Sola³⁰,
K. Bouaouda^{34a}, J. Boudreau¹³⁵, E.V. Bouhova-Thacker⁸⁷, D. Boumediene³⁷, C. Bourdarios¹²⁸,
S.K. Boutle⁵⁵, A. Boveia¹²², J. Boyd³⁵, D. Boye^{32b}, I.R. Boyko⁷⁷, A.J. Bozson⁹¹, J. Bracinik²¹,
N. Brahimi⁹⁹, A. Brandt⁸, G. Brandt¹⁷⁹, O. Brandt^{59a}, F. Braren⁴⁴, U. Bratzler¹⁶¹, B. Brau¹⁰⁰, J.E. Brau¹²⁷,
W.D. Breaden Madden⁵⁵, K. Brendlinger⁴⁴, L. Brenner⁴⁴, R. Brenner¹⁶⁹, S. Bressler¹⁷⁷, B. Brickwedde⁹⁷,
D.L. Briglin²¹, D. Britton⁵⁵, D. Britzger^{59b}, I. Brock²⁴, R. Brock¹⁰⁴, G. Brooijmans³⁸, T. Brooks⁹¹,
W.K. Brooks^{144b}, E. Brost¹¹⁹, J.H. Broughton²¹, P.A. Bruckman de Renstrom⁸², D. Bruncko^{28b},
A. Bruni^{23b}, G. Bruni^{23b}, L.S. Bruni¹¹⁸, S. Bruno^{71a,71b}, B.H. Brunt³¹, M. Bruschi^{23b}, N. Bruscino¹³⁵,
P. Bryant³⁶, L. Bryngemark⁴⁴, T. Buanes¹⁷, Q. Buat³⁵, P. Buchholz¹⁴⁸, A.G. Buckley⁵⁵, I.A. Budagov⁷⁷,
F. Buehrer⁵⁰, M.K. Bugge¹³⁰, O. Bulekov¹¹⁰, D. Bullock⁸, T.J. Burch¹¹⁹, S. Burdin⁸⁸, C.D. Burgard¹¹⁸,
A.M. Burger⁵, B. Burghgrave¹¹⁹, K. Burka⁸², S. Burke¹⁴¹, I. Burmeister⁴⁵, J.T.P. Burr¹³¹, V. Büscher⁹⁷,
E. Buschmann⁵¹, P. Bussey⁵⁵, J.M. Butler²⁵, C.M. Buttar⁵⁵, J.M. Butterworth⁹², P. Butti³⁵,
W. Buttinger³⁵, A. Buzatu¹⁵⁵, A.R. Buzykaev^{120b,120a}, G. Cabras^{23b,23a}, S. Cabrera Urbán¹⁷¹,
D. Caforio¹³⁸, H. Cai¹⁷⁰, V.M.M. Cairo², O. Cakir^{4a}, N. Calace⁵², P. Calafiura¹⁸, A. Calandri⁹⁹,
G. Calderini¹³², P. Calfayan⁶³, G. Callea^{40b,40a}, L.P. Caloba^{78b}, S. Calvente Lopez⁹⁶, D. Calvet³⁷,
S. Calvet³⁷, T.P. Calvet¹⁵², M. Calvetti^{69a,69b}, R. Camacho Toro¹³², S. Camarda³⁵, P. Camarri^{71a,71b},
D. Cameron¹³⁰, R. Caminal Armadans¹⁰⁰, C. Camincher³⁵, S. Campana³⁵, M. Campanelli⁹²,
A. Camplani³⁹, A. Campoverde¹⁴⁸, V. Canale^{67a,67b}, M. Cano Bret^{58c}, J. Cantero¹²⁵, T. Cao¹⁵⁸, Y. Cao¹⁷⁰,
M.D.M. Capeans Garrido³⁵, I. Caprini^{27b}, M. Caprini^{27b}, M. Capua^{40b,40a}, R.M. Carbone³⁸,
R. Cardarelli^{71a}, F.C. Cardillo¹⁴⁶, I. Carli¹³⁹, T. Carli³⁵, G. Carlino^{67a}, B.T. Carlson¹³⁵, L. Carminati^{66a,66b},
R.M.D. Carney^{43a,43b}, S. Caron¹¹⁷, E. Carquin^{144b}, S. Carrá^{66a,66b}, G.D. Carrillo-Montoya³⁵, D. Casadei^{32b},
M.P. Casado^{14,f}, A.F. Casha¹⁶⁴, D.W. Casper¹⁶⁸, R. Castelijm¹¹⁸, F.L. Castillo¹⁷¹, V. Castillo Gimenez¹⁷¹,
N.F. Castro^{136a,136e}, A. Catinaccio³⁵, J.R. Catmore¹³⁰, A. Cattai³⁵, J. Caudron²⁴, V. Cavaliere²⁹,
E. Cavallaro¹⁴, D. Cavalli^{66a}, M. Cavalli-Sforza¹⁴, V. Cavasinni^{69a,69b}, E. Celebi^{12b}, F. Ceradini^{72a,72b},
L. Cerda Alberich¹⁷¹, A.S. Cerqueira^{78a}, A. Cerri¹⁵³, L. Cerrito^{71a,71b}, F. Cerutti¹⁸, A. Cervelli^{23b,23a},

S.A. Cetin ^{12b}, A. Chafaq ^{34a}, D Chakraborty ¹¹⁹, S.K. Chan ⁵⁷, W.S. Chan ¹¹⁸, Y.L. Chan ^{61a}, J.D. Chapman ³¹, B. Chargeishvili ^{156b}, D.G. Charlton ²¹, C.C. Chau ³³, C.A. Chavez Barajas ¹⁵³, S. Che ¹²², A. Chegwidan ¹⁰⁴, S. Chekanov ⁶, S.V. Chekulaev ^{165a}, G.A. Chelkov ^{77.aq}, M.A. Chelstowska ³⁵, C. Chen ^{58a}, C.H. Chen ⁷⁶, H. Chen ²⁹, J. Chen ^{58a}, J. Chen ³⁸, S. Chen ¹³³, S.J. Chen ^{15c}, X. Chen ^{15b.ap}, Y. Chen ⁸⁰, Y-H. Chen ⁴⁴, H.C. Cheng ¹⁰³, H.J. Cheng ^{15d}, A. Cheplakov ⁷⁷, E. Cheremushkina ¹⁴⁰, R. Cherkaoui El Moursli ^{34e}, E. Cheu ⁷, K. Cheung ⁶², L. Chevalier ¹⁴², V. Chiarella ⁴⁹, G. Chiarelli ^{69a}, G. Chiodini ^{65a}, A.S. Chisholm ^{35.21}, A. Chitan ^{27b}, I. Chiu ¹⁶⁰, Y.H. Chiu ¹⁷³, M.V. Chizhov ⁷⁷, K. Choi ⁶³, A.R. Chomont ¹²⁸, S. Chouridou ¹⁵⁹, Y.S. Chow ¹¹⁸, V. Christodoulou ⁹², M.C. Chu ^{61a}, J. Chudoba ¹³⁷, A.J. Chuinard ¹⁰¹, J.J. Chwastowski ⁸², L. Chytka ¹²⁶, D. Cinca ⁴⁵, V. Cindro ⁸⁹, I.A. Cioară ²⁴, A. Ciocio ¹⁸, F. Ciotto ^{67a.67b}, Z.H. Citron ¹⁷⁷, M. Citterio ^{66a}, A. Clark ⁵², M.R. Clark ³⁸, P.J. Clark ⁴⁸, C. Clement ^{43a.43b}, Y. Coadou ⁹⁹, M. Cobl ^{64a.64c}, A. Coccaro ^{53b.53a}, J. Cochran ⁷⁶, H. Cohen ¹⁵⁸, A.E.C. Coimbra ¹⁷⁷, L. Colasurdo ¹¹⁷, B. Cole ³⁸, A.P. Colijn ¹¹⁸, J. Collot ⁵⁶, P. Conde Muiño ^{136a.136b}, E. Coniavitis ⁵⁰, S.H. Connell ^{32b}, I.A. Connelly ⁹⁸, S. Constantinescu ^{27b}, F. Conventi ^{67a.as}, A.M. Cooper-Sarkar ¹³¹, F. Cormier ¹⁷², K.J.R. Cormier ¹⁶⁴, L.D. Corpe ⁹², M. Corradi ^{70a.70b}, E.E. Corrigan ⁹⁴, F. Corriveau ^{101.ac}, A. Cortes-Gonzalez ³⁵, M.J. Costa ¹⁷¹, F. Costanza ⁵, D. Costanzo ¹⁴⁶, G. Cottin ³¹, G. Cowan ⁹¹, B.E. Cox ⁹⁸, J. Crane ⁹⁸, K. Cranmer ¹²¹, S.J. Crawley ⁵⁵, R.A. Creager ¹³³, G. Cree ³³, S. Crépe-Renaudin ⁵⁶, F. Crescioli ¹³², M. Cristinziani ²⁴, V. Croft ¹²¹, G. Crosetti ^{40b.40a}, A. Cueto ⁹⁶, T. Cuhadar Donszelmann ¹⁴⁶, A.R. Cukierman ¹⁵⁰, S. Czekierda ⁸², P. Czodrowski ³⁵, M.J. Da Cunha Sargedas De Sousa ^{58b.136b}, C. Da Via ⁹⁸, W. Dabrowski ^{81a}, T. Dado ^{28a.x}, S. Dahbi ^{34e}, T. Dai ¹⁰³, F. Dallaire ¹⁰⁷, C. Dallapiccola ¹⁰⁰, M. Dam ³⁹, G. D'amen ^{23b.23a}, J. Damp ⁹⁷, J.R. Dandoy ¹³³, M.F. Daneri ³⁰, N.P. Dang ^{178.j}, N.D. Dann ⁹⁸, M. Danninger ¹⁷², V. Dao ³⁵, G. Darbo ^{53b}, S. Darmora ⁸, O. Dartsis ⁵, A. Dattagupta ¹²⁷, T. Daubney ⁴⁴, S. D'Auria ⁵⁵, W. Davey ²⁴, C. David ⁴⁴, T. Davidek ¹³⁹, D.R. Davis ⁴⁷, E. Dawe ¹⁰², I. Dawson ¹⁴⁶, K. De ⁸, R. De Asmundis ^{67a}, A. De Benedetti ¹²⁴, M. De Beurs ¹¹⁸, S. De Castro ^{23b.23a}, S. De Cecco ^{70a.70b}, N. De Groot ¹¹⁷, P. de Jong ¹¹⁸, H. De la Torre ¹⁰⁴, F. De Lorenzi ⁷⁶, A. De Maria ^{51.s}, D. De Pedis ^{70a}, A. De Salvo ^{70a}, U. De Sanctis ^{71a.71b}, M. De Santis ^{71a.71b}, A. De Santo ¹⁵³, K. De Vasconcelos Corga ⁹⁹, J.B. De Vivie De Regie ¹²⁸, C. Debenedetti ¹⁴³, D.V. Dedovich ⁷⁷, N. Dehghanian ³, M. Del Gaudio ^{40b.40a}, J. Del Peso ⁹⁶, Y. Delabat Diaz ⁴⁴, D. Delgove ¹²⁸, F. Deliot ¹⁴², C.M. Delitzsch ⁷, M. Della Pietra ^{67a.67b}, D. Della Volpe ⁵², A. Dell'Acqua ³⁵, L. Dell'Asta ²⁵, M. Delmastro ⁵, C. Delporte ¹²⁸, P.A. Delsart ⁵⁶, D.A. DeMarco ¹⁶⁴, S. Demers ¹⁸⁰, M. Demichev ⁷⁷, S.P. Denisov ¹⁴⁰, D. Denysiuk ¹¹⁸, L. D'Eramo ¹³², D. Derendarz ⁸², J.E. Derkaoui ^{34d}, F. Derue ¹³², P. Dervan ⁸⁸, K. Desch ²⁴, C. Deterre ⁴⁴, K. Dette ¹⁶⁴, M.R. Devesa ³⁰, P.O. Deviveiros ³⁵, A. Dewhurst ¹⁴¹, S. Dhaliwal ²⁶, F.A. Di Bello ⁵², A. Di Ciaccio ^{71a.71b}, L. Di Ciaccio ⁵, W.K. Di Clemente ¹³³, C. Di Donato ^{67a.67b}, A. Di Girolamo ³⁵, G. Di Gregorio ^{69a.69b}, B. Di Micco ^{72a.72b}, R. Di Nardo ¹⁰⁰, K.F. Di Petrillo ⁵⁷, R. Di Sipio ¹⁶⁴, D. Di Valentino ³³, C. Diaconu ⁹⁹, M. Diamond ¹⁶⁴, F.A. Dias ³⁹, T. Dias Do Vale ^{136a}, M.A. Diaz ^{144a}, J. Dickinson ¹⁸, E.B. Diehl ¹⁰³, J. Dietrich ¹⁹, S. Díez Cornell ⁴⁴, A. Dimitrievska ¹⁸, J. Dingfelder ²⁴, F. Dittus ³⁵, F. Djama ⁹⁹, T. Djobava ^{156b}, J.I. Djuvsland ^{59a}, M.A.B. Do Vale ^{78c}, M. Dobre ^{27b}, D. Dodsworth ²⁶, C. Doglioni ⁹⁴, J. Dolejsi ¹³⁹, Z. Dolezal ¹³⁹, M. Donadelli ^{78d}, J. Donini ³⁷, A. D'Onofrio ⁹⁰, M. D'Onofrio ⁸⁸, J. Dopke ¹⁴¹, A. Doria ^{67a}, M.T. Dova ⁸⁶, A.T. Doyle ⁵⁵, E. Drechsler ⁵¹, E. Dreyer ¹⁴⁹, T. Dreyer ⁵¹, Y. Du ^{58b}, F. Dubinin ¹⁰⁸, M. Dubovsky ^{28a}, A. Dubreuil ⁵², E. Duchovni ¹⁷⁷, G. Duckeck ¹¹², A. Ducourthial ¹³², O.A. Ducu ^{107.w}, D. Duda ¹¹³, A. Dudarev ³⁵, A.C. Dudder ⁹⁷, E.M. Duffield ¹⁸, L. Duflot ¹²⁸, M. Dührssen ³⁵, C. Dülsen ¹⁷⁹, M. Dumancic ¹⁷⁷, A.E. Dumitriu ^{27b.d}, A.K. Duncan ⁵⁵, M. Dunford ^{59a}, A. Duperrin ⁹⁹, H. Duran Yildiz ^{4a}, M. Düren ⁵⁴, A. Durglishvili ^{156b}, D. Duschinger ⁴⁶, B. Dutta ⁴⁴, D. Duvnjak ¹, M. Dyndal ⁴⁴, S. Dysch ⁹⁸, B.S. Dziedzic ⁸², C. Eckardt ⁴⁴, K.M. Ecker ¹¹³, R.C. Edgar ¹⁰³, T. Eifert ³⁵, G. Eigen ¹⁷, K. Einsweiler ¹⁸, T. Ekelof ¹⁶⁹, M. El Kacimi ^{34c}, R. El Kosseifi ⁹⁹, V. Ellajosyula ⁹⁹, M. Ellert ¹⁶⁹, F. Ellinghaus ¹⁷⁹, A.A. Elliot ⁹⁰, N. Ellis ³⁵, J. Elmsheuser ²⁹, M. Elsing ³⁵, D. Emelianov ¹⁴¹, Y. Enari ¹⁶⁰, J.S. Ennis ¹⁷⁵, M.B. Epland ⁴⁷, J. Erdmann ⁴⁵, A. Ereditato ²⁰, S. Errede ¹⁷⁰, M. Escalier ¹²⁸, C. Escobar ¹⁷¹, O. Estrada Pastor ¹⁷¹, A.I. Etienne ¹⁴², E. Etzion ¹⁵⁸, H. Evans ⁶³, A. Ezhilov ¹³⁴, M. Ezzi ^{34e}, F. Fabbri ⁵⁵, L. Fabbri ^{23b.23a}, V. Fabiani ¹¹⁷, G. Facini ⁹², R.M. Faisca Rodrigues Pereira ^{136a}, R.M. Fakhruddinov ¹⁴⁰, S. Falciano ^{70a}, P.J. Falke ⁵, S. Falke ⁵, J. Faltova ¹³⁹, Y. Fang ^{15a}, M. Fanti ^{66a.66b}, A. Farbin ⁸, A. Farilla ^{72a}, E.M. Farina ^{68a.68b}, T. Farooque ¹⁰⁴, S. Farrell ¹⁸, S.M. Farrington ¹⁷⁵, P. Farthouat ³⁵, F. Fassi ^{34e}, P. Fassnacht ³⁵, D. Fassouliotis ⁹, M. Fauci Giannelli ⁴⁸, A. Favareto ^{53b.53a}, W.J. Fawcett ³¹, L. Fayard ¹²⁸, O.L. Fedin ^{134.o}, W. Fedorko ¹⁷², M. Feickert ⁴¹, S. Feigl ¹³⁰, L. Feligioni ⁹⁹, C. Feng ^{58b}, E.J. Feng ³⁵,

