

Pathwise regularization of the stochastic heat equation with multiplicative noise through irregular perturbation

Rémi Catellier^{1,a} and Fabian A. Harang^{2,3,b}

¹Université Côte d'Azur, CNRS, LJAD, France, ^aremi.catellier@univ-cotedazur,fr
 ²Department of Economics, BI Norwegian Business School, 0442, Oslo, Norway
 ³Department of Mathematics, University of Oslo, P.O. box 1053, Blindern, 0316, OSLO, Norway, ^bfabian.a.harang@bi.no

Received 11 January 2021; revised 18 January 2022; accepted 10 July 2022

Abstract. Existence and uniqueness of solutions to the stochastic heat equation with multiplicative spatial noise is studied. In the spirit of pathwise regularization by noise, we show that a perturbation by a sufficiently irregular continuous path establish wellposedness of such equations, even when the drift and diffusion coefficients are given as generalized functions or distributions. In addition we prove regularity of the averaged field associated to a Lévy fractional stable motion, and use this as an example of a perturbation regularizing the multiplicative stochastic heat equation.

Résumé. L'existence et l'unicité des solutions de l'équation de la chaleur stochastique avec bruit spatial multiplicatif sont étudiées. Dans l'esprit de la régularisation par le bruit trajectorielle des équations différentielles ordinaires, nous montrons qu'une perturbation additive par des fonctions suffisamment irrégulières permet d'obtenir existence et unicité, et ce même lorsque les coefficients sont des fonctions généralisées. De plus, nous étudions la régularité du champ moyennisé associé aux processus fractionnaires de Lévy stables. Nous utilisons de tels processus comme exemples de perturbations régularisantes pour l'équation de la chaleur stochastique multiplicative.

MSC2020 subject classifications: Primary 60H50; 60H15; secondary 60L20

Keywords: Pathwise regularization by noise; Stochastic heat equation; Generalized parabolic Anderson model; Fractional Lévy processes

1. Introduction

The stochastic heat equation with multiplicative noise (mSHE), is given on the form

(1.1)
$$\partial_t u = \Delta u + b(u) + g(u)\xi, \quad u_0 \in \mathcal{C}^{\beta}.$$

where ξ is a space time noise on \mathbb{R}^d or \mathbb{T}^d and \mathcal{C}^β is the Besov Hölder space of \mathbb{R}^d or \mathbb{T}^d , and *b* and *g* are sufficiently smooth functions. This equation is a fundamental stochastic partial differential equation, and is applied for modelling in a diverse selection of natural sciences, ranging from chemistry and biology to physics. The existence and uniqueness of (1.1) is typically proven under the condition that both *b* and *g* are Lipschitz functions of linear growth (see e.g. [27]). When taking a pathwise approach to a solution theory, even more regularity of these non linear functions may be required (see for instance [4] where $g \in C^3$ is required and Appendix B for a proof a pathwise wellposedness in a simple context). Of course if $g \equiv 0$, then (1.1) is known as the (deterministic) non-linear heat equation, for which uniqueness fails in general under weaker conditions on *b* than Lipschitz and linear growth.

Motivated by this, a natural question to ask is if it is possible to prove existence and uniqueness of (1.2) under weaker conditions on *b* and *g*. Inspired by the well known regularization by noise phenomena in stochastic differential equations, one may think that the same principles of regularization would extend to the case of stochastic partial differential equations like (1.2). In this article, we aim at giving some insights into this question by investigating (1.2) in a fully pathwise manner under perturbation by a measurable time (only) dependent path. In fact, we will prove that a perturbation by a sufficiently irregular path yields wellposedness of the mSHE, even for distributional (generalized functions) coefficients *g* and *b*. In

the next section we give a more detailed description of the specific equation under consideration, and the techniques that we apply in order to prove this regularizing effect of the perturbation.

1.1. Methodology

Inspired by the theory of regularization by noise for ordinary or stochastic differential equations, we show that a suitably chosen measurable path $\omega : [0, T] \to \mathbb{R}^d$ provides a regularizing effect on (1.1), by considering the formal equation

(1.2)
$$\partial_t u = \Delta u + b(u) + g(u)\xi + \dot{\omega}_t, \quad u_0 \in \mathcal{C}^{\beta}$$

where ξ is a spatial (distributional) noise taking values in \mathbb{R}^d , and $\dot{\omega}_t$ is the distributional derivative of a continuous path ω . To this end, we formulate (1.2) in terms of the non-linear Young framework, developed in [8,15,17]. We extend this framework to the infinite dimensional setting adapted to Volterra type integrals appearing when considering the mild formulation of (1.2). The integration framework developed here is strongly based on the recently developed *Volterra sewing lemma* of [26], and does not require any semi-group property of the Volterra operator. The integral can therefore be applied to several different problems relating to infinite dimensional Volterra integration, and thus we believe that this construction is interesting in itself.

To motivate the methodology of the current paper, consider again (1.2) and set $\theta = u - w$ with $\theta_0 = u_0 \in C^{\beta}$ then formally θ solves the following integral equation

$$\theta_t = P_t \theta_0 + \int_0^t P_{t-s} b(\theta_s + \omega_s) \,\mathrm{d}s + \int_0^t P_{t-s} \xi g(\theta_s + \omega_s) \,\mathrm{d}s,$$

where *P* is the fundamental solution operator associated with the heat equation, and the product $P_t \theta_0$ is interpreted as spatial convolution. For simplicity, we will carry out most of our analysis with $b \equiv 0$, as this term is indeed easier to handle than the term with multiplicative noise. The equation we will then consider is given by

(1.3)
$$\theta_t = P_t \theta_0 + \int_0^t P_{t-s} \xi g(\theta_s + \omega_s) \, \mathrm{d}s$$

In Section 4.2 we provide a detailed description of how our results can easily be extended to include the drift term with distributional *b*, by simply appealing to the the analysis carried out for the purely multiplicative equation (1.3). Associated to the path ω and the distribution *g*, define the averaged distribution $T^{\omega}g:[0,T] \times \mathbb{R}^d \to \mathbb{R}^d$ by the mapping

(1.4)
$$(t,x) \mapsto \int_0^t g(x+\omega_s) \,\mathrm{d}s$$

There are several papers devoted to the study of the regularity of the mapping $(t, x) \mapsto T^{\omega}g(t, x)$ under certain assumptions on the distribution g and the noise ω (typically chosen to be a sample path of an irregular stochastic process), see e.g. [16,24]). Of course, assuming $g \in H^{\alpha}$ is a Sobolev distribution and interpreting the integral in (1.4) as a Bochner integral, it is not difficult to check that $(t, x) \mapsto T^{\omega}g$ is contained in $C_T^1 H^{\alpha}$. However, exploiting specific structures of the pathwise or probabilistic behavior of certain paths ω , one may hope to get better regularity of this mapping.

In particular, it was already observed in [8] and further elaborated on in [25] that if ω is a sample path of a fractional Brownian motion with Hurst index $H \in (0, 1)$, and g is a Sobolev distribution in H^{α} , then $T^{\omega}g \in C_T^{\gamma}H^{\alpha+\frac{1}{2H}}$ for some $\gamma > \frac{1}{2}$. Thus if H is very small, we see that the path ω provides a strong regularizing effect on the averaged field $T^{\omega}g$. In this article we will prove regularity of the averaged field associated to distributions g in some weighted Besov spaces and paths sample paths of a fractional Lévy process ω , extending the results from [24] to weighted spaces.

After proving that for certain paths ω the distribution $T^{\omega}g$ is in fact a regular function, we then consider (1.3) as a non-linear Young equation of the form

(1.5)
$$\theta_t = P_t \theta_0 + \int_0^t P_{t-s} \xi T_{\mathrm{ds}}^\omega g(\theta_s), \quad \theta_0 \in \mathcal{C}^\beta$$

Here, the integral is interpreted in an infinite dimensional non-linear Young–Volterra sense. That is, to suit our purpose, we extend the non-linear Young integral to an infinite dimensional setting as well as allowing for the action of a Volterra operator on the integrand. The integral is then constructed as the Banach valued element

$$\int_0^t P_{t-s} T^{\omega}_{\mathrm{d}s} g(\theta_s) := \lim_{|\mathcal{P}| \to 0} \sum_{[u,v] \in \mathcal{P}} P_{t-u} \xi T^{\omega}_{u,v} g(\theta_u),$$

where $T_{u,v}^{\omega}g := T_v^{\omega}g - T_u^{\omega}g$. We stress that in contrast to the non-linear Young integral used for example in [8,15,17,25], the above integral is truly an infinite dimensional object, and extra care must be taken when building it from the averaged function $T^{\omega}g$. Indeed, for each $t \ge 0$, $\theta_t \in C^{\beta}(\mathbb{R}^d; \mathbb{R})$ and so the function $T^{\omega}g$ is then lifted to be a functional on C^{β} . We show that this lift comes at the cost of an extra degree of assumed regularity on the averaged function $T^{\omega}g$. Furthermore, due to the assumption that $\xi \in C^{-\vartheta}$ for $\vartheta > 0$ (i.e. ξ is assumed to be truly distributional) we need to make use of the product in Besov space in order to make the product of $\xi T_{u,v}^{\omega}g(\theta_u)$ well defined.

