

SPECTROSCOPY OF LOW-LYING STATES IN  $^{140}\text{Sm}^*$ 

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Electromagnetic transition strengths and spectroscopic quadrupole moments for  $^{140}\text{Sm}$  were measured by means of multi-step Coulomb excitation with radioactive beam at the ISOLDE facility at CERN. A complementary experiment was performed at the Heavy Ion Laboratory in Warsaw to assign spins for non-yrast states using the angular correlation technique. Based on the new experimental data previous spin assignments need to be revised.

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## 1. Introduction

The open-shell nuclei with  $Z > 50$  and  $N < 82$  are known to have some of the largest ground-state deformations in the nuclear chart [1]. The shapes of the nuclei in this region are expected to be prolate, except for a small island of nuclei with  $Z > 62$  and  $N \approx 78$ , which are predicted to be oblate [2]. Nuclei near  $^{140}\text{Sm}$  are, therefore, expected to be located in a transitional region between deformed and spherical shapes (as a function of neutron number) and between prolate and oblate shapes (as a function of proton number), as well as shape coexistence may be expected to occur. Indeed, the rotational bands built on the isomeric  $10^+$  states with  $\pi(h_{11/2})^2$  and  $\nu(h_{11/2})^{-2}$  configuration are interpreted as prolate and oblate structures, respectively [3]. Furthermore, an excited state at 990 keV was tentatively assigned as  $(0^+)$  in  $^{140}\text{Sm}$  [4], which could be interpreted as a sign for shape coexistence. The measurement of spectroscopic quadrupole moments and transition strengths for low-lying states represents a sensitive test for theoretical predictions in this region. Electromagnetic transition rates were unknown due to the difficulty of measuring lifetimes below long-lived isomeric states. To determine the transition rates between low-lying states and measure spectroscopic quadrupole moments, a low-energy Coulomb excitation experiment with a radioactive  $^{140}\text{Sm}$  beam was performed.

## 2. Coulomb excitation experiment

A quasi-pure  $^{140}\text{Sm}$  beam with an average intensity of  $2 \times 10^5$  particles per second and energy 2.85 MeV per nucleon was produced at ISOLDE using resonant laser ionization and scattered on a  $^{94}\text{Mo}$  target. The MINIBALL (HPGe) spectrometer [5] coupled to an annular DSSSD array was used for  $\gamma$ -particle coincidence detection. The DSSSD was placed 25.2 mm from the target and covered the angular range from  $19.7^\circ$  to  $58.4^\circ$  in the laboratory frame. Three excited states in  $^{140}\text{Sm}$  were populated during the experiment: the  $2^+$  and  $4^+$  states of the ground-state band and the tentatively assigned  $(0^+)$  state at 990 keV excitation energy. Figure 1 shows the  $\gamma$ -ray spectrum taken in coincidence with recoiling  $^{94}\text{Mo}$  target nuclei detected in the silicon detector. The spectrum was background subtracted and Doppler corrected for the velocity of the scattered  $^{140}\text{Sm}$  projectile. The excitation of the  $2^+$  state at 871 keV in  $^{94}\text{Mo}$  appears as a broad structure. The yield from the target excitation is obtained after applying a proper Doppler correction. Since the matrix elements for  $^{94}\text{Mo}$  are well known, the yield from the target excitation can be used to normalize the yield from the projectile excitation.

For the analysis, the data was divided into five subsets corresponding to different ranges of scattering angles, and a set of matrix elements was obtained by using a least-square fitting procedure to reproduce the experi-