M. Feng⁴⁷, M.J. Fenton⁵⁵, A.B. Fenyuk¹⁴⁰, L. Feremenga⁸, J. Ferrando⁴⁴, A. Ferrari¹⁶⁹, P. Ferrari¹¹⁸,
 R. Ferrari^{68a}, D.E. Ferreira de Lima^{59b}, A. Ferrer¹⁷¹, D. Ferrere⁵², C. Ferretti¹⁰³, F. Fiedler⁹⁷, A. Filipčić⁸⁹,
 F. Filthaut¹¹⁷, K.D. Finelli²⁵, M.C.N. Fiolhais^{136a,136c,a}, L. Fiorini¹⁷¹, C. Fischer¹⁴, W.C. Fisher¹⁰⁴,
 N. Flaschel⁴⁴, I. Fleck¹⁴⁸, P. Fleischmann¹⁰³, R.R.M. Fletcher¹³³, T. Flick¹⁷⁹, B.M. Flierl¹¹², L.M. Flores¹³³,
 L.R. Flores Castillo^{61a}, F.M. Follega^{73a,73b}, N. Fomin¹⁷, G.T. Forcolin^{73a,73b}, A. Formica¹⁴², F.A. Förster¹⁴,
 A.C. Forti⁹⁸, A.G. Foster²¹, D. Fournier¹²⁸, H. Fox⁸⁷, S. Fracchia¹⁴⁶, P. Francavilla^{69a,69b},
 M. Franchini^{23b,23a}, S. Franchino^{59a}, D. Francis³⁵, L. Franconi¹⁴³, M. Franklin⁵⁷, M. Frate¹⁶⁸,
 M. Fraternali^{68a,68b}, A.N. Fray⁹⁰, D. Freeborn⁹², S.M. Fressard-Batraneanu³⁵, B. Freund¹⁰⁷,
 W.S. Freund^{78b}, E.M. Freundlich⁴⁵, D.C. Frizzell¹²⁴, D. Froidevaux³⁵, J.A. Frost¹³¹, C. Fukunaga¹⁶¹,
 E. Fullana Torregrosa¹⁷¹, T. Fusayasu¹¹⁴, J. Fuster¹⁷¹, O. Gabizon¹⁵⁷, A. Gabrielli^{23b,23a}, A. Gabrielli¹⁸,
 G.P. Gach^{81a}, S. Gadatsch⁵², P. Gadow¹¹³, G. Gagliardi^{53b,53a}, L.G. Gagnon¹⁰⁷, C. Galea^{27b},
 B. Galhardo^{136a,136c}, E.J. Gallas¹³¹, B.J. Gallop¹⁴¹, P. Gallus¹³⁸, G. Galster³⁹, R. Gamboa Goni⁹⁰,
 K.K. Gan¹²², S. Ganguly¹⁷⁷, J. Gao^{58a}, Y. Gao⁸⁸, Y.S. Gao^{150,l}, C. García¹⁷¹, J.E. García Navarro¹⁷¹,
 J.A. García Pascual^{15a}, M. Garcia-Sciveres¹⁸, R.W. Gardner³⁶, N. Garelli¹⁵⁰, V. Garonne¹³⁰,
 K. Gasnikova⁴⁴, A. Gaudiello^{53b,53a}, G. Gaudio^{68a}, I.L. Gavrilenko¹⁰⁸, A. Gavrilyuk¹⁰⁹, C. Gay¹⁷²,
 G. Gaycken²⁴, E.N. Gazis¹⁰, C.N.P. Gee¹⁴¹, J. Geisen⁵¹, M. Geisen⁹⁷, M.P. Geisler^{59a}, K. Gellerstedt^{43a,43b},
 C. Gemme^{53b}, M.H. Genest⁵⁶, C. Geng¹⁰³, S. Gentile^{70a,70b}, S. George⁹¹, D. Gerbaudo¹⁴, G. Gessner⁴⁵,
 S. Ghasemi¹⁴⁸, M. Ghasemi Bostanabad¹⁷³, M. Ghneimat²⁴, B. Giacobbe^{23b}, S. Giagu^{70a,70b},
 N. Giangiacomi^{23b,23a}, P. Giannetti^{69a}, A. Giannini^{67a,67b}, S.M. Gibson⁹¹, M. Gignac¹⁴³, D. Gillberg³³,
 G. Gilles¹⁷⁹, D.M. Gingrich^{3,ar}, M.P. Giordani^{64a,64c}, F.M. Giorgi^{23b}, P.F. Giraud¹⁴², P. Giromini⁵⁷,
 G. Giugliarelli^{64a,64c}, D. Giugni^{66a}, F. Giuli¹³¹, M. Giulini^{59b}, S. Gkaitatzis¹⁵⁹, I. Gkialas^{9,i},
 E.L. Gkoukousis¹⁴, P. Gkoutoumis¹⁰, L.K. Gladilin¹¹¹, C. Glasman⁹⁶, J. Glatzer¹⁴, P.C.F. Glaysheer⁴⁴,
 A. Glazov⁴⁴, M. Goblirsch-Kolb²⁶, J. Godlewski⁸², S. Goldfarb¹⁰², T. Golling⁵², D. Golubkov¹⁴⁰,
 A. Gomes^{136a,136b,136d}, R. Goncalves Gama^{78a}, R. Gonçalo^{136a}, G. Gonella⁵⁰, L. Gonella²¹,
 A. Gongadze⁷⁷, F. Gonnella²¹, J.L. Gonski⁵⁷, S. González de la Hoz¹⁷¹, S. Gonzalez-Sevilla⁵²,
 L. Goossens³⁵, P.A. Gorbounov¹⁰⁹, H.A. Gordon²⁹, B. Gorini³⁵, E. Gorini^{65a,65b}, A. Gorišek⁸⁹,
 A.T. Goshaw⁴⁷, C. Gössling⁴⁵, M.I. Gostkin⁷⁷, C.A. Gottardo²⁴, C.R. Goudet¹²⁸, D. Goujdami^{34c},
 A.G. Goussiou¹⁴⁵, N. Govender^{32b,b}, C. Goy⁵, E. Gozani¹⁵⁷, I. Grabowska-Bold^{81a}, P.O.J. Gradin¹⁶⁹,
 E.C. Graham⁸⁸, J. Gramling¹⁶⁸, E. Gramstad¹³⁰, S. Grancagnolo¹⁹, V. Gratchev¹³⁴, P.M. Gravila^{27f},
 F.G. Gravili^{65a,65b}, C. Gray⁵⁵, H.M. Gray¹⁸, Z.D. Greenwood^{93,ai}, C. Grefe²⁴, K. Gregersen⁹⁴,
 I.M. Gregor⁴⁴, P. Grenier¹⁵⁰, K. Grevtsov⁴⁴, N.A. Grieser¹²⁴, J. Griffiths⁸, A.A. Grillo¹⁴³, K. Grimm¹⁵⁰,
 S. Grinstein^{14,y}, Ph. Gris³⁷, J.-F. Grivaz¹²⁸, S. Groh⁹⁷, E. Gross¹⁷⁷, J. Grosse-Knetter⁵¹, G.C. Grossi⁹³,
 Z.J. Grout⁹², C. Grud¹⁰³, A. Grummer¹¹⁶, L. Guan¹⁰³, W. Guan¹⁷⁸, J. Guenther³⁵, A. Guerguichon¹²⁸,
 F. Guescini^{165a}, D. Guest¹⁶⁸, R. Gugel⁵⁰, B. Gui¹²², T. Guillemin⁵, S. Guindon³⁵, U. Gul⁵⁵, C. Gumpert³⁵,
 J. Guo^{58c}, W. Guo¹⁰³, Y. Guo^{58a,r}, Z. Guo⁹⁹, R. Gupta⁴¹, S. Gurbuz^{12c}, G. Gustavino¹²⁴, B.J. Gutelman¹⁵⁷,
 P. Gutierrez¹²⁴, C. Gutscheow⁹², C. Guyot¹⁴², M.P. Guzik^{81a}, C. Gwenlan¹³¹, C.B. Gwilliam⁸⁸, A. Haas¹²¹,
 C. Haber¹⁸, H.K. Hadavand⁸, N. Haddad^{34e}, A. Hadeef^{58a}, S. Hageböck²⁴, M. Hagihara¹⁶⁶,
 H. Hakobyan^{181,*}, M. Haleem¹⁷⁴, J. Haley¹²⁵, G. Halladjian¹⁰⁴, G.D. Hallewell⁹⁹, K. Hamacher¹⁷⁹,
 P. Hamal¹²⁶, K. Hamano¹⁷³, A. Hamilton^{32a}, G.N. Hamity¹⁴⁶, K. Han^{58a,ah}, L. Han^{58a}, S. Han^{15d},
 K. Hanagaki^{79,u}, M. Hance¹⁴³, D.M. Handl¹¹², B. Haney¹³³, R. Hankache¹³², P. Hanke^{59a}, E. Hansen⁹⁴,
 J.B. Hansen³⁹, J.D. Hansen³⁹, M.C. Hansen²⁴, P.H. Hansen³⁹, K. Hara¹⁶⁶, A.S. Hard¹⁷⁸, T. Harenberg¹⁷⁹,
 S. Harkusha¹⁰⁵, P.F. Harrison¹⁷⁵, N.M. Hartmann¹¹², Y. Hasegawa¹⁴⁷, A. Hasib⁴⁸, S. Hassani¹⁴²,
 S. Haug²⁰, R. Hauser¹⁰⁴, L. Hauswald⁴⁶, L.B. Havener³⁸, M. Havranek¹³⁸, C.M. Hawkes²¹,
 R.J. Hawking³⁵, D. Hayden¹⁰⁴, C. Hayes¹⁵², C.P. Hays¹³¹, J.M. Hays⁹⁰, H.S. Hayward⁸⁸, S.J. Haywood¹⁴¹,
 M.P. Heath⁴⁸, V. Hedberg⁹⁴, L. Heelan⁸, S. Heer²⁴, K.K. Heidegger⁵⁰, J. Heilman³³, S. Heim⁴⁴,
 T. Heim¹⁸, B. Heinemann^{44,am}, J.J. Heinrich¹¹², L. Heinrich¹²¹, C. Heinz⁵⁴, J. Hejbal¹³⁷, L. Helary³⁵,
 A. Held¹⁷², S. Hellesund¹³⁰, S. Hellman^{43a,43b}, C. Helsens³⁵, R.C.W. Henderson⁸⁷, Y. Heng¹⁷⁸,
 S. Henkelmann¹⁷², A.M. Henriques Correia³⁵, G.H. Herbert¹⁹, H. Herde²⁶, V. Herget¹⁷⁴,
 Y. Hernández Jiménez^{32c}, H. Herr⁹⁷, M.G. Herrmann¹¹², G. Herten⁵⁰, R. Hertenberger¹¹², L. Hervas³⁵,
 T.C. Herwig¹³³, G.G. Hesketh⁹², N.P. Hessey^{165a}, J.W. Hetherly⁴¹, S. Higashino⁷⁹, E. Higón-Rodríguez¹⁷¹,
 K. Hildebrand³⁶, E. Hill¹⁷³, J.C. Hill³¹, K.K. Hill²⁹, K.H. Hiller⁴⁴, S.J. Hillier²¹, M. Hils⁴⁶, I. Hinchliffe¹⁸,
 M. Hirose¹²⁹, D. Hirschbuehl¹⁷⁹, B. Hiti⁸⁹, O. Hladik¹³⁷, D.R. Hlaluku^{32c}, X. Hoad⁴⁸, J. Hobbs¹⁵²,