Similar to the theory of rough paths, our analysis can be divided into two parts: (i) a deterministic (analytic) step and (ii) a probabilistic step. We give a short description of the two steps here:

(i) Let *E* be a separable Banach space. We develop an abstract framework of existence and uniqueness of Banach valued equations

(1.6)
$$\theta_t = p_t + \int_0^t S_{t-s} X_{ds}(\theta_s),$$

where *S* is a suitable (possibly singular) Volterra operator, and $X : [0, T] \times E \to E$ is a function which is $\frac{1}{2}$ + Hölder regular in time, and suitably regular in its spatial argument (to be specified later), and $p : [0, T] \to E$ is a sufficiently regular function. To this end, we use a simple extension of the Volterra sewing lemma developed in [26] to construct the non-linear Young–Volterra integral appearing in (1.6) as the following

$$\int_0^t S_{t-s} X_{\mathrm{d}s}(\theta_s) := \lim_{|\mathcal{P}| \to 0} \sum_{[u,v] \in \mathcal{P}[0,t]} S_{t-u} X_{u,v}(\theta_u),$$

where $\mathcal{P}[0, t]$ is a partition of [0, t] with mesh size $|\mathcal{P}|$ converging to zero.

- +

(ii) The second step is then to consider $\{\omega_t\}_{t\in[0,T]}$ to be a stochastic process on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$, and we need to show that the averaged function $T^{\omega}g$ is indeed a sufficiently regular function \mathbb{P} -a.s., even when g is a true distribution. This is done by probabilistic methods.

At last we relate the abstract function X from (2.21) to the averaged function T^{ω} , and then a combination of the previous two steps gives us existence and uniqueness of (1.5), and thus also (1.3), which then is used to make sense of (1.2) through the translation $u = v + \omega$.

1.2. Short overview of existing literature

In order to prove existence and uniqueness of (1.1), there are two main directions to follow: the classical probabilistic setting based on Itô type theory, or the pathwise approach based on rough paths or similar techniques. Using the first approach, one typically requires that *b* and *g* are Lipschitz and of linear growth (see e.g. [27]). Certain extensions of this regime for the nonlinear function *g* has been shown over the years both in the case of white and colored space-time noise ξ . For example, in [36] the authors proves pathwise uniqueness of (1.1) for spatial dimension $d \ge 1$ in the case when *g* is only γ -Hölder continuous, where the regularity parameter γ depends on covariance structure of the colored noise. The Hölder regularity requirement obtained in [36] was recently improved in [40]. The above results was extended to space-time white noise ξ for spatial dimension d = 1 in [35] where pathwise uniqueness for γ -Hölder continuous coefficients *g* with $\gamma > \frac{3}{4}$ was proved.

Similar results are proven in [34,37,43], where certain special choices of the nonlinear function g are considered. In particular in [34] the authors prove, using probabilistic arguments, that there exists a non-zero solution to the equation

(1.7)
$$\partial_t u = \Delta u + |u|^{\gamma} \xi + \psi, \quad u(0, \cdot) = 0,$$

where ξ is a space time white noise in dimension 1 + 1, ψ is a non-zero, smooth non-negative compactly supported function and $0 \le \gamma < \frac{3}{4}$. They also prove that when $\gamma > \frac{3}{4}$ uniqueness holds.

Note also that when ξ is a deterministic function and not a distribution, wellposedness is a well-know topic. In particular if ξ is a non-negative continuous and bounded function, Fujita and Watanabe [14] prove that Osgood condition on g are "nearly" necessary and sufficient to guaranty uniqueness. In particular when g is only Hölder continuous, one can not expect to have uniqueness. One can also consult [6] and the reference therein for further attempt in that direction.

When using pathwise techniques to solve (stochastic) nonlinear heat equation, as usual (see [10] for counterexamples in a rough path context), one typically needs to require even higher regularity on b and g in order guarantee existence and uniqueness (see e.g. [4] where three times differentiability is assumed and the Appendix B for a simple proof).

$$\partial_t u = \partial_x^2 u + b(u) + \partial_t \partial_x \omega, \quad (t, x) \in [0, T] \times [0, 1]$$

exists uniquely in a strong sense, even when b is only bounded and measurable and $\partial_t \partial_x \omega$ is understood as a white noise on $[0, T] \times [0, 1]$. Similar assumptions was considered in [7], but there the authors also proved path-by-path uniqueness in the sense of Davie [10].

These results was recently improved in [1], where the authors prove pathwise existence and uniqueness even in the case of distributional drift functions b, and in particular they study the skew stochastic heat equation. The techniques used there are based on the stochastic sewing lemma and may be considered of a similar nature to the ones that will be considered here. However, the equation under study is very different, since in the present paper we consider the case of multiplicative noise. Still, the above results prove that the addition of noise seems to give similar regularizing effects in the stochastic heat equation as is observed for regular ODEs.

1.3. Main results

Before presenting our main results, let us first give a definition of what we will call a solution to (1.2). Let us remind that the definition of admissible weight and Besov spaces is written in Appendix A.2 and A.1, the precise definition of non-linear Young–Volterra equation is given in Section 2.1 and the definition of averaged field is given in Section 3. The equations considered in the below results can either be considered on \mathbb{R} or \mathbb{T} . Indeed, all the following results are true in both settings. Furthermore, for the sake of using the space white noise, it might be instructive to think of the space as the torus \mathbb{T} . For notational simplicity we choose to write the result in the case of \mathbb{R} .

Definition 1. Let $\omega : [0, T] \to \mathbb{R}$ be a measurable path, and consider a $g \in S'$ such that the averaged field $T^{\omega}g \in C_T^{\gamma}C^{\kappa}(w)$ for some $\gamma > \frac{1}{2}$ and $\kappa \ge 3$ and an admissible weight $w : \mathbb{R}^d \to \mathbb{R}$. Suppose that $0 < \vartheta < \beta < 1$ and suppose that ξ is a spatial noise contained in $C^{-\vartheta}$. Take $\rho = \frac{\beta + \vartheta}{2}$ and suppose that $\gamma - \rho > 1 - \gamma$. Let $u_0 \in C^{\beta}$. We say that there exists a unique solution to the equation

(1.8)
$$u_t = P_t u_0 + \int_0^t P_{t-s} \xi g(u_s) \, \mathrm{d}s + \omega_t, \quad t \in [0, T]$$

if $u \in \omega + \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$ for any $1 - \gamma < \varsigma < \gamma - \rho$, and there exists a unique $\theta \in \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$ such that $u = \omega + \theta$ and θ solves the non-linear Young equation

(1.9)
$$\theta_t = P_t u_0 + \int_0^t P_{t-s} \xi T_{ds}^{\omega} g(\theta_s)$$

Here the integral is understood as a non-linear Young–Volterra integral, as constructed in Section 2.1 and the space \mathscr{C}_T^5 is defined in Definition 10.

Theorem 2. Consider parameters $\gamma > \frac{1}{2}$, $\kappa \ge 3$ and $\vartheta \in (0, 1)$. Let $\beta > 0$ be such that $\beta \in (\vartheta, 2 - \vartheta)$, and assume that $\gamma - \rho > 1 - \gamma$ where $\rho = \frac{\beta + \vartheta}{2}$. Let $\omega : [0, T] \to \mathbb{R}$ be a measurable path, and consider a distribution $g \in S'(\mathbb{R})$ such that the associated averaged field $T^{\omega}g \in C_T^{\gamma}C^{\kappa}(w)$ for some admissible weight function $w : \mathbb{R}^d \to \mathbb{R}_+$. Let ξ be a spatial noise contained in $C^{-\vartheta}$. Then there exists a time $\tau \in (0, T]$ such that there exists a unique solution u with initial data $u_0 \in C^{\beta}$ in the sense of Definition 1 to Equation (1.8) on the interval $[0, \tau]$. If w is globally bounded, then the solution is global, in the sense that a unique solution u to (1.8) is defined on all $[0, \tau]$ for any $\tau \in (0, T]$.