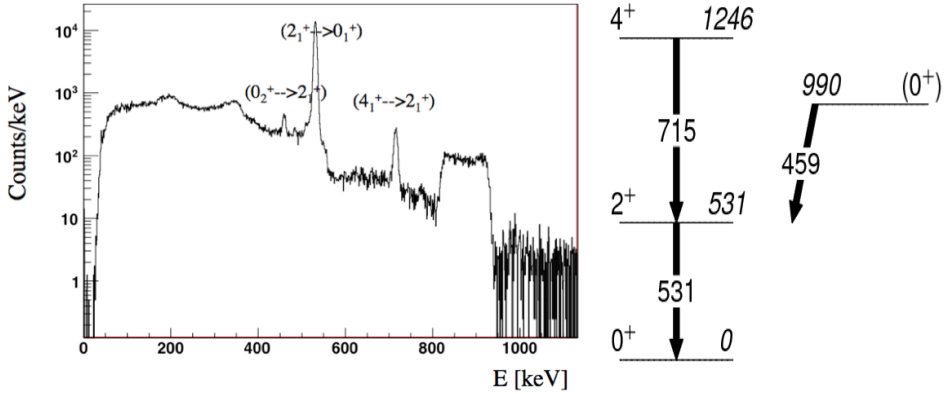


Fig. 1. Background subtracted  $\gamma$ -ray spectrum, Doppler corrected for the velocity of the scattered  $^{140}\text{Sm}$  projectile, and partial level scheme [6].

mentally observed  $\gamma$ -ray yields with the coupled channels codes GOSIA and GOSIA2 [7]. In this way, the  $B(E2; 2_1^+ \rightarrow 0_1^+)$  and  $B(E2; 4_1^+ \rightarrow 2_1^+)$  values could be obtained with high precision. The spectroscopic quadrupole moment of the  $2^+$  state was found to be small and negative. Although the error bars are large, the result suggests a small prolate deformation. The sensitivity to the quadrupole moment was enhanced by using the lifetime of the first excited  $2^+$  state, which was measured recently in a complementary experiment at the Heavy Ion Laboratory in Warsaw with the recoil distance Doppler shift technique using Bucharest–Köln Plunger [8]. The GOSIA analysis resulted in a very large  $B(E2; 0_2^+ \rightarrow 2_1^+)$  value of more than 200 Weisskopf units, which is difficult to understand. A new experiment was performed at the Heavy Ion Laboratory (Warsaw) to further investigate the nature of the non-yrast states and, in particular, to verify the spin assignments, which in the case of the 990 keV state was based on the non-observation of a direct transition to the ground state [4].

### 3. Angular correlation experiment

Excited states in  $^{140}\text{Sm}$  were populated following the  $^{140}\text{Eu} \rightarrow ^{140}\text{Sm}$  and  $^{140}\text{Gd} \rightarrow ^{140}\text{Eu} \rightarrow ^{140}\text{Sm}$   $\beta$ -decays. The  $^{140}\text{Eu}$  and  $^{140}\text{Gd}$  nuclei were produced in the reactions  $^{112}\text{Cd}(^{32}\text{S}, p3n)$  and  $^{112}\text{Cd}(^{32}\text{S}, 4n)$ , respectively, at a beam energy of 155 MeV. The recoils were stopped with an Au foil in the center of the EAGLE array [9], which comprised 12 escape-suppressed HPGe detectors. The macrostructure of the beam from the U200P cyclotron with 2 ms beam on and 4 ms beam off was used to suppress prompt  $\gamma$  rays and only select decay events during the beam-off periods. The detectors of the EAGLE

array were placed at five relative angles of 42(2), 70(2), 110(2), 138(2) and 180(2) degrees. This allowed the analysis of angular correlations for cascades originating from non-yrast states, where a coincidence gate was placed on the 531 keV  $2_1^+ \rightarrow 0_1^+$  transition. Angular correlation coefficients  $A_{22}$  and  $A_{44}$  were fitted to the observed intensities at different relative angles for the 459–531 keV and the 1068–531 keV cascades. The resulting coefficients are shown in Fig. 2 together with the expected values for  $0^+ \rightarrow 2^+ \rightarrow 0^+$  and  $2^+ \rightarrow 2^+ \rightarrow 0^+$  cascades. In the latter case, the angular correlation coefficients depend on the M1/E2 mixing ratio  $\delta$  of the  $2^+ \rightarrow 2^+$  transition.

The results of the angular correlation analysis clearly show the  $2^+$  nature of the state at 990 keV, contrary to the previous assignment as a  $0^+$  state. The coefficients are consistent with a mostly pure E2 character of the 459 keV transition. Moreover, the results show that the state at 1599 keV is indeed a  $0^+$  state, contrary to the previous  $2^+$  assignment. The resulting new level scheme is also shown in Fig. 2.