N. Hod ^{165a}, M.C. Hodgkinson ¹⁴⁶, A. Hoecker ³⁵, M.R. Hoferkamp ¹¹⁶, F. Hoenig ¹¹², D. Hohn ²⁴,
 D. Hohov ¹²⁸, T.R. Holmes ³⁶, M. Holzbock ¹¹², M. Homann ⁴⁵, S. Honda ¹⁶⁶, T. Honda ⁷⁹, T.M. Hong ¹³⁵,
 A. Hönle ¹¹³, B.H. Hooberman ¹⁷⁰, W.H. Hopkins ¹²⁷, Y. Horii ¹¹⁵, P. Horn ⁴⁶, A.J. Horton ¹⁴⁹, L.A. Horyn ³⁶,
 J.-Y. Hostachy ⁵⁶, A. Hostiuc ¹⁴⁵, S. Hou ¹⁵⁵, A. Hoummada ^{34a}, J. Howarth ⁹⁸, J. Hoya ⁸⁶, M. Hrabovsky ¹²⁶,
 I. Hristova ¹⁹, J. Hrivnac ¹²⁸, A. Hrynevich ¹⁰⁶, T. Hryn'ova ⁵, P.J. Hsu ⁶², S.-C. Hsu ¹⁴⁵, Q. Hu ²⁹, S. Hu ^{58c},
 Y. Huang ^{15a}, Z. Hubacek ¹³⁸, F. Hubaut ⁹⁹, M. Huebner ²⁴, F. Huegging ²⁴, T.B. Huffman ¹³¹,
 E.W. Hughes ³⁸, M. Huhtinen ³⁵, R.F.H. Hunter ³³, P. Huo ¹⁵², A.M. Hupe ³³, N. Huseynov ^{77,ae},
 J. Huston ¹⁰⁴, J. Huth ⁵⁷, R. Hyneman ¹⁰³, G. Iacobucci ⁵², G. Iakovidis ²⁹, I. Ibragimov ¹⁴⁸,
 L. Iconomidou-Fayard ¹²⁸, Z. Idrissi ^{34e}, P. Iengo ³⁵, R. Ignazzi ³⁹, O. Igonkina ^{118,aa}, R. Iguchi ¹⁶⁰,
 T. Iizawa ⁵², Y. Ikegami ⁷⁹, M. Ikeno ⁷⁹, D. Iliadis ¹⁵⁹, N. Ilic ¹⁵⁰, F. Iltzsche ⁴⁶, G. Introzzi ^{68a,68b},
 M. Iodice ^{72a}, K. Iordanidou ³⁸, V. Ippolito ^{70a,70b}, M.F. Isacson ¹⁶⁹, N. Ishijima ¹²⁹, M. Ishino ¹⁶⁰,
 M. Ishitsuka ¹⁶², W. Islam ¹²⁵, C. Issever ¹³¹, S. Istin ¹⁵⁷, F. Ito ¹⁶⁶, J.M. Iturbe Ponce ^{61a}, R. Iuppa ^{73a,73b},
 A. Ivina ¹⁷⁷, H. Iwasaki ⁷⁹, J.M. Izen ⁴², V. Izzo ^{67a}, P. Jacka ¹³⁷, P. Jackson ¹, R.M. Jacobs ²⁴, V. Jain ²,
 G. Jäkel ¹⁷⁹, K.B. Jakobi ⁹⁷, K. Jakobs ⁵⁰, S. Jakobsen ⁷⁴, T. Jakoubek ¹³⁷, D.O. Jamin ¹²⁵, D.K. Jana ⁹³,
 R. Jansky ⁵², J. Janssen ²⁴, M. Janus ⁵¹, P.A. Janus ^{81a}, G. Jarlskog ⁹⁴, N. Javadov ^{77,ae}, T. Javůrek ³⁵,
 M. Javurkova ⁵⁰, F. Jeanneau ¹⁴², L. Jeanty ¹⁸, J. Jejelava ^{156a,af}, A. Jelinskas ¹⁷⁵, P. Jenni ^{50,c}, J. Jeong ⁴⁴,
 N. Jeong ⁴⁴, S. Jézéquel ⁵, H. Ji ¹⁷⁸, J. Jia ¹⁵², H. Jiang ⁷⁶, Y. Jiang ^{58a}, Z. Jiang ^{150,p}, S. Jiggins ⁵⁰,
 F.A. Jimenez Morales ³⁷, J. Jimenez Pena ¹⁷¹, S. Jin ^{15c}, A. Jinaru ^{27b}, O. Jinnouchi ¹⁶², H. Jivan ^{32c},
 P. Johansson ¹⁴⁶, K.A. Johns ⁷, C.A. Johnson ⁶³, W.J. Johnson ¹⁴⁵, K. Jon-And ^{43a,43b}, R.W.L. Jones ⁸⁷,
 S.D. Jones ¹⁵³, S. Jones ⁷, T.J. Jones ⁸⁸, J. Jongmanns ^{59a}, P.M. Jorge ^{136a,136b}, J. Jovicevic ^{165a}, X. Ju ¹⁸,
 J.J. Junggeburth ¹¹³, A. Juste Rozas ^{14,y}, A. Kaczmarska ⁸², M. Kado ¹²⁸, H. Kagan ¹²², M. Kagan ¹⁵⁰,
 T. Kaji ¹⁷⁶, E. Kajomovitz ¹⁵⁷, C.W. Kalderon ⁹⁴, A. Kaluza ⁹⁷, S. Kama ⁴¹, A. Kamenshchikov ¹⁴⁰, L. Kanjir ⁸⁹,
 Y. Kano ¹⁶⁰, V.A. Kantserov ¹¹⁰, J. Kanzaki ⁷⁹, B. Kaplan ¹²¹, L.S. Kaplan ¹⁷⁸, D. Kar ^{32c}, M.J. Kareem ^{165b},
 E. Karentzos ¹⁰, S.N. Karpov ⁷⁷, Z.M. Karpova ⁷⁷, V. Kartvelishvili ⁸⁷, A.N. Karyukhin ¹⁴⁰, L. Kashif ¹⁷⁸,
 R.D. Kass ¹²², A. Kastanas ^{43a,43b}, Y. Kataoka ¹⁶⁰, C. Kato ^{58d,58c}, J. Katzy ⁴⁴, K. Kawade ⁸⁰, K. Kawagoe ⁸⁵,
 T. Kawamoto ¹⁶⁰, G. Kawamura ⁵¹, E.F. Kay ⁸⁸, V.F. Kazanin ^{120b,120a}, R. Keeler ¹⁷³, R. Kehoe ⁴¹, J.S. Keller ³³,
 E. Kellermann ⁹⁴, J.J. Kempster ²¹, J. Kendrick ²¹, O. Kepka ¹³⁷, S. Kersten ¹⁷⁹, B.P. Kerševan ⁸⁹,
 R.A. Keyes ¹⁰¹, M. Khader ¹⁷⁰, F. Khalil-Zada ¹³, A. Khanov ¹²⁵, A.G. Kharlamov ^{120b,120a},
 T. Kharlamova ^{120b,120a}, E.E. Khoda ¹⁷², A. Khodinov ¹⁶³, T.J. Khoo ⁵², E. Khramov ⁷⁷, J. Khubua ^{156b},
 S. Kido ⁸⁰, M. Kiehn ⁵², C.R. Kilby ⁹¹, Y.K. Kim ³⁶, N. Kimura ^{64a,64c}, O.M. Kind ¹⁹, B.T. King ⁸⁸,
 D. Kirchmeier ⁴⁶, J. Kirk ¹⁴¹, A.E. Kiryunin ¹¹³, T. Kishimoto ¹⁶⁰, D. Kisielewska ^{81a}, V. Kitali ⁴⁴,
 O. Kivernyk ⁵, E. Kladiva ^{28b}, T. Klapdor-Kleingrothaus ⁵⁰, M.H. Klein ¹⁰³, M. Klein ⁸⁸, U. Klein ⁸⁸,
 K. Kleinknecht ⁹⁷, P. Klimek ¹¹⁹, A. Klimentov ²⁹, R. Klingenberg ^{45,*}, T. Klingl ²⁴, T. Klioutchnikova ³⁵,
 F.F. Klitzner ¹¹², P. Kluit ¹¹⁸, S. Kluth ¹¹³, E. Kneringer ⁷⁴, E.B.F.G. Knoops ⁹⁹, A. Knue ⁵⁰, A. Kobayashi ¹⁶⁰,
 D. Kobayashi ⁸⁵, T. Kobayashi ¹⁶⁰, M. Kobel ⁴⁶, M. Kocian ¹⁵⁰, P. Kodys ¹³⁹, P.T. Koenig ²⁴, T. Koffas ³³,
 E. Koffeman ¹¹⁸, N.M. Köhler ¹¹³, T. Koi ¹⁵⁰, M. Kolb ^{59b}, I. Koletsou ⁵, T. Kondo ⁷⁹, N. Kondrashova ^{58c},
 K. Köneke ⁵⁰, A.C. König ¹¹⁷, T. Kono ⁷⁹, R. Konoplich ^{121,aj}, V. Konstantinides ⁹², N. Konstantinidis ⁹²,
 B. Konya ⁹⁴, R. Kopeliansky ⁶³, S. Koperny ^{81a}, K. Korcyl ⁸², K. Kordas ¹⁵⁹, G. Koren ¹⁵⁸, A. Korn ⁹²,
 I. Korolkov ¹⁴, E.V. Korolkova ¹⁴⁶, N. Korotkova ¹¹¹, O. Kortner ¹¹³, S. Kortner ¹¹³, T. Kosek ¹³⁹,
 V.V. Kostyukhin ²⁴, A. Kotwal ⁴⁷, A. Koulouris ¹⁰, A. Kourkouveli-Charalampidi ^{68a,68b}, C. Kourkouvelis ⁹,
 E. Kourlitis ¹⁴⁶, V. Kouskoura ²⁹, A.B. Kowalewska ⁸², R. Kowalewski ¹⁷³, T.Z. Kowalski ^{81a}, C. Kozakai ¹⁶⁰,
 W. Kozanecki ¹⁴², A.S. Kozhin ¹⁴⁰, V.A. Kramarenko ¹¹¹, G. Kramberger ⁸⁹, D. Krasnopevtsev ^{58a},
 M.W. Krasny ¹³², A. Krasznahorkay ³⁵, D. Krauss ¹¹³, J.A. Kremer ^{81a}, J. Kretschmar ⁸⁸, P. Krieger ¹⁶⁴,
 K. Krizka ¹⁸, K. Kroeninger ⁴⁵, H. Kroha ¹¹³, J. Kroll ¹³⁷, J. Kroll ¹³³, J. Krstic ¹⁶, U. Kruchonak ⁷⁷,
 H. Krüger ²⁴, N. Krumnack ⁷⁶, M.C. Kruse ⁴⁷, T. Kubota ¹⁰², S. Kuday ^{4b}, J.T. Kuechler ¹⁷⁹, S. Kuehn ³⁵,
 A. Kugel ^{59a}, F. Kuger ¹⁷⁴, T. Kuhl ⁴⁴, V. Kukhtin ⁷⁷, R. Kukla ⁹⁹, Y. Kulchitsky ¹⁰⁵, S. Kuleshov ^{144b},
 Y.P. Kulinich ¹⁷⁰, M. Kuna ⁵⁶, T. Kunigo ⁸³, A. Kupco ¹³⁷, T. Kupfer ⁴⁵, O. Kuprash ¹⁵⁸, H. Kurashige ⁸⁰,
 L.L. Kurchaninov ^{165a}, Y.A. Kurochkin ¹⁰⁵, M.G. Kurth ^{15d}, E.S. Kuwertz ³⁵, M. Kuze ¹⁶², J. Kvita ¹²⁶,
 T. Kwan ¹⁰¹, A. La Rosa ¹¹³, J.L. La Rosa Navarro ^{78d}, L. La Rotonda ^{40b,40a}, F. La Ruffa ^{40b,40a}, C. Lacasta ¹⁷¹,
 F. Lacava ^{70a,70b}, J. Lacey ⁴⁴, D.P.J. Lack ⁹⁸, H. Lacker ¹⁹, D. Lacour ¹³², E. Ladygin ⁷⁷, R. Lafaye ⁵,
 B. Laforge ¹³², T. Lagouri ^{32c}, S. Lai ⁵¹, S. Lammers ⁶³, W. Lampl ⁷, E. Lançon ²⁹, U. Landgraf ⁵⁰,
 M.P.J. Landon ⁹⁰, M.C. Lanfermann ⁵², V.S. Lang ⁴⁴, J.C. Lange ¹⁴, R.J. Langenberg ³⁵, A.J. Lankford ¹⁶⁸,