Remark 3. The condition that $T^{\omega}g \in C^{\gamma}C^{\kappa}(w)$ for some suitable weight w, in order to get local existence and uniqueness may at this stage seem very abstract. Recall that $T^{\omega}g$ denotes the average field associated to the distribution g and ω as formally defined through (1.4). The regularity of this object has been the subject of several recent articles under various conditions on g and the (typically stochastic) process ω , see e.g. [16,17,24]. However, in the current framework we also allow the averaged field to be contained in certain weighted Besov spaces, which opens up for new types of distributions g. We illustrate this in Examples 7 and 8. Note however that if $\omega \equiv 0$ then $T_t^{\omega}g(x) = g(x)t$, and thus the regularity requirement in Theorem 2 would simply be that the function g is at least three times (weighted) differentiable. Our results are of course sub optimal in this case, as the non-linear Young framework that is used in the current article typically comes at a trade-off between spatial regularity assumptions on the non-linearity and time regularity, see e.g. [17]. In the case of zero noise, the time regularity of the averaged field is always 1 (i.e. differentiable), while the requirement in Theorem 2 reflects a time regularity of the averaged field close to $\frac{1}{2}$, which is the typical time regularity when seeking maximum regularizing effect from a regularizing path ω .

The unique solution found in Theorem 2 can be interpreted to be a "physical" one, in the sense that it is stable under approximations. We summarize this in the following corollary.

Corollary 4. Under the assumptions of Theorem 2, if $\{g_n\}_{n\in\mathbb{N}}$ is a sequence of smooth functions converging to $g \in S'(\mathbb{R})$ such that $T^{\omega}g_n \to T^{\omega}g \in C_T^{\gamma}C^{\kappa}(w)$, then the corresponding sequence of solutions $\{u_n\}_{n\in\mathbb{N}} = \{\omega + \theta^n\}_{n\in\mathbb{N}}$ to equation (1.9) converge to $u = w + \theta$ in the sense that $\theta^n \to \theta$ in $\mathcal{C}_{\tau}^{\varsigma}C^{\beta}$ for any $1 - \gamma < \varsigma < \gamma - \rho$.

In applications, one is typically interested in the regularizing effects provided by specific sample paths of stochastic processes ω . Although the class of regularizing paths is already well developed, we show here as an example the regularizing effect of measurable sample paths of fractional Lévy processes (see Section 5).

Theorem 5. Let $\alpha \in (0, 2]$. Let $H \in (0, 1) \cap \{\alpha^{-1}\}$. Let L^H be a Fractional Lévy Process with Hurst parameter H built from a symmetric α -stable Lévy process as defined in Section 5. Let $\vartheta \in (0, 1)$ and let $\kappa > 3 - \frac{1-\vartheta}{2H}$. There exists $\varepsilon > 0$ small enough such that for all $\xi \in C^{-\vartheta}$ and all $g \in C^{\kappa}(w)$ where w is an admissible weight, almost surely for all $u_0 \in C^{\vartheta + \varepsilon}$ there exists a unique local solution to the mSHE in the sense of Definition 1. If w is bounded, then the solution exists globally.

One can specify the previous theorem by letting ξ to be the space white noise:

Corollary 6. Let d = 1 and let ξ be a space white noise on the Torus \mathbb{T} . Let $\alpha \in (0, 2]$, let $H \in (0, 1) \cap \{\alpha^{-1}\}$ and let $\kappa > 3 - \frac{1}{4H}$. Let w be an admissible weight. There exists $\varepsilon > 0$ small enough such that for all $g \in C^{\kappa}(w)$, almost surely for all $u_0 \in C^{\frac{1}{2}+\varepsilon}$ there exists a unique solution (in the sense of Definition 1) to the mSHE

 $\mathrm{d} u = \Delta u \,\mathrm{d} t + g(u)\xi \,\mathrm{d} t + \mathrm{d} L_t^H, \quad u(0,\cdot) = u_0,$

where L^H is a linear fractional stable motion.

In particular taking $H < \frac{1}{4}$ allows us to take $\kappa < 2$ and go beyond the classical theory using Bony estimates for the product of distribution in Besov spaces. When $H < \frac{1}{8}$ one can deal with non Lipschitz-continuous g and when $H < \frac{1}{12}$, one can deal with distributional field g.

The proofs of the above theorems and corollary can be found in Sections 4 and 5. Let us illustrate Theorem 5 and Corollary 6 with the following example.

Example 7. For $\gamma > 0$, consider the function $g : \mathbb{T} \to \mathbb{T}$ given by $g(x) = |x|^{\kappa}$ for $\kappa \in (0, 1)$. From Proposition 44 it is readily checked that $g \in \mathcal{C}^{\kappa}(w_{\kappa})$, where $w_{\kappa} := (1 + |\cdot|^2)^{\frac{\kappa}{2}}$, and $\mathcal{C}^{\kappa}(w_{\kappa})$ denotes the weighted Hölder space defined in Appendix A.2. Referring to Corollary 6, suppose L^H is a fractional Levy process with $H < \frac{1}{4(3-\kappa)}$. Then there exists a unique local solution to

(1.10)
$$du = \Delta u \, dt + |u|^{\kappa} \xi \, dt + dL_t^H, \quad u(0, \cdot) = u_0 \in \mathcal{C}^{\frac{1}{2}+}$$

in the sense of Definition 1. In particular, if $H < \frac{1}{12}$, there exists a unique solution for any $\kappa \in (0, 1)$. Although not directly comparable with the results in [34] about equation (1.7) due to the fact that ξ in our case is only a spatial noise and not time dependent, it is an interesting observation that the irregularity of the noise provides regularizing effects yielding uniqueness for any $\kappa > 0$ even though the white noise is acting in a multiplicative way on the non-linearity.

Example 8. In the case when $\kappa < 0$ in $g(x) = |x|^{\kappa}$, the results in the previous example may be extended by interpreting g as a homogeneous distribution in $S'(\mathbb{T})$. Indeed, in [23, Appendix A.3] it is proven that $g \in B_{p,\infty}^{\kappa+\frac{d}{p}}$ (note that in this case we do not need the weight in the Besov space). In particular, $g \in C^{\kappa}$, and thus again if $H < \frac{1}{4(3-\kappa)}$ there exists a unique (global) solution to (1.10).

1.4. Outline of the paper

The paper is structured as follows: In Section 2.1 we extend the concept of non–linear Young integration to the infinite dimensional setting, including a Volterra operator, which in later sections will play the role of the inverse Laplacian. We also give a result on existence and uniqueness of abstract equations in Banach spaces.

In Section 3 we give a short overview on the concept of averaged fields and their properties. As this topic is by now well studied in the literature, we only give here the necessary details and provide several references for further information. We also show how the standard averaged field can be viewed as an operator on certain Besov function spaces.

In Section 4 we formulate the multiplicative stochastic heat equation in the non-linear Young–Volterra integration framework, using the concept of averaged fields. We prove existence and uniqueness of these equations, as well as Theorem 2 and Corollary 4. In Section 5 we investigate closer the regularity of averaged fields associated with sample paths of fractional Levy processes, and prove Theorem 5.

At last, in Section 6 we give a short reflection on the main results of the article and provide some thoughts on future extensions of our results.

For the sake of self-containedness, we have included an appendix with preliminaries on weighted Besov spaces, and some results regarding the regularity/singularity of the inverse Laplacian acting on Besov distributions. In addition we give a short proof for existence and uniqueness of (1.2) in the case of twice differentiable g.

1.5. Notation

For $\beta \in \mathbb{R}$ and $d \ge 1$, we define the Hölder Besov space

$$\mathcal{C}^{\beta} = B^{\beta}_{\infty,\infty}(\mathbb{R}^d;\mathbb{R})$$

endowed with its usual norm built upon Paley-Littlewood blocks and denoted by $\|\cdot\|_{\mathcal{C}^{\beta}}$ (see Appendix A.2 and especially Proposition 44 for more on weighted Besov spaces). Note that when working with a weight w and weighted spaces, we denote the spaces $\mathcal{C}^{\beta}(w)$ and the norm $\|\cdot\|_{\mathcal{C}^{\beta}(w)}$. For T > 0 and $\varsigma \in (0, 1)$ and E a Banach space with norm $\|\cdot\|_{E}$ and for $f : [0, T] \mapsto E$,

$$[f]_{\varsigma;E} = \sup_{s \neq t} \frac{\|f_t - f_s\|_E}{|t - s|^{\varsigma}} < \infty,$$

and for $\varsigma > 0$ we define

$$\|f\|_{\varsigma;E} := \sum_{k=0}^{\lceil \varsigma \rceil} \|f_0^{(k)}\|_E + \left[f^{(\lfloor \varsigma \rfloor)}\right]_{\varsigma - \lfloor \varsigma \rfloor;E},$$

where for an integer $k \in \mathbb{N}$, $f^{(k)}$ denotes the k'th derivative of f, and for $\zeta \in \mathbb{R}$, $\lfloor \zeta \rfloor$ denotes the largest integer smaller than ζ . Finally we define

$$\mathcal{C}_{T}^{\varsigma} E = \mathcal{C}^{\varsigma} \big([0, T]; E \big) := \big\{ f : [0, T] \to E : \| f \|_{\varsigma; E} < \infty \big\}.$$

Whenever the underlying space *E* is clear from the context we will use the short hand notation $[f]_{\varsigma}$ etc. to denote the the Hölder semi norm. We will frequently use the increment notation $f_{s,t} := f_t - f_s$. We write $f \leq g$ if there exists a constant C > 0 such that $f \leq Cg$. Furthermore to stress that the constant *C* depends on a parameter *p* we write $l \leq_p g$. We write $f \simeq g$ is $f \leq g$ and $g \leq f$. For $s \leq t$, we denote by $\mathcal{P}([s,t])$ a partition of the interval [s,t]. We write $|\mathcal{P}| = \max_{k \in \{0,...,n-1\}} |t_{k+1} - t_k|$. The Fourier transform of a function *f* will be denoted by \hat{f} with inverse transform denoted by $\mathcal{F}^{-1}(f)$.