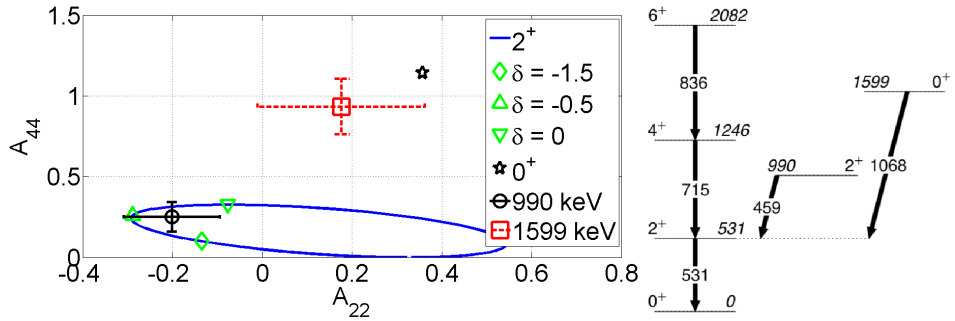


Fig. 2. Angular correlation coefficients for the  $\gamma$ -cascade from the 990 keV (black circle) and from the 1599 keV state (red square). The closed curve shows expected values for a  $2^+ \rightarrow 2^+ \rightarrow 0^+$  cascade for different M1/E2 mixing ratios  $\delta$ . The symbols (gray/green) on the black curve indicate the angular correlation coefficients for specific values of  $\delta$ . The star indicates the expected coefficients for a  $0^+ \rightarrow 2^+ \rightarrow 0^+$  cascade.

#### 4. Conclusions

The information obtained from the angular correlation measurement necessitates a re-analysis of the Coulomb excitation data with the correct spin assignments. The finding that the state at 990 keV has spin-parity  $2^+$  requires that the  $\langle 2_2^+ || M1 || 2_1^+ \rangle$ ,  $\langle 2_2^+ || E2 || 2_2^+ \rangle$  and  $\langle 2_2^+ || E2 || 0_1^+ \rangle$  matrix elements are included in the analysis. Preliminary results of the new GOSIA analysis suggest that the  $B(E2; 4_1^+ \rightarrow 2_1^+)$  and  $B(E2; 2_1^+ \rightarrow 0_1^+)$  values are insensitive to the new spin assignment. The  $B(E2; 2_2^+ \rightarrow 2_1^+)$  value corresponding to

the 459 keV transition, on the other hand, changes significantly toward a smaller value. The  $B(E2; 2_2^+ \rightarrow 0_1^+)$  value is found to be very small, which is in agreement with the non-observation of the transition in the present and previous experiments. Including the M1/E2 mixing ratio from the angular correlation experiment as known spectroscopic data in the Coulomb excitation analysis does not yield more precise matrix elements due to its relatively large uncertainty. The analysis of the Coulomb excitation data is presently being finalized and results will be presented and compared to theoretical calculations in a forthcoming publication.

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#### REFERENCES

- [1] P. Moeller *et al.*, *At. Data. Nucl. Data Tables* **59**, 185 (1995).
- [2] J.-P. Delaroche *et al.*, *Phys. Rev.* **C81**, 014303 (2010).
- [3] M.A. Cardona *et al.*, *Phys. Rev.* **C44**, 891 (1991).
- [4] R.B. Firestone *et al.*, *Phys. Rev.* **C43**, 1066 (1991).
- [5] N. Warr *et al.*, *Eur. Phys. J.* **A49**, 40 (2013).
- [6] N. Nica, *Nucl. Data Sheets* **108**, 1287 (2007).
- [7] T. Czosnyka, D. Cline, C.Y. Wu, *Bull. Am. Phys. Soc.* **28**, 745 (1983).
- [8] F.L. Bello Garrote *et al.*, to be published.
- [9] J. Mierzejewski *et al.*, *Nucl. Instrum. Methods* **A659**, 84 (2011).