F. Lanni²⁹, K. Lantzsch²⁴, A. Lanza^{68a}, A. Lapertosa^{53b,53a}, S. Laplace¹³², J.F. Laporte¹⁴², T. Lari^{66a}, F. Lasagni Manghi^{23b,23a}, M. Lassnig³⁵, T.S. Lau^{61a}, A. Laudrain¹²⁸, M. Lavorgna^{67a,67b}, A.T. Law¹⁴³, M. Lazzaroni^{66a,66b}, B. Le¹⁰², O. Le Dortz¹³², E. Le Guirriec⁹⁹, E.P. Le Quilleuc¹⁴², M. LeBlanc⁷, T. LeCompte⁶, F. Ledroit-Guillon⁵⁶, C.A. Lee²⁹, G.R. Lee^{144a}, L. Lee⁵⁷, S.C. Lee¹⁵⁵, B. Lefebvre¹⁰¹, M. Lefebvre¹⁷³, F. Legger¹¹², C. Leggett¹⁸, K. Lehmann¹⁴⁹, N. Lehmann¹⁷⁹, G. Lehmann Miotto³⁵, W.A. Leight⁴⁴, A. Leisos^{159,v}, M.A.L. Leite^{78d}, R. Leitner¹³⁹, D. Lellouch¹⁷⁷, B. Lemmer⁵¹, K.J.C. Leney⁹², T. Lenz²⁴, B. Lenzi³⁵, R. Leone⁷, S. Leone^{69a}, C. Leonidopoulos⁴⁸, G. Lerner¹⁵³, C. Leroy¹⁰⁷, R. Les¹⁶⁴, A.A.J. Lesage¹⁴², C.G. Lester³¹, M. Levchenko¹³⁴, J. Levêque⁵, D. Levin¹⁰³, L.J. Levinson¹⁷⁷, D. Lewis⁹⁰, B. Li¹⁰³, C.-Q. Li^{58a}, H. Li^{58b}, L. Li^{58c}, M. Li^{15a}, Q. Li^{15d}, Q.Y. Li^{58a}, S. Li^{58d,58c}, X. Li^{58c}, Y. Li¹⁴⁸, Z. Liang^{15a}, B. Liberti^{71a}, A. Liblong¹⁶⁴, K. Lie^{61c}, S. Liem¹¹⁸, A. Limosani¹⁵⁴, C.Y. Lin³¹, K. Lin¹⁰⁴, T.H. Lin⁹⁷, R.A. Linck⁶³, J.H. Lindon²¹, B.E. Lindquist¹⁵², A.L. Lioni⁵², E. Lipeles¹³³, A. Lipniacka¹⁷, M. Lisovyi^{59b}, T.M. Liss^{170,ao}, A. Lister¹⁷², A.M. Litke¹⁴³, J.D. Little⁸, B. Liu⁷⁶, B.L. Liu⁶, H.B. Liu²⁹, H. Liu¹⁰³, J.B. Liu^{58a}, J.K.K. Liu¹³¹, K. Liu¹³², M. Liu^{58a}, P. Liu¹⁸, Y. Liu^{15a}, Y.L. Liu^{58a}, Y.W. Liu^{58a}, M. Livan^{68a,68b}, A. Lleres⁵⁶, J. Llorente Merino^{15a}, S.L. Lloyd⁹⁰, C.Y. Lo^{61b}, F. Lo Sterzo⁴¹, E.M. Lobodzinska⁴⁴, P. Loch⁷, A. Loesle⁵⁰, T. Lohse¹⁹, K. Lohwasser¹⁴⁶, M. Lokajicek¹³⁷, B.A. Long²⁵, J.D. Long¹⁷⁰, R.E. Long⁸⁷, L. Longo^{65a,65b}, K.A. Looper¹²², J.A. Lopez^{144b}, I. Lopez Paz¹⁴, A. Lopez Solis¹⁴⁶, J. Lorenz¹¹², N. Lorenzo Martinez⁵, M. Losada²², P.J. Lösel¹¹², X. Lou⁴⁴, X. Lou^{15a}, A. Lounis¹²⁸, J. Love⁶, P.A. Love⁸⁷, J.J. Lozano Bahilo¹⁷¹, H. Lu^{61a}, M. Lu^{58a}, N. Lu¹⁰³, Y.J. Lu⁶², H.J. Lubatti¹⁴⁵, C. Luci^{70a,70b}, A. Lucotte⁵⁶, C. Luedtke⁵⁰, F. Luehring⁶³, I. Luise¹³², L. Luminari^{70a}, B. Lund-Jensen¹⁵¹, M.S. Lutz¹⁰⁰, P.M. Luzi¹³², D. Lynn²⁹, R. Lysak¹³⁷, E. Lytken⁹⁴, F. Lyu^{15a}, V. Lyubushkin⁷⁷, H. Ma²⁹, L.L. Ma^{58b}, Y. Ma^{58b}, G. Maccarrone⁴⁹, A. Macchiolo¹¹³, C.M. Macdonald¹⁴⁶, J. Machado Miguens^{133,136b}, D. Madaffari¹⁷¹, R. Madar³⁷, W.F. Mader⁴⁶, A. Madsen⁴⁴, N. Madysa⁴⁶, J. Maeda⁸⁰, K. Maekawa¹⁶⁰, S. Maeland¹⁷, T. Maeno²⁹, A.S. Maevskiy¹¹¹, V. Magerl⁵⁰, C. Maidantchik^{78b}, T. Maier¹¹², A. Maio^{136a,136b,136d}, O. Majersky^{28a}, S. Majewski¹²⁷, Y. Makida⁷⁹, N. Makovec¹²⁸, B. Malaescu¹³², Pa. Malecki⁸², V.P. Maleev¹³⁴, F. Malek⁵⁶, U. Mallik⁷⁵, D. Malon⁶, C. Malone³¹, S. Maltezos¹⁰, S. Malyukov³⁵, J. Mamuzic¹⁷¹, G. Mancini⁴⁹, I. Mandić⁸⁹, J. Maneira^{136a}, L. Manhaes de Andrade Filho^{78a}, J. Manjarres Ramos⁴⁶, K.H. Mankinen⁹⁴, A. Mann¹¹², A. Manousos⁷⁴, B. Mansoulie¹⁴², J.D. Mansour^{15a}, M. Mantoani⁵¹, S. Manzoni^{66a,66b}, A. Marantis¹⁵⁹, G. Marceca³⁰, L. March⁵², L. Marchese¹³¹, G. Marchiori¹³², M. Marcisovsky¹³⁷, C.A. Marin Tobon³⁵, M. Marjanovic³⁷, D.E. Marley¹⁰³, F. Marroquim^{78b}, Z. Marshall¹⁸, M.U.F. Martensson¹⁶⁹, S. Marti-Garcia¹⁷¹, C.B. Martin¹²², T.A. Martin¹⁷⁵, V.J. Martin⁴⁸, B. Martin dit Latour¹⁷, M. Martinez^{14,y}, V.I. Martinez Outschoorn¹⁰⁰, S. Martin-Haugh¹⁴¹, V.S. Martoiu^{27b}, A.C. Martyniuk⁹², A. Marzin³⁵, L. Masetti⁹⁷, T. Mashimo¹⁶⁰, R. Mashinistov¹⁰⁸, J. Masik⁹⁸, A.L. Maslennikov^{120b,120a}, L.H. Mason¹⁰², L. Massa^{71a,71b}, P. Massarotti^{67a,67b}, P. Mastrandrea⁵, A. Mastroberardino^{40b,40a}, T. Masubuchi¹⁶⁰, P. Mättig¹⁷⁹, J. Maurer^{27b}, B. Maček⁸⁹, S.J. Maxfield⁸⁸, D.A. Maximov^{120b,120a}, R. Mazini¹⁵⁵, I. Maznas¹⁵⁹, S.M. Mazza¹⁴³, N.C. Mc Fadden¹¹⁶, G. Mc Goldrick¹⁶⁴, S.P. Mc Kee¹⁰³, A. McCarn¹⁰³, T.G. McCarthy¹¹³, L.I. McClymont⁹², E.F. McDonald¹⁰², J.A. MCFayden³⁵, G. Mchedlize⁵¹, M.A. McKay⁴¹, K.D. McLean¹⁷³, S.J. McMahon¹⁴¹, P.C. McNamara¹⁰², C.J. McNicol¹⁷⁵, R.A. McPherson^{173,ac}, J.E. Mdhlluli^{32c}, Z.A. Meadows¹⁰⁰, S. Meehan¹⁴⁵, T.M. Megy⁵⁰, S. Mehlhase¹¹², A. Mehta⁸⁸, T. Meideck⁵⁶, B. Meirose⁴², D. Melini^{171,g}, B.R. Mellado Garcia^{32c}, J.D. Mellenthin⁵¹, M. Melo^{28a}, F. Meloni⁴⁴, A. Melzer²⁴, S.B. Menary⁹⁸, E.D. Mendes Gouveia^{136a}, L. Meng⁸⁸, X.T. Meng¹⁰³, A. Mengarelli^{23b,23a}, S. Menke¹¹³, E. Meoni^{40b,40a}, S. Mergelmeyer¹⁹, C. Merlassino²⁰, P. Mermod⁵², L. Merola^{67a,67b}, C. Meroni^{66a}, F.S. Merritt³⁶, A. Messina^{70a,70b}, J. Metcalfe⁶, A.S. Mete¹⁶⁸, C. Meyer¹³³, J. Meyer¹⁵⁷, J.-P. Meyer¹⁴², H. Meyer Zu Theenhausen^{59a}, F. Miano¹⁵³, R.P. Middleton¹⁴¹, L. Mijović⁴⁸, G. Mikenberg¹⁷⁷, M. Mikestikova¹³⁷, M. Mikuž⁸⁹, M. Milesi¹⁰², A. Milic¹⁶⁴, D.A. Millar⁹⁰, D.W. Miller³⁶, A. Milov¹⁷⁷, D.A. Milstead^{43a,43b}, A.A. Minaenko¹⁴⁰, M. Miñano Moya¹⁷¹, I.A. Minashvili^{156b}, A.I. Mincer¹²¹, B. Mindur^{81a}, M. Mineev⁷⁷, Y. Minegishi¹⁶⁰, Y. Ming¹⁷⁸, L.M. Mir¹⁴, A. Mirto^{65a,65b}, K.P. Mistry¹³³, T. Mitani¹⁷⁶, J. Mitrevski¹¹², V.A. Mitsou¹⁷¹, A. Miucci²⁰, P.S. Miyagawa¹⁴⁶, A. Mizukami⁷⁹, J.U. Mjörnmark⁹⁴, T. Mkrtchyan¹⁸¹, M. Mlynarikova¹³⁹, T. Moa^{43a,43b}, K. Mochizuki¹⁰⁷, P. Mogg⁵⁰, S. Mohapatra³⁸, S. Molander^{43a,43b}, R. Moles-Valls²⁴, M.C. Mondragon¹⁰⁴, K. Mönig⁴⁴, J. Monk³⁹, E. Monnier⁹⁹, A. Montalbano¹⁴⁹, J. Montejo Berlingen³⁵, F. Monticelli⁸⁶, S. Monzani^{66a}, N. Morange¹²⁸, D. Moreno²², M. Moreno Llacer³⁵, P. Morettini^{53b}, M. Morgenstern¹¹⁸,