2. Non-linear Young–Volterra theory in Banach spaces

To motivate the coming analysis of how parabolic equations through non-linear Young–Volterra theory, we will first recall some basic properties of the fundamental solution of the α -fractional heat equation. For $\alpha \in (0, 2]$, let $t \mapsto P_t^{\frac{\alpha}{2}}$ be the fundamental solution associated to the α -fractional heat equation

$$\partial_t P^{\frac{\alpha}{2}} = -(-\Delta)^{\frac{\alpha}{2}} P^{\frac{\alpha}{2}}, \quad P_0^{\frac{\alpha}{2}} = \delta_0.$$

For a function $y : \mathbb{R}^d \to \mathbb{R}$ let $P_t^{\frac{\alpha}{2}} y$ denotes the convolution between the fundamental solution at time $t \ge 0$ and y. Note that for $\alpha = 2$, we obtain the classical heat equation, and $P_t := P_t^{-1}$ then denotes the convolution operator with the Gaussian kernel.

Remark 9. We consider here the α -fractional heat equation, as the estimates outlined below will prove useful for our subsequent analysis of the regularizing effects of α -stable Fractional motions, but for most part of our analysis we will consider $\alpha = 2$, such that $(-\Delta)^{\frac{\alpha}{2}} = \Delta$ corresponding to the classical Laplace operator.

Towards a pathwise analysis of stochastic parabolic equations, one encounter the problem that for a function $y \in C^{\kappa}(\mathbb{R}^d)$ with $\kappa \ge 0$, the mapping $t \mapsto P_t^{\frac{\alpha}{2}} y$ is smooth in time everywhere except when approaching 0 where it is only continuous. In particular, from standard (fractional) heat kernel estimates (see Corollary 50 in the Appendix) we know that for $s \le t \in [0, T]$, any $\theta \in [0, 1]$ and $\rho \in (0, \alpha]$ the following inequality holds

(2.1)
$$\left\| \left(P_t^{\frac{\alpha}{2}} - P_s^{\frac{\alpha}{2}} \right) y \right\|_{\mathcal{C}^{\kappa+\alpha\rho}} \lesssim \|y\|_{\mathcal{C}^{\kappa}} |t-s|^{\theta} s^{-\theta-\rho}.$$

Note the special case when $\rho = 0$. Then P_t is a linear operator on C^{κ} which is θ -Hölder continuous on an interval $[\varepsilon, T] \subset [0, T]$ for any $\varepsilon > 0$ and $\theta \in [0, 1]$, but might only be continuous when approaching the point t = 0. We therefore need to extend the concept of Hölder spaces in order to take into account this type of loss of regularity near the origin.

Definition 10. Let *E* be a Banach space. We define the space of continuous paths on (0, T] which are Hölder of order $\zeta \in (0, 1)$ on (0, T] and only continuous at zero in the following way

$$\mathscr{C}_T^{\varsigma} E := \left\{ y : [0, T] \to E | [y]_{\varsigma} < \infty \right\}$$

where we define the semi-norm

$$[y]_{\varsigma} := \sup_{s \le t \in [0,T]; \zeta \in [0,\varsigma]} \frac{|y_{s,t}|_E}{|t-s|^{\zeta} s^{-\zeta}}.$$

The space $\mathscr{C}_T^{\varsigma} E$ is a Banach space when equipped with the norm $||y||_{\varsigma} := |y_0|_E + [y]_{\varsigma}$.

Singular Hölder type spaces introduced above has recently been extensively studied in [5]. The space introduced above can be seen as a special case of the more general spaces considered there. We use the convention of taking supremum over $\theta \in [0, \varsigma]$, as this will make computations in subsequent sections simpler. It is well known that if a function $y \in C_T^{\gamma} E$ for some $\gamma \in [0, 1)$, then $y \in C_T^{\theta} E$ for all $\theta \in [0, \gamma]$ (note that $\theta = 0$ implies that y is bounded, which is always true for Hölder continuous functions on bounded domains). Thus it follows that also

$$\sup_{s \le t \in [0,T]; \theta \in [0,\gamma]} \frac{|y_{s,t}|_E}{|t-s|^{\theta}} < \infty.$$

Remark 11. It is readily seen that $\mathscr{C}_T^{\varsigma} E$ consists of all functions which are Hölder continuous on (0, T], but is only continuous in the point {0}. Note therefore that the following inclusion $\mathscr{C}_T^{\varsigma} E \subset \mathscr{C}_T^{\varsigma} E$ holds. Indeed, for any $s, t \in [0, T]$ and $\theta \in [0, \varsigma]$

$$\frac{|y_{s,t}|_E}{|t-s|^{\theta}} = \frac{s^{-\theta}|y_{s,t}|_E}{s^{-\theta}|t-s|^{\theta}}$$

and thus in particular, for $|t - s| \le 1$

$$\sup_{\theta \in [0,\varsigma]} \frac{|y_{s,t}|_E}{s^{-\theta}|t-s|^{\theta}} \le \sup_{\theta \in [0,\varsigma]} T^{\theta} \frac{|y_{s,t}|_E}{|t-s|^{\theta}} \le \frac{|y_{s,t}|_E}{|t-s|^{\varsigma}}.$$

Now taking supremum on both sides over $s < t \in [0, T] \subset [0, 1]$, we obtain the claimed result for intervals $[0, T] \subset [0, 1]$. Extending to intervals of length greater than 1 follows by standard procedures.

Remark 12. It may be instructive for the reader to keep in mind that in subsequent sections, we will take the Banach space *E* to be the Besov-Hölder space $C^{\beta}(\mathbb{R}^d)$ or $C^{\beta}(\mathbb{T}^d)$ for some $\beta \in \mathbb{R}$ and $d \ge 1$.

We will throughout this section work with general Volterra type operators satisfying certain regularity assumptions. We therefore give the following working hypothesis.

Hypothesis 13. For each $t \in [0, T]$, let $S_t \in \mathcal{L}(E)$ be a linear operator on E satisfying the following three regularity conditions for some $\rho > 0$ and any $\theta, \theta' \in [0, 1]$ and any $0 \le s \le t \le \tau' \le \tau \le T$

$$|S_t u|_E \lesssim t^{-\rho} |u|_E$$

(ii)
$$|(S_t - S_s)u|_E \lesssim (t - s)^{\theta} s^{-\theta - \rho} |u|_E$$

(iii)
$$\left| \left((S_{\tau-t} - S_{\tau-s}) - (S_{\tau'-t} - S_{\tau'-s}) \right) u \right|_E \lesssim (\tau - \tau')^{\theta'} (t-s)^{\theta} (\tau'-t)^{-\theta - \theta' - \rho} |u|_E$$

We then say that $t \mapsto S_t$ is a ρ -singular operator.

2.1. Non-linear Young-Volterra integration

(i)

We are now ready to construct a non-linear Young–Volterra integral in Banach spaces. If $X : [0, T] \times E \to E$ is a smooth function in time, and $(P_t)_{t \in [0,T]}$ is a (nice) linear operator of E and if $(y_t)_{t \in [0,T]}$ is a continuous path from [0, T] to E itself, it is quite standard to consider integrals of the following form:

$$\int_0^t P_{t-r} \dot{X}_r(y_r) \,\mathrm{d}r.$$

The aim of this part is to extend the notion of integral for non-smooth drivers X, and to solve integral equations using this extension of the integral; The notion of ρ -singular operator will be useful in the following.