S. Morgenstern⁴⁶, D. Mori¹⁴⁹, M. Morii⁵⁷, M. Morinaga¹⁷⁶, V. Morisbak¹³⁰, A.K. Morley³⁵, G. Mornacchi³⁵, A.P. Morris⁹², J.D. Morris⁹⁰, L. Morvaj¹⁵², P. Moschovakos¹⁰, M. Mosidze^{156b}, H.J. Moss¹⁴⁶, J. Moss^{150,m}, K. Motohashi¹⁶², R. Mount¹⁵⁰, E. Mountricha³⁵, E.J.W. Moyses¹⁰⁰, S. Muanza⁹⁹, F. Mueller¹¹³, J. Mueller¹³⁵, R.S.P. Mueller¹¹², D. Muenstermann⁸⁷, G.A. Mullier⁹⁴, F.J. Munoz Sanchez⁹⁸, P. Murin^{28b}, W.J. Murray^{175,141}, A. Murrone^{66a,66b}, M. Muškinja⁸⁹, C. Mwewa^{32a}, A.G. Myagkov^{140,ak}, J. Myers¹²⁷, M. Myska¹³⁸, B.P. Nachman¹⁸, O. Nackenhorst⁴⁵, K. Nagai¹³¹, K. Nagano⁷⁹, Y. Nagasaka⁶⁰, M. Nagel⁵⁰, E. Nagy⁹⁹, A.M. Nairz³⁵, Y. Nakahama¹¹⁵, K. Nakamura⁷⁹, T. Nakamura¹⁶⁰, I. Nakano¹²³, H. Nanjo¹²⁹, F. Napolitano^{59a}, R.F. Naranjo Garcia⁴⁴, R. Narayan¹¹, D.I. Narrias Villar^{59a}, I. Naryshkin¹³⁴, T. Naumann⁴⁴, G. Navarro²², R. Nayyar⁷, H.A. Neal¹⁰³, P.Y. Nechaeva¹⁰⁸, T.J. Neep¹⁴², A. Negri^{68a,68b}, M. Negrini^{23b}, S. Nektarijevic¹¹⁷, C. Nellist⁵¹, M.E. Nelson¹³¹, S. Nemecek¹³⁷, P. Nemethy¹²¹, M. Nessi^{35,e}, M.S. Neubauer¹⁷⁰, M. Neumann¹⁷⁹, P.R. Newman²¹, T.Y. Ng^{61c}, Y.S. Ng¹⁹, H.D.N. Nguyen⁹⁹, T. Nguyen Manh¹⁰⁷, E. Nibigira³⁷, R.B. Nickerson¹³¹, R. Nicolaidou¹⁴², D.S. Nielsen³⁹, J. Nielsen¹⁴³, N. Nikiforou¹¹, V. Nikolaenko^{140,ak}, I. Nikolic-Audit¹³², K. Nikolopoulos²¹, P. Nilsson²⁹, Y. Ninomiya⁷⁹, A. Nisati^{70a}, N. Nishu^{58c}, R. Nisius¹¹³, I. Nitsche⁴⁵, T. Nitta¹⁷⁶, T. Nobe¹⁶⁰, Y. Noguchi⁸³, M. Nomachi¹²⁹, I. Nomidis¹³², M.A. Nomura²⁹, T. Nooney⁹⁰, M. Nordberg³⁵, N. Norjoharuddeen¹³¹, T. Novak⁸⁹, O. Novgorodova⁴⁶, R. Novotny¹³⁸, L. Nozka¹²⁶, K. Ntekas¹⁶⁸, E. Nurse⁹², F. Nuti¹⁰², F.G. Oakham^{33,ar}, H. Oberlack¹¹³, T. Obermann²⁴, J. Ocariz¹³², A. Ochi⁸⁰, I. Ochoa³⁸, J.P. Ochoa-Ricoux^{144a}, K. O'Connor²⁶, S. Oda⁸⁵, S. Odaka⁷⁹, S. Oerdek⁵¹, A. Oh⁹⁸, S.H. Oh⁴⁷, C.C. Ohm¹⁵¹, H. Oide^{53b,53a}, M.L. Ojeda¹⁶⁴, H. Okawa¹⁶⁶, Y. Okazaki⁸³, Y. Okumura¹⁶⁰, T. Okuyama⁷⁹, A. Olariu^{27b}, L.F. Oleiro Seabra^{136a}, S.A. Olivares Pino^{144a}, D. Oliveira Damazio²⁹, J.L. Oliver¹, M.J.R. Olsson³⁶, A. Olszewski⁸², J. Olszowska⁸², D.C. O'Neil¹⁴⁹, A. Onofre^{136a,136e}, K. Onogi¹¹⁵, P.U.E. Onyisi¹¹, H. Oppen¹³⁰, M.J. Oreglia³⁶, G.E. Orellana⁸⁶, Y. Oren¹⁵⁸, D. Orestano^{72a,72b}, E.C. Orgill⁹⁸, N. Orlando^{61b}, A.A. O'Rourke⁴⁴, R.S. Orr¹⁶⁴, B. Osculati^{53b,53a,*}, V. O'Shea⁵⁵, R. Ospanov^{58a}, G. Otero y Garzon³⁰, H. Otono⁸⁵, M. Ouchrif^{34d}, F. Ould-Saada¹³⁰, A. Ouraou¹⁴², Q. Ouyang^{15a}, M. Owen⁵⁵, R.E. Owen²¹, V.E. Ozcan^{12c}, N. Ozturk⁸, J. Pacalt¹²⁶, H.A. Pacey³¹, K. Pachal¹⁴⁹, A. Pacheco Pages¹⁴, L. Pacheco Rodriguez¹⁴², C. Padilla Aranda¹⁴, S. Pagan Griso¹⁸, M. Paganini¹⁸⁰, G. Palacino⁶³, S. Palazzo^{40b,40a}, S. Palestini³⁵, M. Palka^{81b}, D. Pallin³⁷, I. Panagoulas¹⁰, C.E. Pandini³⁵, J.G. Panduro Vazquez⁹¹, P. Pani³⁵, G. Panizzo^{64a,64c}, L. Paolozzi⁵², T.D. Papadopoulou¹⁰, K. Papageorgiou^{9,i}, A. Paramonov⁶, D. Paredes Hernandez^{61b}, S.R. Paredes Saenz¹³¹, B. Parida¹⁶³, A.J. Parker⁸⁷, K.A. Parker⁴⁴, M.A. Parker³¹, F. Parodi^{53b,53a}, J.A. Parsons³⁸, U. Parzefall⁵⁰, V.R. Pascuzzi¹⁶⁴, J.M.P. Pasner¹⁴³, E. Pasqualucci^{70a}, S. Passaggio^{53b}, F. Pastore⁹¹, P. Pasuwan^{43a,43b}, S. Pataria⁹⁷, J.R. Pater⁹⁸, A. Pathak^{178,j}, T. Pauly³⁵, B. Pearson¹¹³, M. Pedersen¹³⁰, L. Pedraza Diaz¹¹⁷, R. Pedro^{136a,136b}, S.V. Peleganchuk^{120b,120a}, O. Penc¹³⁷, C. Peng^{15d}, H. Peng^{58a}, B.S. Peralva^{78a}, M.M. Perego¹⁴², A.P. Pereira Peixoto^{136a}, D.V. Perepelitsa²⁹, F. Peri¹⁹, L. Perini^{66a,66b}, H. Pernegger³⁵, S. Perrella^{67a,67b}, V.D. Peshekhonov^{77,*}, K. Peters⁴⁴, R.F.Y. Peters⁹⁸, B.A. Petersen³⁵, T.C. Petersen³⁹, E. Petit⁵⁶, A. Petridis¹, C. Petridou¹⁵⁹, P. Petroff¹²⁸, M. Petrov¹³¹, F. Petrucci^{72a,72b}, M. Pettee¹⁸⁰, N.E. Pettersson¹⁰⁰, A. Peyaud¹⁴², R. Pezoa^{144b}, T. Pham¹⁰², F.H. Phillips¹⁰⁴, P.W. Phillips¹⁴¹, M.W. Phipps¹⁷⁰, G. Piacquadio¹⁵², E. Pianori¹⁸, A. Picazio¹⁰⁰, M.A. Pickering¹³¹, R.H. Pickles⁹⁸, R. Piegaia³⁰, J.E. Pilcher³⁶, A.D. Pilkington⁹⁸, M. Pinamonti^{71a,71b}, J.L. Pinfold³, M. Pitt¹⁷⁷, L. Pizzimento^{71a,71b}, M-A. Pleier²⁹, V. Pleskot¹³⁹, E. Plotnikova⁷⁷, D. Pluth⁷⁶, P. Podberezko^{120b,120a}, R. Poettgen⁹⁴, R. Poggi⁵², L. Poggioli¹²⁸, I. Pogrebnyak¹⁰⁴, D. Pohl²⁴, I. Pokharel⁵¹, G. Polesello^{68a}, A. Poley¹⁸, A. Policicchio^{70a,70b}, R. Polifka³⁵, A. Polini^{23b}, C.S. Pollard⁴⁴, V. Polychronakos²⁹, D. Ponomarenko¹¹⁰, L. Pontecorvo^{70a}, G.A. Popeneciu^{27d}, D.M. Portillo Quintero¹³², S. Pospisil¹³⁸, K. Potamianos⁴⁴, I.N. Potrap⁷⁷, C.J. Potter³¹, H. Potti¹¹, T. Poulsen⁹⁴, J. Poveda³⁵, T.D. Powell¹⁴⁶, M.E. Pozo Astigarraga³⁵, P. Pralavorio⁹⁹, S. Prell⁷⁶, D. Price⁹⁸, M. Primavera^{65a}, S. Prince¹⁰¹, N. Proklova¹¹⁰, K. Prokofiev^{61c}, F. Prokoshin^{144b}, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{81a}, A. Puri¹⁷⁰, P. Puzo¹²⁸, J. Qian¹⁰³, Y. Qin⁹⁸, A. Quadt⁵¹, M. Queitsch-Maitland⁴⁴, A. Qureshi¹, P. Rados¹⁰², F. Ragusa^{66a,66b}, G. Rahal⁹⁵, J.A. Raine⁵², S. Rajagopalan²⁹, A. Ramirez Morales⁹⁰, T. Rashid¹²⁸, S. Raspopov⁵, M.G. Ratti^{66a,66b}, D.M. Rauch⁴⁴, F. Rauscher¹¹², S. Rave⁹⁷, B. Ravina¹⁴⁶, I. Ravinovich¹⁷⁷, J.H. Rawling⁹⁸, M. Raymond³⁵, A.L. Read¹³⁰, N.P. Readoff⁵⁶, M. Reale^{65a,65b}, D.M. Rebuzzi^{68a,68b}, A. Redelbach¹⁷⁴, G. Redlinger²⁹, R. Reece¹⁴³, R.G. Reed^{32c}, K. Reeves⁴², L. Rehnisch¹⁹, J. Reichert¹³³, D. Reikher¹⁵⁸, A. Reiss⁹⁷, C. Rembser³⁵, H. Ren^{15d},

M. Rescigno ^{70a}, S. Resconi ^{66a}, E.D. Resseguie ¹³³, S. Rettie ¹⁷², E. Reynolds ²¹, O.L. Rezanova ^{120b,120a},
 P. Reznicek ¹³⁹, E. Ricci ^{73a,73b}, R. Richter ¹¹³, S. Richter ⁴⁴, E. Richter-Was ^{81b}, O. Ricken ²⁴, M. Ridel ¹³²,
 P. Rieck ¹¹³, C.J. Riegel ¹⁷⁹, O. Rifki ⁴⁴, M. Rijssenbeek ¹⁵², A. Rimoldi ^{68a,68b}, M. Rimoldi ²⁰, L. Rinaldi ^{23b},
 G. Ripellino ¹⁵¹, B. Ristić ⁸⁷, E. Ritsch ³⁵, I. Riu ¹⁴, J.C. Rivera Vergara ^{144a}, F. Rizatdinova ¹²⁵, E. Rizvi ⁹⁰,
 C. Rizzi ¹⁴, R.T. Roberts ⁹⁸, S.H. Robertson ^{101.ac}, D. Robinson ³¹, J.E.M. Robinson ⁴⁴, A. Robson ⁵⁵,
 E. Rocco ⁹⁷, C. Roda ^{69a,69b}, Y. Rodina ⁹⁹, S. Rodriguez Bosca ¹⁷¹, A. Rodriguez Perez ¹⁴,
 D. Rodriguez Rodriguez ¹⁷¹, A.M. Rodríguez Vera ^{165b}, S. Roe ³⁵, C.S. Rogan ⁵⁷, O. Røhne ¹³⁰, R. Röhrig ¹¹³,
 C.P.A. Roland ⁶³, J. Roloff ⁵⁷, A. Romaniouk ¹¹⁰, M. Romano ^{23b,23a}, N. Rompotis ⁸⁸, M. Ronzani ¹²¹,
 L. Roos ¹³², S. Rosati ^{70a}, K. Rosbach ⁵⁰, P. Rose ¹⁴³, N-A. Rosien ⁵¹, B.J. Rosser ¹³³, E. Rossi ⁴⁴,
 E. Rossi ^{72a,72b}, E. Rossi ^{67a,67b}, L.P. Rossi ^{53b}, L. Rossini ^{66a,66b}, J.H.N. Rosten ³¹, R. Rosten ¹⁴, M. Rotaru ^{27b},
 J. Rothberg ¹⁴⁵, D. Rousseau ¹²⁸, D. Roy ^{32c}, A. Rozanov ⁹⁹, Y. Rozen ¹⁵⁷, X. Ruan ^{32c}, F. Rubbo ¹⁵⁰,
 F. Rühr ⁵⁰, A. Ruiz-Martinez ¹⁷¹, Z. Rurikova ⁵⁰, N.A. Rusakovich ⁷⁷, H.L. Russell ¹⁰¹, J.P. Rutherford ⁷,
 E.M. Rüttinger ^{44,k}, Y.F. Ryabov ¹³⁴, M. Rybar ¹⁷⁰, G. Rybkin ¹²⁸, S. Ryu ⁶, A. Ryzhov ¹⁴⁰, G.F. Rzehorz ⁵¹,
 P. Sabatini ⁵¹, G. Sabato ¹¹⁸, S. Sacerdoti ¹²⁸, H.F-W. Sadrozinski ¹⁴³, R. Sadykov ⁷⁷, F. Safai Tehrani ^{70a},
 P. Saha ¹¹⁹, M. Sahinsoy ^{59a}, A. Sahu ¹⁷⁹, M. Saimpert ⁴⁴, M. Saito ¹⁶⁰, T. Saito ¹⁶⁰, H. Sakamoto ¹⁶⁰,
 A. Sakharov ^{121,aj}, D. Salamani ⁵², G. Salamanna ^{72a,72b}, J.E. Salazar Loyola ^{144b}, P.H. Sales De Bruin ¹⁶⁹,
 D. Salihagic ¹¹³, A. Salnikov ¹⁵⁰, J. Salt ¹⁷¹, D. Salvatore ^{40b,40a}, F. Salvatore ¹⁵³, A. Salvucci ^{61a,61b,61c},
 A. Salzburger ³⁵, J. Samarati ³⁵, D. Sammel ⁵⁰, D. Sampsonidis ¹⁵⁹, D. Sampsonidou ¹⁵⁹, J. Sánchez ¹⁷¹,
 A. Sanchez Pineda ^{64a,64c}, H. Sandaker ¹³⁰, C.O. Sander ⁴⁴, M. Sandhoff ¹⁷⁹, C. Sandoval ²²,
 D.P.C. Sankey ¹⁴¹, M. Sannino ^{53b,53a}, Y. Sano ¹¹⁵, A. Sansoni ⁴⁹, C. Santoni ³⁷, H. Santos ^{136a},
 I. Santoyo Castillo ¹⁵³, A. Santra ¹⁷¹, A. Saponov ⁷⁷, J.G. Saraiva ^{136a,136d}, O. Sasaki ⁷⁹, K. Sato ¹⁶⁶,
 E. Sauvan ⁵, P. Savard ^{164.ar}, N. Savic ¹¹³, R. Sawada ¹⁶⁰, C. Sawyer ¹⁴¹, L. Sawyer ^{93.ai}, C. Sbarra ^{23b},
 A. Sbrizzi ^{23b,23a}, T. Scanlon ⁹², J. Schaarschmidt ¹⁴⁵, P. Schacht ¹¹³, B.M. Schachtner ¹¹², D. Schaefer ³⁶,
 L. Schaefer ¹³³, J. Schaeffer ⁹⁷, S. Schaepe ³⁵, U. Schäfer ⁹⁷, A.C. Schaffer ¹²⁸, D. Schaile ¹¹²,
 R.D. Schamberger ¹⁵², N. Scharmberg ⁹⁸, V.A. Schegelsky ¹³⁴, D. Scheirich ¹³⁹, F. Schenck ¹⁹,
 M. Schernau ¹⁶⁸, C. Schiavi ^{53b,53a}, S. Schier ¹⁴³, L.K. Schildgen ²⁴, Z.M. Schillaci ²⁶, E.J. Schioppa ³⁵,
 M. Schioppa ^{40b,40a}, K.E. Schleicher ⁵⁰, S. Schlenker ³⁵, K.R. Schmidt-Sommerfeld ¹¹³, K. Schmieden ³⁵,
 C. Schmitt ⁹⁷, S. Schmitt ⁴⁴, S. Schmitz ⁹⁷, J.C. Schmoeckel ⁴⁴, U. Schnoor ⁵⁰, L. Schoeffel ¹⁴²,
 A. Schoening ^{59b}, E. Schopf ¹³¹, M. Schott ⁹⁷, J.F.P. Schouwenberg ¹¹⁷, J. Schovancova ³⁵, S. Schramm ⁵²,
 A. Schulte ⁹⁷, H-C. Schultz-Coulon ^{59a}, M. Schumacher ⁵⁰, B.A. Schumm ¹⁴³, Ph. Schune ¹⁴²,
 A. Schwartzman ¹⁵⁰, T.A. Schwarz ¹⁰³, Ph. Schwemling ¹⁴², R. Schwienhorst ¹⁰⁴, A. Sciandra ²⁴,
 G. Sciolla ²⁶, M. Scornajenghi ^{40b,40a}, F. Scuri ^{69a}, F. Scutti ¹⁰², L.M. Scyboz ¹¹³, J. Searcy ¹⁰³,
 C.D. Sebastiani ^{70a,70b}, P. Seema ¹⁹, S.C. Seidel ¹¹⁶, A. Seiden ¹⁴³, T. Seiss ³⁶, J.M. Seixas ^{78b},
 G. Sekhniaidze ^{67a}, K. Sekhon ¹⁰³, S.J. Sekula ⁴¹, N. Semprini-Cesari ^{23b,23a}, S. Sen ⁴⁷, S. Senkin ³⁷,
 C. Serfon ¹³⁰, L. Serin ¹²⁸, L. Serkin ^{64a,64b}, M. Sessa ^{58a}, H. Severini ¹²⁴, F. Sforza ¹⁶⁷, A. Sfyrla ⁵²,
 E. Shabalina ⁵¹, J.D. Shahinian ¹⁴³, N.W. Shaikh ^{43a,43b}, L.Y. Shan ^{15a}, R. Shang ¹⁷⁰, J.T. Shank ²⁵,
 M. Shapiro ¹⁸, A.S. Sharma ¹, A. Sharma ¹³¹, P.B. Shatalov ¹⁰⁹, K. Shaw ¹⁵³, S.M. Shaw ⁹⁸,
 A. Shcherbakova ¹³⁴, Y. Shen ¹²⁴, N. Sherafati ³³, A.D. Sherman ²⁵, P. Sherwood ⁹², L. Shi ^{155.an},
 S. Shimizu ⁷⁹, C.O. Shimmin ¹⁸⁰, M. Shimojima ¹¹⁴, I.P.J. Shipsey ¹³¹, S. Shirabe ⁸⁵, M. Shiyakova ⁷⁷,
 J. Shlomi ¹⁷⁷, A. Shmeleva ¹⁰⁸, D. Shoaleh Saadi ¹⁰⁷, M.J. Shochet ³⁶, S. Shojaii ¹⁰², D.R. Shope ¹²⁴,
 S. Shrestha ¹²², E. Shulga ¹¹⁰, P. Sicho ¹³⁷, A.M. Sickles ¹⁷⁰, P.E. Sidebo ¹⁵¹, E. Sideras Haddad ^{32c},
 O. Sidiropoulou ³⁵, A. Sidoti ^{23b,23a}, F. Siegert ⁴⁶, Dj. Sijacki ¹⁶, J. Silva ^{136a}, M. Silva Jr. ¹⁷⁸,
 M.V. Silva Oliveira ^{78a}, S.B. Silverstein ^{43a}, S. Simion ¹²⁸, E. Simioni ⁹⁷, M. Simon ⁹⁷, R. Simoniello ⁹⁷,
 P. Sinervo ¹⁶⁴, N.B. Sinev ¹²⁷, M. Sioli ^{23b,23a}, G. Siragusa ¹⁷⁴, I. Siral ¹⁰³, S.Yu. Sivoklokov ¹¹¹,
 J. Sjölin ^{43a,43b}, P. Skubic ¹²⁴, M. Slater ²¹, T. Slavicek ¹³⁸, M. Slawinska ⁸², K. Sliwa ¹⁶⁷, R. Slovak ¹³⁹,
 V. Smakhtin ¹⁷⁷, B.H. Smart ⁵, J. Smiesko ^{28a}, N. Smirnov ¹¹⁰, S.Yu. Smirnov ¹¹⁰, Y. Smirnov ¹¹⁰,
 L.N. Smirnova ¹¹¹, O. Smirnova ⁹⁴, J.W. Smith ⁵¹, M.N.K. Smith ³⁸, M. Smizanska ⁸⁷, K. Smolek ¹³⁸,
 A. Smykiewicz ⁸², A.A. Snesarev ¹⁰⁸, I.M. Snyder ¹²⁷, S. Snyder ²⁹, R. Sobie ^{173.ac}, A.M. Soffa ¹⁶⁸,
 A. Soffer ¹⁵⁸, A. Søgaard ⁴⁸, D.A. Soh ¹⁵⁵, G. Sokhrannyi ⁸⁹, C.A. Solans Sanchez ³⁵, M. Solar ¹³⁸,
 E.Yu. Soldatov ¹¹⁰, U. Soldevila ¹⁷¹, A.A. Solodkov ¹⁴⁰, A. Soloshenko ⁷⁷, O.V. Solovyanov ¹⁴⁰,
 V. Solovyev ¹³⁴, P. Sommer ¹⁴⁶, H. Son ¹⁶⁷, W. Song ¹⁴¹, W.Y. Song ^{165b}, A. Sopczak ¹³⁸, F. Sopkova ^{28b},
 C.L. Sotiropoulou ^{69a,69b}, S. Sottocornola ^{68a,68b}, R. Soualah ^{64a,64c,h}, A.M. Soukharev ^{120b,120a}, D. South ⁴⁴,