Lemma 14. Consider parameters $\gamma > \frac{1}{2}$, $0 \le \rho \le \gamma$ and $0 < \varsigma < \gamma - \rho$ and assume $\varsigma + \gamma > 1$. Let *E* be a Banach space and let $\gamma \in \mathscr{C}_T^{\varsigma} E$. Suppose that $X : [0, T] \times E \to E$ satisfies for any $x, y \in E$ and $s \le t \in [0, T]$

(i)
$$|X_{s,t}(x)|_E \lesssim H(|x|_E)|t-s|^{\gamma}$$

(ii) $|X_{s,t}(x) - X_{s,t}(y)|_E \lesssim H(|x|_E \vee |y|_E)|x-y|_E|t-s|^{\gamma}$

where *H* is a positive locally bounded function on \mathbb{R}_+ . Let $(S_t)_{t \in [0,T]} \in \mathcal{L}(E)$ be a ρ -singular linear operator on *E*, satisfying Hypothesis 13. We then define the non-linear Young–Volterra integral by

(2.2)
$$\Theta(y)_t := \lim_{\substack{\mathcal{P} \in \mathcal{P}([0,t]) \\ |\mathcal{P}| \to 0}} \sum_{\substack{[u,v] \in \mathcal{P}}} S_{t-u} X_{u,v}(y_u)$$

The integration map Θ is a continuous non-linear operator from $\mathscr{C}_T^5 E \to \mathscr{C}_T^5 E$, and the following inequality holds

(2.3)
$$\left|\Theta(y)_t - \Theta(y)_s\right|_E \lesssim \sup_{0 \le z \le \|y\|_{\infty}} H(z) \left(1 + [y]_{\varsigma}\right) (t-s)^{\varsigma} T^{\gamma-\rho+\varsigma},$$

Furthermore, for a linear operator $A \in \mathcal{L}(E)$, the following commutative property holds

$$A\Theta(y)_t = \lim_{|\mathcal{P}| \to} \sum_{[u,v] \in \mathcal{P}} AS_{t-u} X_{u,v}(y_u).$$

Proof. Let us first assume that for any $0 \le s \le t \le \tau' \le \tau \le T$ the following operator is well-defined:

$$\Theta_s^t(y)_\tau := \lim_{\substack{\mathcal{P} \in \mathcal{P}([s,t]) \\ |\mathcal{P}| \to 0}} \sum_{[u,v] \in \mathcal{P}} S_{\tau-u} X_{u,v}(y_u).$$

Note that in this setting we have

$$\Theta(\mathbf{y})_t = \Theta_0^t(\mathbf{y})_t$$

and the increment satisfies

$$\Theta(\mathbf{y})_{s,t} = \Theta_s^t(\mathbf{y})_t + \Theta_0^s(\mathbf{y})_{s,t}$$

Hence, in order to have the bound (2.3), it is enough to have a bound on $\Theta_s^t(y)_{\tau}$ and on $\Theta_s^t(y)_{\tau',\tau}$.

To this end, we will begin to show the existence of the integrals $\Theta_s^t(y)_{\tau}$ and $\Theta_s^t(y)_{\tau',\tau}$ together with the suitable bounds. Both these terms are constructed in the same way, however since

(2.4)
$$\Theta_s^t(y)_{\tau',\tau} = \lim_{\substack{\mathcal{P} \in \mathcal{P}([s,t]) \\ |\mathcal{P}| \to 0}} \sum_{\substack{[u,v] \in \mathcal{P}}} (S_{\tau-u} - S_{\tau'-u}) X_{u,v}(y_u),$$

this term is a bit more involved as it has an increment of the kernel *P* in the summand. We will therefore show existence as well as a suitable bound for this term and leave the specifics of the first term as a simple exercise for the reader. Everything will be proven in a similar manner as the sewing lemma from the theory of rough paths. More specifically, the recently developed Volterra sewing lemma from [26] provides the correct techniques to this specific setting. The uniqueness, and additivity (i.e. that $\Theta(y)_{s,t} = \Theta_s^t(y)_t + \Theta_0^s(y)_{s,t}$) of the mapping follows directly from the standard arguments given for example in [26] or [13, Lem. 4.2].

Consider now a dyadic partition \mathcal{P}^n of [s, t] defined iteratively such that $\mathcal{P}^0 = \{[s, t]\}$, and for $n \ge 0$

$$\mathcal{P}^{n+1} := \bigcup_{[u,v]\in\mathcal{P}^n} \{[u,m],[m,v]\},\$$

where $m := \frac{u+v}{2}$. It follows that \mathcal{P}^n consists of 2^n sub-intervals [u, v], each of length $2^{-n}|t-s|$. Define the approximating sum

(2.5)
$$\mathcal{I}_n := \sum_{[u,v]\in\mathcal{P}^n} (S_{\tau-u} - S_{\tau'-u}) X_{u,v}(y_u),$$

and observe that for $n \in \mathbb{N}$, we have

(2.6)
$$\mathcal{I}_{n+1} - \mathcal{I}_n = -\sum_{[u,v]\in\mathcal{P}^n} \delta_m \Big[(S_{\tau-u} - S_{\tau'-u}) X_{u,v}(y_u) \Big],$$

where $m = \frac{u+v}{2}$ and for a two variable function f, we use that $\delta_m f_{u,v} := f_{u,v} - f_{u,m} - f_{m,v}$. By elementary algebraic manipulations we see that

(2.7)
$$\delta_m [(S_{\tau-u} - S_{t-u})X_{u,v}(y_u)] = (S_{\tau-u} - S_{\tau'-u})(X_{m,v}(y_u) - X_{m,v}(y_m)) + (S_{\tau-u} - S_{\tau'-u} - S_{\tau-m} + S_{\tau'-m})X_{m,v}(y_u).$$

We first investigate the second term on the right hand side above. Invoking (iii) of Hypothesis 13 and invoking assumption (i) on *X*, we observe that for any $\theta, \theta' \in [0, 1]$

$$(S_{\tau-u} - S_{t-u} - S_{\tau-m} + S_{t-m})X_{m,v}(y_u)\big|_E$$

\$\lesssim H(|y_u|_E)|\tau - \tau'|^{\theta'}|\tau' - m|^{-\theta'-\rho-\theta}|m-u|^{\theta}|\tu-m|^{\gamma'}.

Let us now fixe $\theta' = \varsigma \in [0, \gamma - \rho)$, and choose $\theta \in [0, 1]$ such that $\gamma + \theta > 1$ and $\theta + \varsigma + \rho < 1$. Note that this is always possible due to the fact that $\gamma - \rho - \varsigma > 0$. Furthermore, we note that for any partition \mathcal{P} of [s, t] we have

(2.8)
$$\sum_{[u,v]\in\mathcal{P}} \left|\tau'-m\right|^{-\theta-\rho-\varsigma} \left|v-m\right| \lesssim \int_{s}^{t} \left|\tau'-r\right|^{-\theta-\rho-\varsigma} \mathrm{d}r \lesssim |t-s|^{1-\varsigma-\rho-\theta},$$

where we have used that m = (u + v)/2. From this, it follows that for any $\theta \in [0, \zeta]$ the following inequality holds

(2.9)
$$\sum_{\substack{[u,v]\in\mathcal{P}^n\\ \lesssim \sup_{0\leq z\leq \|y\|_{\infty}}} H(z) |\mathcal{P}^n|^{\gamma+\theta-1} |\tau-\tau'|^{\varsigma} |t-s|^{1-\theta-\rho-\varsigma}.$$

Let us now move on to the first term in (2.7). By invoking the bounds on *P* from (ii) of Hypothesis 13 and assumption (ii) on *X*, we observe that for any $\theta \ge 0$ and any $0 \le \zeta \le \zeta$,

$$\left| (S_{\tau-u} - S_{\tau'-u}) \left(X_{m,v}(y_u) - X_{m,v}(y_m) \right) \right|_E$$

$$\lesssim \left| \tau - \tau' \right|^{\theta} \left| \tau' - u \right|^{-\rho-\theta} |m-v|^{\gamma} |m-u|^{\zeta} u^{-\zeta} H \left(|y_u|_E \vee |y_m|_E \right) [y]_{\zeta}.$$

Similarly as shown in (2.9), we now take $\theta = \zeta = \zeta$ and we consider a sum over a partition \mathcal{P} over $[s, t] \subset [0, T]$, and see that since $\zeta + \gamma > 1$, and for $\rho + \zeta < 1$,

$$\sum_{[u,v]\in\mathcal{P}} \left|\tau-\tau'\right|^{\varsigma} \left|\tau'-m\right|^{-\rho-\varsigma} |m-v|^{\gamma} |m-u|^{\varsigma} u^{-\varsigma} \le |\mathcal{P}|^{\varsigma+\gamma-1} \left|\tau-\tau'\right|^{\varsigma} \int_{s}^{t} \left|\tau'-r\right|^{-\rho-\varsigma} r^{-\varsigma} dr$$

where again m = (u + v)/2. Furthermore, when $s < t < \tau'$, we have

$$\begin{split} \int_{s}^{t} |\tau' - r|^{-\rho - \varsigma} r^{-\varsigma} \, \mathrm{d}r &\leq \int_{s}^{t} |\tau' - r|^{-\rho - \varsigma} r^{-\varsigma} \, \mathrm{d}r \\ &\leq (t - s)^{1 - \rho - 2\varsigma} \int_{0}^{1} (1 - r)^{-(\rho + \varsigma)} r^{-\varsigma} \, \mathrm{d}r \\ &\lesssim (t - s)^{1 - \rho - 2\varsigma}. \end{split}$$

We therefore obtain when specifying $\theta = \zeta$,

(2.10)
$$\sum_{[u,v]\in\mathcal{P}^n} |(S_{\tau-u} - S_{\tau'-u})(X_{m,v}(y_u) - X_{m,v}(y_m))|_E \\ \lesssim |\mathcal{P}^n|^{\varsigma+\gamma-1} |\tau - \tau'|^{\varsigma} |t-s|^{1-\rho-2\varsigma} \sup_{0 \le z ||y||_{\infty}} H(z)[y]_{\varsigma}$$

Combining (2.9) and (2.10), and using that for $|\mathcal{P}^n| = 2^{-n}|t - s|$, it follows from (2.6) that

$$\left|\mathcal{I}_{n+1}(s,t) - \mathcal{I}_{n}(s,t)\right|_{E} \lesssim \sup_{0 \le z \le \|y\|_{\infty,[s,t]}} H(z) (1+[y]_{\varsigma}) 2^{-n(\gamma-(\rho+\varsigma))} (\tau-\tau')^{\varsigma} (t-s)^{\gamma-\rho-\varsigma},$$

For $m > n \in \mathbb{N}$ thanks to the triangle inequality, and the estimate above, we get

(2.11)
$$\left| \mathcal{I}_{m}(s,t) - \mathcal{I}_{n}(s,t) \right|_{E} \lesssim \sup_{0 \le z \le \|y\|_{\infty,[s,t]}} H(z) \left(1 + [y]_{\varsigma} \right) \left(\tau - \tau' \right)^{\varsigma} (t-s)^{\gamma - \rho - \varsigma} \psi_{n,m},$$

where $\psi_{n,m} = \sum_{i=n}^{m} 2^{-i(\gamma - (\rho + \varsigma))}$, and it follows that $\{\mathcal{I}_n\}_{n \in \mathbb{N}}$ is Cauchy in *E*. It follows that there exists a limit $\mathcal{I} = \lim_{n \to \infty} \mathcal{I}_n E$. Moreover, from (2.11) we find that the following inequality holds

(2.12)
$$\left| \mathcal{I} - (S_{\tau-s} - S_{\tau'-s}) X_{s,t}(y_s) \right|_E \lesssim \sup_{0 \le z \le \|y\|_{\infty}} H(z) \left(1 + [y]_{\varsigma} \right) \left(\tau - \tau' \right)^{\varsigma} (t-s)^{\gamma-\rho-\varsigma} \psi_{0,\infty}.$$

the proof that \mathcal{I} is equal to $\Theta_s^t(y)_{\tau,\tau'}$ defined in (2.4) (where in particular $\Theta(y)$ is defined independent of the partition \mathcal{P} of [s, t]) follows by standard arguments for the sewing lemma, see [26] in the Volterra case. Finally, we get for any $0 \le s \le t \le \tau' \le \tau \le T$,

(2.13)
$$\left| \Theta_{s}^{t}(y)_{\tau',\tau} - (S_{\tau-s} - S_{\tau'-s})X_{s,t}(y_{s}) \right|_{E} \lesssim \sup_{0 \le z \le \|y\|_{\infty}} H(z)T^{\gamma-\rho+\varsigma}|t-s|^{\varsigma}.$$

Using the same techniques, in order to prove that $\Theta_s^t(y)_t$ exists for all $0 \le s \le t \le \tau \le T$ one has to control

$$S_{\tau-m}(X_{m,\nu}(y_m) - X_{m,\nu}(y_u)) + (S_{\tau-u} - S_{\tau-m})X_{m,\nu}(y_u).$$

Performing exactly the same computation, we have

(2.14)
$$\left|\Theta_{s}^{t}(y)_{\tau}-S_{\tau-s}X_{s,t}(y_{s})\right|_{E}\lesssim \sup_{0\leq z\leq \|y\|_{\infty}}H(z)T^{\gamma-\varsigma+\rho}|t-s|^{\varsigma}\left(1+[h]_{\varsigma}\right).$$

Combining (2.13) and (2.14), and using the fact that

$$\Theta(y)_t - \Theta(y)_s = \Theta_{s,t}(y)_t + \Theta_0^s(y)_{s,t}$$

$$\left|\Theta(y)_{t,s} - S_{t-s}X_{s,t}(y_s) - (S_t - S_s)X_{0,s}(y_0)\right|_E \le \left|\Theta_s^t(y)_t - (S_t - S_s)X_{s,t}(y_s)\right|_E + \left|\Theta_0^s(y)_{s,t} - S_{t-s}\xi X_{0,s}(y_0)\right|_E$$

it is readily checked that the following inequality holds

(i)

(ii)

$$\left|\Theta(y)_{t} - \Theta(y)_{s} - S_{t-s}X_{s,t}(y_{s}) - (S_{t} - S_{s})X_{0,s}(y_{0})\right|_{E} \lesssim \sup_{0 \le z \le \|y\|_{\infty}} H(z) \left(1 + [y]_{\varsigma}\right) T^{\gamma - \rho + \varsigma} (t-s)^{\varsigma}.$$

Finally note that since $\zeta < \gamma - \rho$,

$$\left|S_{t-s}X_{s,t}(y_s)\right|_E \lesssim H\left(|y_s|\right)|t-s|^{\gamma-\rho} \lesssim H\left(|y_s|\right)T^{\gamma-\rho+\varsigma}|t-s|^{\varsigma},$$

and

$$\left| (S_t - S_s) X_{0,s}(y_0) \right|_E \lesssim H(|y_0|) s^{\gamma - \rho - \varsigma} |t - s|^{\varsigma}$$

For the last claim, if $A \in \mathcal{L}(E)$ is a linear operator, then one re-define \mathcal{I}_n in (2.5) to be given as

$$\mathcal{I}_n(A) := \sum_{[u,v]\in\mathcal{P}^n} A(S_{\tau-u} - S_{t-u}) X_{u,v}(y_u),$$

and by linearity of A we see that $\mathcal{I}_n(A) = A\mathcal{I}_n$. Taking the limits, using the above established inequalities, we find that $\lim_{n\to\infty} \|\mathcal{I}_n(A) - A\mathcal{I}_n\|_E = 0$.

With the construction of the non-linear Young Volterra integral, we will in later applications need certain stability estimates.

Proposition 15 (Stability of Θ). Let γ , ς , ρ be given as in of Lemma 14. Assume that for $i = 1, 2, X^i : [0, T] \times E \to E$ satisfies for any $x, y \in E$ and $s \leq t \in [0, T]$

$$\left|X_{s,t}^{i}(x)\right|_{E} + \left|\nabla X_{s,t}^{i}(x)\right|_{\mathcal{L}(E)} \lesssim H\left(|x|_{E}\right)|t-s|^{\gamma}$$

(2.15)

(iii)
$$\left|\nabla X_{s,t}^{i}(x) - \nabla X_{s,t}^{i}(y)\right|_{\mathcal{L}(E)} \lesssim H\left(|x|_{E} \vee |y|_{E}\right)|x-y|_{E}|t-s|^{\gamma},$$

 $\left|X_{s,t}^{i}(x) - X_{s,t}^{i}(y)\right|_{E} \lesssim H\left(|x|_{E} \vee |y|_{E}\right)|x - y|_{E}|t - s|^{\gamma}$

where *H* is a positive locally bounded function, and ∇ is understood as a linear operator on *E* in the Fréchet sense. Furthermore, suppose there exists a positive and locally bounded function $H_{X^1-X^2}$ such that

(2.16)
(i)
$$|X_{s,t}^{1}(x) - X_{s,t}^{2}(x)|_{E} \lesssim H_{X^{1}-X^{2}}(|x|_{E})|t-s|^{\gamma}$$

(ii) $|(X_{s,t}^{1} - X_{s,t}^{2})(x) - (X_{s,t}^{1} - X_{s,t}^{2})(y)|_{E} \lesssim H_{X^{1}-X^{2}}(|x|_{E} \vee |y|_{E})|x-y|_{E}|t-s|^{\gamma}$

Let Θ^1 denote the non-linear integral operator constructed in Lemma 14 with respect to X^1 , and similarly let Θ^2 denote the integral operator with respect to X^2 . Then for two paths $y, \tilde{y} \in \mathscr{C}_T^5 \mathcal{C}^{\beta+2\rho}$,

(2.17)
$$\left[\Theta^{1}(y^{1}) - \Theta^{2}(y^{2}) \right]_{\varsigma} \lesssim_{P} \left[\sup_{0 \le z \le \|y^{1}\| \lor \|y^{2}\|} H(z) \left(\left[y^{1} \right]_{\varsigma} + \left[y^{2} \right]_{\varsigma} \right) \left(\left| y^{1}_{0} - y^{2}_{0} \right|_{E} + \left[y^{1} - y^{2} \right]_{\varsigma} \right) \right. \\ \left. + \sup_{0 \le z \le \|y^{1}\|_{\infty} \lor \|y^{2}\|_{\infty}} H_{X^{1} - X^{2}}(z) \left[y^{1} \right]_{\varsigma} \lor \left[y^{2} \right]_{\varsigma} \right] T^{\gamma - \rho + \varsigma}.$$