B.C. Sowden⁹¹, S. Spagnolo^{65a,65b}, M. Spalla¹¹³, M. Spangenberg¹⁷⁵, F. Spanò⁹¹, D. Sperlich¹⁹,
 F. Spettel¹¹³, T.M. Spieker^{59a}, R. Spighi^{23b}, G. Spigo³⁵, L.A. Spiller¹⁰², D.P. Spiteri⁵⁵, M. Spousta¹³⁹,
 A. Stabile^{66a,66b}, R. Stamen^{59a}, S. Stamm¹⁹, E. Stanecka⁸², R.W. Stanek⁶, C. Stanescu^{72a},
 B. Stanislaus¹³¹, M.M. Stanitzki⁴⁴, B. Stapf¹¹⁸, S. Stapnes¹³⁰, E.A. Starchenko¹⁴⁰, G.H. Stark³⁶, J. Stark⁵⁶,
 S.H. Stark³⁹, P. Staroba¹³⁷, P. Starovoitov^{59a}, S. Stärz³⁵, R. Staszewski⁸², M. Stegler⁴⁴, P. Steinberg²⁹,
 B. Stelzer¹⁴⁹, H.J. Stelzer³⁵, O. Stelzer-Chilton^{165a}, H. Stenzel⁵⁴, T.J. Stevenson⁹⁰, G.A. Stewart⁵⁵,
 M.C. Stockton¹²⁷, G. Stoicea^{27b}, P. Stolte⁵¹, S. Stonjek¹¹³, A. Straessner⁴⁶, J. Strandberg¹⁵¹,
 S. Strandberg^{43a,43b}, M. Strauss¹²⁴, P. Strizenec^{28b}, R. Ströhmer¹⁷⁴, D.M. Strom¹²⁷, R. Stroynowski⁴¹,
 A. Strubig⁴⁸, S.A. Stucci²⁹, B. Stugu¹⁷, J. Stupak¹²⁴, N.A. Styles⁴⁴, D. Su¹⁵⁰, J. Su¹³⁵, S. Suchek^{59a},
 Y. Sugaya¹²⁹, M. Suk¹³⁸, V.V. Sulim¹⁰⁸, M.J. Sullivan⁸⁸, D.M.S. Sultan⁵², S. Sultansoy^{4c}, T. Sumida⁸³,
 S. Sun¹⁰³, X. Sun³, K. Suruliz¹⁵³, C.J.E. Suster¹⁵⁴, M.R. Sutton¹⁵³, S. Suzuki⁷⁹, M. Svatos¹³⁷,
 M. Swiatlowski³⁶, S.P. Swift², A. Sydorenko⁹⁷, I. Sykora^{28a}, T. Sykora¹³⁹, D. Ta⁹⁷, K. Tackmann^{44,z},
 J. Taenzer¹⁵⁸, A. Taffard¹⁶⁸, R. Tafirout^{165a}, E. Tahirovic⁹⁰, N. Taiblum¹⁵⁸, H. Takai²⁹, R. Takashima⁸⁴,
 E.H. Takasugi¹¹³, K. Takeda⁸⁰, T. Takeshita¹⁴⁷, Y. Takubo⁷⁹, M. Talby⁹⁹, A.A. Talyshev^{120b,120a},
 J. Tanaka¹⁶⁰, M. Tanaka¹⁶², R. Tanaka¹²⁸, B.B. Tannenwald¹²², S. Tapia Araya^{144b}, S. Tapprogge⁹⁷,
 A. Tarek Abouelfadl Mohamed¹³², S. Tarem¹⁵⁷, G. Tarna^{27b,d}, G.F. Tartarelli^{66a}, P. Tas¹³⁹,
 M. Tasevsky¹³⁷, T. Tashiro⁸³, E. Tassi^{40b,40a}, A. Tavares Delgado^{136a,136b}, Y. Tayalati^{34e}, A.C. Taylor¹¹⁶,
 A.J. Taylor⁴⁸, G.N. Taylor¹⁰², P.T.E. Taylor¹⁰², W. Taylor^{165b}, A.S. Tee⁸⁷, P. Teixeira-Dias⁹¹, H. Ten Kate³⁵,
 P.K. Teng¹⁵⁵, J.J. Teoh¹¹⁸, S. Terada⁷⁹, K. Terashi¹⁶⁰, J. Terron⁹⁶, S. Terzo¹⁴, M. Testa⁴⁹,
 R.J. Teuscher^{164,ac}, S.J. Thais¹⁸⁰, T. Theveneaux-Pelzer⁴⁴, F. Thiele³⁹, D.W. Thomas⁹¹, J.P. Thomas²¹,
 A.S. Thompson⁵⁵, P.D. Thompson²¹, L.A. Thomsen¹⁸⁰, E. Thomson¹³³, Y. Tian³⁸, R.E. Ticse Torres⁵¹,
 V.O. Tikhomirov^{108,al}, Yu.A. Tikhonov^{120b,120a}, S. Timoshenko¹¹⁰, P. Tipton¹⁸⁰, S. Tisserant⁹⁹,
 K. Todome¹⁶², S. Todorova-Nova⁵, S. Todt⁴⁶, J. Tojo⁸⁵, S. Tokár^{28a}, K. Tokushuku⁷⁹, E. Tolley¹²²,
 K.G. Tomiwa^{32c}, M. Tomoto¹¹⁵, L. Tompkins^{150,p}, K. Toms¹¹⁶, B. Tong⁵⁷, P. Tornambe⁵⁰, E. Torrence¹²⁷,
 H. Torres⁴⁶, E. Torrò Pastor¹⁴⁵, C. Tosciri¹³¹, J. Toth^{99,ab}, F. Touchard⁹⁹, D.R. Tovey¹⁴⁶, C.J. Treado¹²¹,
 T. Trefzger¹⁷⁴, F. Tresoldi¹⁵³, A. Tricoli²⁹, I.M. Trigger^{165a}, S. Trincaz-Duvoid¹³², M.F. Tripijana¹⁴,
 W. Trischuk¹⁶⁴, B. Trocmé⁵⁶, A. Trofymov¹²⁸, C. Troncon^{66a}, M. Trovatelli¹⁷³, F. Trovato¹⁵³,
 L. Truong^{32b}, M. Trzebinski⁸², A. Trzupek⁸², F. Tsai⁴⁴, J.C-L. Tseng¹³¹, P.V. Tsiarehka¹⁰⁵, A. Tsigotis¹⁵⁹,
 N. Tsirintanis⁹, V. Tsiskaridze¹⁵², E.G. Tskhadadze^{156a}, I.I. Tsukerman¹⁰⁹, V. Tsulaia¹⁸, S. Tsuno⁷⁹,
 D. Tsybychev^{152,163}, Y. Tu^{61b}, A. Tudorache^{27b}, V. Tudorache^{27b}, T.T. Tulbure^{27a}, A.N. Tuna⁵⁷,
 S. Turchikhin⁷⁷, D. Turgeman¹⁷⁷, I. Turk Cakir^{4b,t}, R. Turra^{66a}, P.M. Tuts³⁸, E. Tzovara⁹⁷,
 G. Uccielli^{23b,23a}, I. Ueda⁷⁹, M. Ughetto^{43a,43b}, F. Ukegawa¹⁶⁶, G. Unal³⁵, A. Undrus²⁹, G. Unel¹⁶⁸,
 F.C. Ungaro¹⁰², Y. Unno⁷⁹, K. Uno¹⁶⁰, J. Urban^{28b}, P. Urquijo¹⁰², P. Urrejola⁹⁷, G. Usai⁸, J. Usui⁷⁹,
 L. Vacavant⁹⁹, V. Vacek¹³⁸, B. Vachon¹⁰¹, K.O.H. Vadla¹³⁰, A. Vaidya⁹², C. Valderanis¹¹²,
 E. Valdes Santurio^{43a,43b}, M. Valente⁵², S. Valentini^{23b,23a}, A. Valero¹⁷¹, L. Valéry⁴⁴, R.A. Vallance²¹,
 A. Vallier⁵, J.A. Valls Ferrer¹⁷¹, T.R. Van Daalen¹⁴, H. Van der Graaf¹¹⁸, P. Van Gemmeren⁶,
 J. Van Nieuwkoop¹⁴⁹, I. Van Vulpen¹¹⁸, M. Vanadia^{71a,71b}, W. Vandelli³⁵, A. Vaniachine¹⁶³,
 P. Vankov¹¹⁸, R. Vari^{70a}, E.W. Varnes⁷, C. Varni^{53b,53a}, T. Varol⁴¹, D. Varouchas¹²⁸, K.E. Varvell¹⁵⁴,
 G.A. Vasquez^{144b}, J.G. Vasquez¹⁸⁰, F. Vazeille³⁷, D. Vazquez Furelos¹⁴, T. Vazquez Schroeder¹⁰¹,
 J. Veatch⁵¹, V. Vecchio^{72a,72b}, L.M. Veloce¹⁶⁴, F. Veloso^{136a,136c}, S. Veneziano^{70a}, A. Ventura^{65a,65b},
 M. Venturi¹⁷³, N. Venturi³⁵, V. Vercesi^{68a}, M. Verducci^{72a,72b}, C.M. Vergel Infante⁷⁶, C. Vergis²⁴,
 W. Verkerke¹¹⁸, A.T. Vermeulen¹¹⁸, J.C. Vermeulen¹¹⁸, M.C. Vetterli^{149,ar}, N. Viaux Maira^{144b},
 M. Vicente Barreto Pinto⁵², I. Vichou^{170,*}, T. Vickey¹⁴⁶, O.E. Vickey Boeriu¹⁴⁶, G.H.A. Viehhauser¹³¹,
 S. Viel¹⁸, L. Vigani¹³¹, M. Villa^{23b,23a}, M. Villaplana Perez^{66a,66b}, E. Vilucchi⁴⁹, M.G. Vincker³³,
 V.B. Vinogradov⁷⁷, A. Vishwakarma⁴⁴, C. Vittori^{23b,23a}, I. Vivarelli¹⁵³, S. Vlachos¹⁰, M. Vogel¹⁷⁹,
 P. Vokac¹³⁸, G. Volpi¹⁴, S.E. von Buddenbrock^{32c}, E. Von Toerne²⁴, V. Vorobel¹³⁹, K. Vorobev¹¹⁰,
 M. Vos¹⁷¹, J.H. Vosseveld⁸⁸, N. Vranjes¹⁶, M. Vranjes Milosavljevic¹⁶, V. Vrba¹³⁸, M. Vreeswijk¹¹⁸,
 T. Šfiligoj⁸⁹, R. Vuillermet³⁵, I. Vukotic³⁶, T. Ženiš^{28a}, L. Živković¹⁶, P. Wagner²⁴, W. Wagner¹⁷⁹,
 J. Wagner-Kuhr¹¹², H. Wahlberg⁸⁶, S. Wahrmund⁴⁶, K. Wakamiya⁸⁰, V.M. Walbrecht¹¹³, J. Walder⁸⁷,
 R. Walker¹¹², S.D. Walker⁹¹, W. Walkowiak¹⁴⁸, V. Wallangen^{43a,43b}, A.M. Wang⁵⁷, C. Wang^{58b,d},
 F. Wang¹⁷⁸, H. Wang¹⁸, H. Wang³, J. Wang¹⁵⁴, J. Wang^{59b}, P. Wang⁴¹, Q. Wang¹²⁴, R.-J. Wang¹³²,
 R. Wang^{58a}, R. Wang⁶, S.M. Wang¹⁵⁵, W.T. Wang^{58a}, W. Wang^{15c,ad}, W.X. Wang^{58a}, Y. Wang^{58a},

Z. Wang^{58c}, C. Wanotayaroj⁴⁴, A. Warburton¹⁰¹, C.P. Ward³¹, D.R. Wardrope⁹², A. Washbrook⁴⁸, P.M. Watkins²¹, A.T. Watson²¹, M.F. Watson²¹, G. Watts¹⁴⁵, S. Watts⁹⁸, B.M. Waugh⁹², A.F. Webb¹¹, S. Webb⁹⁷, C. Weber¹⁸⁰, M.S. Weber²⁰, S.A. Weber³³, S.M. Weber^{59a}, A.R. Weidberg¹³¹, B. Weinert⁶³, J. Weingarten⁵¹, M. Weirich⁹⁷, C. Weiser⁵⁰, P.S. Wells³⁵, T. Wenaus²⁹, T. Wengler³⁵, S. Wenig³⁵, N. Wermes²⁴, M.D. Werner⁷⁶, P. Werner³⁵, M. Wessels^{59a}, T.D. Weston²⁰, K. Whalen¹²⁷, N.L. Whallon¹⁴⁵, A.M. Wharton⁸⁷, A.S. White¹⁰³, A. White⁸, M.J. White¹, R. White^{144b}, D. Whiteson¹⁶⁸, B.W. Whitmore⁸⁷, F.J. Wickens¹⁴¹, W. Wiedenmann¹⁷⁸, M. Wielers¹⁴¹, C. Wiglesworth³⁹, L.A.M. Wiik-Fuchs⁵⁰, F. Wilk⁹⁸, H.G. Wilkens³⁵, L.J. Wilkins⁹¹, H.H. Williams¹³³, S. Williams³¹, C. Willis¹⁰⁴, S. Willocq¹⁰⁰, J.A. Wilson²¹, I. Wingerter-Seez⁵, E. Winkels¹⁵³, F. Winklmeier¹²⁷, O.J. Winston¹⁵³, B.T. Winter²⁴, M. Wittgen¹⁵⁰, M. Wobisch⁹³, A. Wolf⁹⁷, T.M.H. Wolf¹¹⁸, R. Wolff⁹⁹, M.W. Wolter⁸², H. Wolters^{136a,136c}, V.W.S. Wong¹⁷², N.L. Woods¹⁴³, S.D. Worm²¹, B.K. Wosiek⁸², K.W. Woźniak⁸², K. Wraight⁵⁵, M. Wu³⁶, S.L. Wu¹⁷⁸, X. Wu⁵², Y. Wu^{58a}, T.R. Wyatt⁹⁸, B.M. Wynne⁴⁸, S. Xella³⁹, Z. Xi¹⁰³, L. Xia¹⁷⁵, D. Xu^{15a}, H. Xu^{58a}, L. Xu²⁹, T. Xu¹⁴², W. Xu¹⁰³, B. Yabsley¹⁵⁴, S. Yacoob^{32a}, K. Yajima¹²⁹, D.P. Yallup⁹², D. Yamaguchi¹⁶², Y. Yamaguchi¹⁶², A. Yamamoto⁷⁹, T. Yamanaka¹⁶⁰, F. Yamane⁸⁰, M. Yamatani¹⁶⁰, T. Yamazaki¹⁶⁰, Y. Yamazaki⁸⁰, Z. Yan²⁵, H.J. Yang^{58c,58d}, H.T. Yang¹⁸, S. Yang⁷⁵, Y. Yang¹⁶⁰, Z. Yang¹⁷, W.-M. Yao¹⁸, Y.C. Yap⁴⁴, Y. Yasu⁷⁹, E. Yatsenko^{58c,58d}, J. Ye⁴¹, S. Ye²⁹, I. Yeletsikh⁷⁷, E. Yigitbasi²⁵, E. Yildirim⁹⁷, K. Yorita¹⁷⁶, K. Yoshihara¹³³, C.J.S. Young³⁵, C. Young¹⁵⁰, J. Yu⁸, J. Yu⁷⁶, X. Yue^{59a}, S.P.Y. Yuen²⁴, B. Zabinski⁸², G. Zacharis¹⁰, E. Zaffaroni⁵², R. Zaidan¹⁴, A.M. Zaitsev^{140,ak}, T. Zakareishvili^{156b}, N. Zakharchuk³³, J. Zalieckas¹⁷, S. Zambito⁵⁷, D. Zanzi³⁵, D.R. Zaripovas⁵⁵, S.V. Zeiřner⁴⁵, C. Zeitnitz¹⁷⁹, G. Zemaityte¹³¹, J.C. Zeng¹⁷⁰, Q. Zeng¹⁵⁰, O. Zenin¹⁴⁰, D. Zerwas¹²⁸, M. Zgubič¹³¹, D.F. Zhang^{58b}, D. Zhang¹⁰³, F. Zhang¹⁷⁸, G. Zhang^{58a}, H. Zhang^{15c}, J. Zhang⁶, L. Zhang^{15c}, L. Zhang^{58a}, M. Zhang¹⁷⁰, P. Zhang^{15c}, R. Zhang^{58a}, R. Zhang²⁴, X. Zhang^{58b}, Y. Zhang^{15d}, Z. Zhang¹²⁸, P. Zhao⁴⁷, X. Zhao⁴¹, Y. Zhao^{58b,128,ah}, Z. Zhao^{58a}, A. Zhemchugov⁷⁷, Z. Zheng¹⁰³, D. Zhong¹⁷⁰, B. Zhou¹⁰³, C. Zhou¹⁷⁸, L. Zhou⁴¹, M.S. Zhou^{15d}, M. Zhou¹⁵², N. Zhou^{58c}, Y. Zhou⁷, C.G. Zhu^{58b}, H.L. Zhu^{58a}, H. Zhu^{15a}, J. Zhu¹⁰³, Y. Zhu^{58a}, X. Zhuang^{15a}, K. Zhukov¹⁰⁸, V. Zhulanov^{120b,120a}, A. Zibell¹⁷⁴, D. Zieminska⁶³, N.I. Zimine⁷⁷, S. Zimmermann⁵⁰, Z. Zinonos¹¹³, M. Zinser⁹⁷, M. Ziolkowski¹⁴⁸, G. Zobernig¹⁷⁸, A. Zoccoli^{23b,23a}, K. Zoch⁵¹, T.G. Zorbas¹⁴⁶, R. Zou³⁶, M. Zur Nedden¹⁹, L. Zwalinski³⁵