Proof. Define the following two functions

$$\Theta_{s,t}^{1}(y^{1}) - \Theta_{s,t}^{2}(y^{2}) = (\Theta_{s,t}^{1}(y^{1}) - \Theta_{s,t}^{2}(y^{1})) + (\Theta_{s,t}^{2}(y^{1}) - \Theta_{s,t}^{2}(y^{2})) =: D_{X^{1},X^{2}}(s,t) + D_{y^{1},y^{2}}(s,t).$$

We treat D_{X^1,X^2} and D_{y^1,y^2} separately, and begin to consider D_{y^1,y^2} . Since X^i is differentiable and satisfies (i)–(iii) in (2.15), for i = 1, 2 we have

$$X_{s,t}^{i}(y_{s}^{1}) - X_{s,t}^{i}(y_{s}^{2}) = \mathcal{X}_{s,t}^{i}(y_{s}^{1}, y_{s}^{2})(y_{s}^{1} - y_{s}^{2}),$$

where $\mathcal{X}^i(y_s^1, y_s^2) := \int_0^1 \nabla X_{s,t}^i(qy_s^1 + (1-q)y_s^2) \, dq$. In order to prove (2.17), we proceed with the exact same strategy as outlined in the proof of Lemma 14. That is, we use the same proof as the proof of Lemma 14 to first prove appropriate bounds for $D_{y_1} v^2(s, t)$ and similarly for $D_{X_s} v^2(s, t)$ afterwards. To this end, changing the integrand in (2.5) so that

$$\mathcal{I}_{n}(s,t) := \sum_{[u,v] \in \mathcal{P}^{n}[s,t]} (S_{\tau-u} - S_{t-u}) \mathcal{X}_{u,v}^{i} (y_{u}^{1}, y_{u}^{2}) (y_{u}^{1} - y_{u}^{2})$$

we continue along the lines of the proof in Lemma 14 to show that \mathcal{I}^n is Cauchy. As the strategy of this proof is identical to that of Lemma 14 we will here only point out the important differences. By appealing to the condition (iii) in (2.15),we observe in particular that

$$\left| \mathcal{X}^{i}(y_{s}^{1}, y_{s}^{2}) - \mathcal{X}^{i}(y_{u}^{1}, y_{u}^{2}) \right|_{\mathcal{L}(E)} \lesssim |t - s|^{\gamma} \sup_{0 \le z \le \|y^{1}\|_{\infty} \lor \|y^{2}\|_{\infty}} H(z) \left(\left[y^{1} \right]_{\varsigma} + \left[y^{2} \right]_{\varsigma} \right).$$

Furthermore, it is readily checked that

$$|y_s^1 - y_s^2|_E \lesssim |y_0^1 - y_0^2|_E + [y - \tilde{y}]_{\varsigma} T^{\varsigma}.$$

Following along the lines of the proof of Lemma 14, one can then check that for $m > n \in \mathbb{N}$

(2.18)
$$\begin{aligned} & |\mathcal{I}_m(s,t) - \mathcal{I}_n(s,t)|_E \\ & \lesssim \sup_{0 \le z \le \|y^1\|_{\infty} \lor \|y^2\|_{\infty}} H(z) \big([y^1]_{\varsigma} + [y^2]_{\varsigma} \big) \big(|y_0^1 - y_0^2|_E + [y^1 - y^2]_{\varsigma} T^{\varsigma} \big) \big(\tau - \tau' \big)^{\varsigma} (t-s)^{\gamma - \rho - \varsigma} \psi_{n,m}, \end{aligned}$$

where $\psi_{n,m}$ is defined as below (2.11). With this inequality at hand, the remainder of the proof can be verified in a similar way as in the proof of Lemma 14, and we obtain from this lemma that

$$\|D_{y^{1},y^{2}}\|_{\mathscr{C}_{T}^{\varsigma}E} \lesssim C \sup_{0 \le z \le \|y^{1}\|_{\infty} \lor \|y^{2}\|_{\infty}} (z) ([y^{1}]_{\varsigma} + [y^{2}]_{\varsigma}) (\|y^{1}_{0} - y^{2}_{0}\|_{E} + [y^{1} - y^{2}]_{\varsigma} T^{\varsigma})$$

Next we move on to prove a similar bound of D_{X^1,X^2} as defined in (2.1). Set $Z = X^1 - X^2$. By (2.16) it follows that Z satisfies the conditions of Lemma 14, and then from (2.3) it follows that

$$[D_{X_1,X^2}]_{\varsigma} \lesssim \sup_{0 \le z \le \|y^1\|_{\infty} \lor \|y^2\|_{\infty}} H_{X^1 - X^2}(z) [y^1]_{\varsigma} \lor [y^2]_{\varsigma} T^{\gamma - \varsigma - \rho}.$$

2.2. Existence and uniqueness

We begin to prove local existence and uniqueness for an abstract type of equation with values in a Banach space. The equation in itself does not require the use of the non-linear Young–Volterra integral, and is formulated for general operators $\Theta : [0, T] \times \mathscr{C}_T^5 E \to \mathscr{C}_T^5 E$ satisfying certain regularity conditions. We will apply these results in later sections in combination with the non-linear Young–Volterra integral operator Θ created in the previous section, and thus the reader is welcome to already think of Θ as being a non-linear Young integral operator as constructed in Lemma 14.

Theorem 16 (Local existence and uniqueness). Let $\Theta : [0, T] \times \mathscr{C}_T^{\varsigma} E \to \mathscr{C}_T^{\varsigma} E$ be a function which satisfies for $y, \tilde{y} \in \mathscr{C}_T^{\varsigma} E$ and some $\epsilon > 0$

(2.19)
$$\begin{split} & \left[\Theta(y)\right]_{\varsigma} \leq C\left(\|y\|_{\infty}\right)\left(1+[y]_{\varsigma}\right)T^{\varepsilon} \\ & \left[\Theta(y)-\Theta(\tilde{y})\right]_{\varsigma} \leq C\left(\|y\|_{\infty} \vee \|\tilde{y}\|_{\infty}\right)\left([y]_{\varsigma}+[\tilde{y}]_{\varsigma}\right)\left(|y_{0}-\tilde{y}_{0}|_{E}+[y-\tilde{y}]_{\varsigma}\right)T^{\varepsilon} \end{split}$$

where C is a positive and increasing locally bounded function. Consider $p \in \mathscr{C}_T^{\varsigma} E$ and let $\tau > 0$ be such that

(2.20)
$$\tau \leq \left[4\left(1+[p]_{\varsigma}\right)C\left(1+|p_{0}|+[p]_{\varsigma}\right)\right]^{-\frac{1}{\varepsilon}}$$

Then there exists a unique solution to the equation

(2.21)
$$y_t = p_t + \Theta(y)_t, \quad p \in \mathscr{C}_T^{\varsigma} E$$

in $\mathcal{B}_{\tau}(p)$, where $\mathcal{B}_{\tau}(p)$ is a unit ball in $\mathscr{C}_{\tau}^{\varsigma} E$, centered at p.

Proof. To prove existence and uniqueness, we will apply a standard fixed point argument. Define the solution map $\Gamma_{\tau}: \mathscr{C}_{\tau}^{\varsigma} E \to \mathscr{C}_{\tau}^{\varsigma} E$ given by

$$\Gamma_{\tau}(\mathbf{y}) := \left\{ p_t + \Theta(\mathbf{y})_t | t \in [0, \tau] \right\}.$$

Since $\Theta : \mathscr{C}^{\varsigma}_{\tau} E \to \mathscr{C}^{\varsigma}_{\tau} E$ and $p \in \mathscr{C}^{\varsigma}_{\tau} E$ it follows that $\Gamma_{\tau}(y) \in \mathscr{C}^{\varsigma}_{\tau} E$. We now prove that the solution map Γ_{τ} is an invariant map and a contraction on a unit ball $\mathcal{B}_{\tau}(p) \subset \mathscr{C}^{\varsigma}_{\tau} E$ centered at $p \in \mathscr{C}^{\varsigma}_{\tau} E$. In particular, we define

$$\mathcal{B}_{\tau}(p) := \left\{ y \in \mathscr{C}_{\tau}^{5} | y_{t} = p_{t} + z_{t}, \text{ with } z \in \mathscr{C}_{\tau}^{5} E, z_{0} = 0, [y - p]_{5} \le 1 \right\}.$$

We begin with the invariance. From the first condition in (2.19), it is readily checked that for $y \in \mathcal{B}_{\tau}(P)$

(2.22)
$$\left[\Gamma_{\tau}(y) - p \right]_{\varsigma} \leq C \left(1 + |p_0| + [p]_{\varsigma} \right) \left(1 + [y]_{\varsigma} \right) \tau^{\varepsilon},$$

where we have used that for $y \in \mathcal{B}_{\tau}(P)$, $||y||_{\infty} \leq 1 + |p_0| + [p]_{\varsigma}$ Choosing a parameter $\tau_1 > 0$ such that

$$\tau_1 \le (2C(1+|p_0|+[p]_{\varsigma})^{-\frac{1}{\varepsilon}})$$

it follows that $\Gamma_{\tau_1}(\mathcal{B}_{\tau_1}(p)) \subset \mathcal{B}_{\tau_1}(p)$ and we say that Γ_{τ_1} leaves the ball $\mathcal{B}_{\tau_1}(p)$ invariant.