¹ Department of Physics, University of Adelaide, Adelaide, Australia

² Physics Department, SUNY Albany, Albany, NY, United States of America

³ Department of Physics, University of Alberta, Edmonton, AB, Canada

⁴ (a) Department of Physics, Ankara University, Ankara; (b) Istanbul Aydin University, Istanbul; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey

⁵ LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France

⁶ High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States of America

⁷ Department of Physics, University of Arizona, Tucson, AZ, United States of America

⁸ Department of Physics, University of Texas at Arlington, Arlington, TX, United States of America

⁹ Physics Department, National and Kapodistrian University of Athens, Athens, Greece

¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece

¹¹ Department of Physics, University of Texas at Austin, Austin, TX, United States of America

¹² (a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; (b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; (c) Department of Physics, Bogazici University, Istanbul; (d) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey

¹³ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

¹⁴ Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain

¹⁵ (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Physics Department, Tsinghua University, Beijing; (c) Department of Physics, Nanjing University, Nanjing;

(d) University of Chinese Academy of Science (UCAS), Beijing, China

¹⁶ Institute of Physics, University of Belgrade, Belgrade, Serbia

¹⁷ Department for Physics and Technology, University of Bergen, Bergen, Norway

¹⁸ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States of America

¹⁹ Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany

²⁰ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

²¹ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom

²² Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia

²³ (a) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna; (b) INFN Sezione di Bologna, Italy

²⁴ Physikalisches Institut, Universität Bonn, Bonn, Germany

²⁵ Department of Physics, Boston University, Boston, MA, United States of America

²⁶ Department of Physics, Brandeis University, Waltham, MA, United States of America

²⁷ (a) Transilvania University of Brasov, Brasov; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; (c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; (d) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; (e) University Politehnica Bucharest, Bucharest; (f) West University in Timisoara, Timisoara, Romania

²⁸ (a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

²⁹ Physics Department, Brookhaven National Laboratory, Upton, NY, United States of America

³⁰ Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina

³¹ Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom

- ³² (a) Department of Physics, University of Cape Town, Cape Town; (b) Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; (c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ³³ Department of Physics, Carleton University, Ottawa, ON, Canada
- ³⁴ (a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies – Université Hassan II, Casablanca; (b) Centre National de l'Energie des Sciences Techniques Nucleaires (CNESTEN), Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA, Marrakech; (d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (e) Faculté des sciences, Université Mohammed V, Rabat, Morocco
- ³⁵ CERN, Geneva, Switzerland
- ³⁶ Enrico Fermi Institute, University of Chicago, Chicago, IL, United States of America
- ³⁷ LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
- ³⁸ Nevis Laboratory, Columbia University, Irvington, NY, United States of America
- ³⁹ Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- ⁴⁰ (a) Dipartimento di Fisica, Università della Calabria, Rende; (b) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
- ⁴¹ Physics Department, Southern Methodist University, Dallas, TX, United States of America
- ⁴² Physics Department, University of Texas at Dallas, Richardson, TX, United States of America
- ⁴³ (a) Department of Physics, Stockholm University; (b) Oskar Klein Centre, Stockholm, Sweden
- ⁴⁴ Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
- ⁴⁵ Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- ⁴⁶ Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- ⁴⁷ Department of Physics, Duke University, Durham, NC, United States of America
- ⁴⁸ SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- ⁴⁹ INFN e Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵⁰ Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
- ⁵¹ II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
- ⁵² Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
- ⁵³ (a) Dipartimento di Fisica, Università di Genova, Genova; (b) INFN Sezione di Genova, Italy
- ⁵⁴ II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵⁵ SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- ⁵⁶ LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
- ⁵⁷ Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States of America
- ⁵⁸ (a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; (b) Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; (c) School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; (d) Tsung-Dao Lee Institute, Shanghai, China
- ⁵⁹ (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ⁶⁰ Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- ⁶¹ (a) Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; (b) Department of Physics, University of Hong Kong, Hong Kong; (c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- ⁶² Department of Physics, National Tsing Hua University, Hsinchu, Taiwan
- ⁶³ Department of Physics, Indiana University, Bloomington, IN, United States of America
- ⁶⁴ (a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; (b) ICTP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
- ⁶⁵ (a) INFN Sezione di Lecce; (b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- ⁶⁶ (a) INFN Sezione di Milano; (b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- ⁶⁷ (a) INFN Sezione di Napoli; (b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- ⁶⁸ (a) INFN Sezione di Pavia; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- ⁶⁹ (a) INFN Sezione di Pisa; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- ⁷⁰ (a) INFN Sezione di Roma; (b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
- ⁷¹ (a) INFN Sezione di Roma Tor Vergata; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- ⁷² (a) INFN Sezione di Roma Tre; (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
- ⁷³ (a) INFN-TIFPA; (b) Università degli Studi di Trento, Trento, Italy
- ⁷⁴ Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- ⁷⁵ University of Iowa, Iowa City, IA, United States of America
- ⁷⁶ Department of Physics and Astronomy, Iowa State University, Ames, IA, United States of America
- ⁷⁷ Joint Institute for Nuclear Research, Dubna, Russia
- ⁷⁸ (a) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; (b) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; (c) Universidade Federal de São João del Rei (UFSJ), São João del Rei; (d) Instituto de Física, Universidade de São Paulo, São Paulo, Brazil
- ⁷⁹ KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- ⁸⁰ Graduate School of Science, Kobe University, Kobe, Japan
- ⁸¹ (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
- ⁸² Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
- ⁸³ Faculty of Science, Kyoto University, Kyoto, Japan
- ⁸⁴ Kyoto University of Education, Kyoto, Japan
- ⁸⁵ Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan
- ⁸⁶ Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- ⁸⁷ Physics Department, Lancaster University, Lancaster, United Kingdom
- ⁸⁸ Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- ⁸⁹ Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
- ⁹⁰ School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- ⁹¹ Department of Physics, Royal Holloway University of London, Egham, United Kingdom
- ⁹² Department of Physics and Astronomy, University College London, London, United Kingdom
- ⁹³ Louisiana Tech University, Ruston, LA, United States of America
- ⁹⁴ Fysiska institutionen, Lunds universitet, Lund, Sweden
- ⁹⁵ Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
- ⁹⁶ Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain
- ⁹⁷ Institut für Physik, Universität Mainz, Mainz, Germany
- ⁹⁸ School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- ⁹⁹ CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- ¹⁰⁰ Department of Physics, University of Massachusetts, Amherst, MA, United States of America
- ¹⁰¹ Department of Physics, McGill University, Montreal, QC, Canada
- ¹⁰² School of Physics, University of Melbourne, Victoria, Australia

- 103 Department of Physics, University of Michigan, Ann Arbor, MI, United States of America
- 104 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, United States of America
- 105 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
- 106 Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus
- 107 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
- 108 P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
- 109 Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
- 110 National Research Nuclear University MEPhI, Moscow, Russia
- 111 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- 112 Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- 113 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- 114 Nagasaki Institute of Applied Science, Nagasaki, Japan
- 115 Graduate School of Science and Kobayashi–Maskawa Institute, Nagoya University, Nagoya, Japan
- 116 Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States of America
- 117 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- 118 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- 119 Department of Physics, Northern Illinois University, DeKalb, IL, United States of America
- 120 ^(a) Budker Institute of Nuclear Physics, SB RAS, Novosibirsk; ^(b) Novosibirsk State University, Novosibirsk, Russia
- 121 Department of Physics, New York University, New York, NY, United States of America
- 122 Ohio State University, Columbus, OH, United States of America
- 123 Faculty of Science, Okayama University, Okayama, Japan
- 124 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States of America
- 125 Department of Physics, Oklahoma State University, Stillwater, OK, United States of America
- 126 Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc, Czech Republic
- 127 Center for High Energy Physics, University of Oregon, Eugene, OR, United States of America
- 128 LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
- 129 Graduate School of Science, Osaka University, Osaka, Japan
- 130 Department of Physics, University of Oslo, Oslo, Norway
- 131 Department of Physics, Oxford University, Oxford, United Kingdom
- 132 LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France
- 133 Department of Physics, University of Pennsylvania, Philadelphia, PA, United States of America
- 134 Konstantinov Nuclear Physics Institute of National Research Centre “Kurchatov Institute”, PNPI, St. Petersburg, Russia
- 135 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States of America
- 136 ^(a) Laboratório de Instrumentação e Física Experimental de Partículas – LIP; ^(b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; ^(c) Departamento de Física, Universidade de Coimbra, Coimbra; ^(d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; ^(e) Departamento de Física, Universidade do Minho, Braga; ^(f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada, Spain; ^(g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
- 137 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
- 138 Czech Technical University in Prague, Prague, Czech Republic
- 139 Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic
- 140 State Research Center Institute for High Energy Physics, NRC KI, Protvino, Russia
- 141 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- 142 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
- 143 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States of America
- 144 ^(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; ^(b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- 145 Department of Physics, University of Washington, Seattle, WA, United States of America
- 146 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- 147 Department of Physics, Shinshu University, Nagano, Japan
- 148 Department Physik, Universität Siegen, Siegen, Germany
- 149 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
- 150 SLAC National Accelerator Laboratory, Stanford, CA, United States of America
- 151 Physics Department, Royal Institute of Technology, Stockholm, Sweden
- 152 Departments of Physics and Astronomy, Stony Brook University, Stony Brook, NY, United States of America
- 153 Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
- 154 School of Physics, University of Sydney, Sydney, Australia
- 155 Institute of Physics, Academia Sinica, Taipei, Taiwan
- 156 ^(a) E. Andronikashvili Institute of Physics, Iv. Javakishvili Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
- 157 Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel
- 158 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
- 159 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
- 160 International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan
- 161 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
- 162 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
- 163 Tomsk State University, Tomsk, Russia
- 164 Department of Physics, University of Toronto, Toronto, ON, Canada
- 165 ^(a) TRIUMF, Vancouver, BC; ^(b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
- 166 Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
- 167 Department of Physics and Astronomy, Tufts University, Medford, MA, United States of America
- 168 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States of America
- 169 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
- 170 Department of Physics, University of Illinois, Urbana, IL, United States of America
- 171 Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia – CSIC, Valencia, Spain
- 172 Department of Physics, University of British Columbia, Vancouver, BC, Canada
- 173 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
- 174 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany
- 175 Department of Physics, University of Warwick, Coventry, United Kingdom
- 176 Waseda University, Tokyo, Japan
- 177 Department of Particle Physics, Weizmann Institute of Science, Rehovot, Israel
- 178 Department of Physics, University of Wisconsin, Madison, WI, United States of America
- 179 Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany

¹⁸⁰ Department of Physics, Yale University, New Haven, CT, United States of America

¹⁸¹ Yerevan Physics Institute, Yerevan, Armenia

- ^a Also at Borough of Manhattan Community College, City University of New York, NY; United States of America.
- ^b Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town; South Africa.
- ^c Also at CERN, Geneva; Switzerland.
- ^d Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
- ^e Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
- ^f Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.
- ^g Also at Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); Spain.
- ^h Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah; United Arab Emirates.
- ⁱ Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.
- ^j Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America.
- ^k Also at Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom.
- ^l Also at Department of Physics, California State University, Fresno CA; United States of America.
- ^m Also at Department of Physics, California State University, Sacramento CA; United States of America.
- ⁿ Also at Department of Physics, King's College London, London; United Kingdom.
- ^o Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg; Russia.
- ^p Also at Department of Physics, Stanford University; United States of America.
- ^q Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.
- ^r Also at Department of Physics, University of Michigan, Ann Arbor MI; United States of America.
- ^s Also at Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.
- ^t Also at Giresun University, Faculty of Engineering, Giresun; Turkey.
- ^u Also at Graduate School of Science, Osaka University, Osaka; Japan.
- ^v Also at Hellenic Open University, Patras; Greece.
- ^w Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; Romania.
- ^x Also at II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.
- ^y Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain.
- ^z Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.
- ^{aa} Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen; Netherlands.
- ^{ab} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary.
- ^{ac} Also at Institute of Particle Physics (IPP); Canada.
- ^{ad} Also at Institute of Physics, Academia Sinica, Taipei; Taiwan.
- ^{ae} Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
- ^{af} Also at Institute of Theoretical Physics, Ilija State University, Tbilisi; Georgia.
- ^{ag} Also at Istanbul University, Dept. of Physics, Istanbul; Turkey.
- ^{ah} Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.
- ^{ai} Also at Louisiana Tech University, Ruston LA; United States of America.
- ^{aj} Also at Manhattan College, New York NY; United States of America.
- ^{ak} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.
- ^{al} Also at National Research Nuclear University MEPhI, Moscow; Russia.
- ^{am} Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.
- ^{an} Also at School of Physics, Sun Yat-sen University, Guangzhou; China.
- ^{ao} Also at The City College of New York, New York NY; United States of America.
- ^{ap} Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China.
- ^{aq} Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.
- ^{ar} Also at TRIUMF, Vancouver BC; Canada.
- ^{as} Also at Università di Napoli Parthenope, Napoli; Italy.
- * Deceased.