Next, we prove that Γ_{τ} is a contraction on $\mathcal{B}_{\tau}(p)$. From the second condition in (2.19), it follows that for two elements $y, \tilde{y} \in \mathcal{B}_{\tau}(p)$ we have

$$\left[\Gamma(\mathbf{y}) - \Gamma(\tilde{\mathbf{y}})\right]_{\varsigma} \leq 2\left(1 + [p]_{\varsigma}\right)C\left(1 + |p_0| + [p]_{\varsigma}\right)[\mathbf{y} - \tilde{\mathbf{y}}]_{\varsigma}\tau^{\varepsilon},$$

where we have used that $y_0 = \tilde{y}_0$, and

$$[y]_{\varsigma} \vee [\tilde{y}]_{\varsigma} \le 1 + [p]_{\varsigma}$$
 and $\|y\|_{\infty} \vee \|\tilde{y}\|_{\infty} \le 1 + |p_0| + [p]_{\varsigma}$.

Again, choosing a parameter $\tau_2 > 0$ such that

$$\tau_2 \le \left(4\left(1 + [p]_{\varsigma}\right)C\left(1 + [p_0] + [p]_{\varsigma}\right)\right)^{-\frac{1}{6}}$$

it follows that

$$\left[\Gamma(y) - \Gamma(\tilde{y})\right]_{\varsigma;\tau_2} \le \frac{1}{2} [y - \tilde{y}]_{\varsigma;\tau_2}.$$

Since $\tau_2 \leq \tau_1$, we conclude that the solution map Γ_{τ_2} both is an invariant map and a contraction on the unit ball $\mathcal{B}_{\tau_2}(p)$. It follows by Picard-Lindlöfs fixed point theorem that a unique solution to (2.21) exists in $\mathcal{B}_{\tau_2}(p)$.

The next theorem shows that if the locally bounded function C appearing in the conditions on Θ in (2.19) of Theorem 16, is uniformly bounded, then there exists a unique global solution to (2.21).

Theorem 17 (Global existence and uniqueness). Let $\Theta : [0, T] \times \mathscr{C}_T^{\varsigma} E \to \mathscr{C}_T^{\varsigma} E$ satisfy (2.19) for a positive, globally bounded function C, i.e. there exists a constant M > 0 such that $\sup_{x \in \mathbb{R}_+} C(x) \leq M$. Furthermore, suppose Θ is time-additive, in the sense that $\Theta_t = \Theta_s + \Theta_{s,t}$ for any $s \leq t \in [0, T]$. Then for any $p \in \mathscr{C}_T^{\varsigma} E$ there exists a unique solution $y \in \mathscr{C}_T^{\varsigma} E$ to the equation

$$y_t = p_t + \Theta(y)_t, \quad t \in [0, T].$$

Proof. By Theorem 16 we know that there exists a unique solution to (2.21) on an interval $[0, \tau]$, where τ satisfies (2.20), and *C* is replaced by the bounding constant *M*, i.e.

(2.23)
$$\tau \leq \left[4\left(1+\left[p\right]_{\varsigma}\right)M\right]^{-\frac{1}{\varsigma}}.$$

By a slight modification of the proof in Theorem 16 it is readily checked that the existence and uniqueness of

$$y_t = p_t + \Theta_{a,t}(y), \quad t \in [a, a + \tau],$$

holds on any interval $[a, a + \tau] \subset [0, T]$, i.e. the solution is constructed in $\mathcal{B}_{[a, a+\tau]}(p)$.

Now, we want iterate solutions to (2.21) to the domain [0, T], by "gluing together" solutions on the integrals $[0, \tau], [\tau, 2\tau] \dots \subset [0, T]$. Using the time-additivity property of Θ , note that for $t \in [\tau, 2\tau]$, we have

$$y_t = p_t + \Theta_t(y) = p_t + \Theta_a(y) + \Theta_{a,t}(y).$$

Thus, set $\tilde{p}_t = p_t + \Theta_a(y|_{[0,\tau]})$, where $y|_{[0,\tau]}$ denotes the solution to (2.21) restricted to $[0,\tau]$. Note that

$$[\tilde{p}]_{\varsigma} = [p]_{\varsigma}$$

since the Hölder seminorm is invariant to constants, and $t \mapsto \Theta_a(y|_{[0,\tau]})$ is constant. Therefore, there exists a unique solution to (2.21) in $\mathcal{B}_{[\tau,2\tau]}(\tilde{p})$ where τ is the same as in (2.23). We can repeat this to all intervals $[k\tau, (k+1)\tau] \subset [0, T]$. At last, invoking the scalability of Hölder norms (see [13, Exc. 4.24]), it follows that there exists a unique solution to (2.21) in $\mathscr{C}_T^{\varsigma} E$.

3. Averaged fields

We give here a quick overview of the concept of averaged fields and averaging operators. We begin with the following definition:

Definition 18. Let ω be a measurable path from [0, T] to \mathbb{R}^d , and let $g \in \mathcal{S}'(\mathbb{R}^d; \mathbb{R})$. We define the average of g against ω as the element of $C^0([0, T]; \mathcal{S}'(\mathbb{R}^d; \mathbb{R}))$ defined for all $s \leq t \in [0, T]$ and all test functions ϕ by

$$\langle \phi, T_{s,t}^{\omega} g \rangle = \int_{s}^{t} \langle \phi(\cdot - \omega_{r}), g \rangle \mathrm{d}r.$$

Introduced in the analysis of regularization by noise in [8], the concept of averaged fields and averaging operators is by now a well studied topic. For example, the recent analysis of Galeati and Gubinelli [17] provides a good overview of the analytic properties in the context of regularization by noise. See also [16,18,25] for further details on probabilistic and analytical aspects of averaged fields and averaging operators, and their connection to the concept of occupation measures. In the current article, we investigate these operators from an infinite-dimensional perspective in order to apply them in the context of (S)PDEs, and thus some extra considerations needs to be taken into account. In addition, we include in Section 5 a construction of the averaged field associated to a fractional Lévy process. The regularizing properties of Volterra-Lévy processes was recently investigated in [24], where the authors constructed the averaged field using the concept of local times. The construction given in the current article provides better regularity properties of the resulting averaged field than those obtained in [24], but comes at the cost of lost generality. More precisely, constructing the averaged field associated to a measurable stochastic process $\omega : \Omega \times [0, T] \to \mathbb{R}^d$ acting on a distribution $b \in \mathcal{S}'(\mathbb{R}^d)$, the exceptional set $\Omega' \subset \Omega$ on which the function $T^{\omega}g$ is a sufficiently regular field (say, Hölder in time and differentiable in space..) depends on $b \in \mathcal{S}'(\mathbb{R}^d)$, i.e. $\Omega' = \Omega'(b)$. Thus choosing one construction or the other depends on the problem at hand, and which properties are important to retain. We will not investigate these differences in more details in this article, and will view the analysis of the SPDE from a purely deterministic point of view. Let us give the following statement which can be viewed as a short summary of some of the results appearing in [17] and [25]:

Proposition 19. There exists a δ -Hölder continuous path $\omega : [0, T] \to \mathbb{R}$ such that for any given $g \in C^{\eta}$ with $\eta > 3 - \frac{1}{2\delta}$, the corresponding averaged field $T^{\omega}g$ is contained in $C_T^{\gamma}C^{\kappa}$ for some $\kappa \geq 3$ and $\gamma > \frac{1}{2}$. Moreover, there exists a continuous $\omega : [0, T] \to \mathbb{R}$ such that for any $g \in S'(\mathbb{R})$, the averaged field $T^{\omega}g$ is contained in $C_T^{\gamma}C^{\kappa}(w)$ for some weight $w : \mathbb{R} \to \mathbb{R}_+$ and any $\gamma \in (\frac{1}{2}, 1)$ and any $\kappa \in \mathbb{R}$.

Proof. The first statement can be seen as a simple version of [17, Thm. 1]. The second is proven in [25, Prop. 24]. \Box

Remark 20. In fact, the statement of [17, Thm. 1] is much stronger; given a $g \in C^{\eta}$ with $\eta \in \mathbb{R}$, then for almost all δ -Hölder continuous paths $\omega : [0, T] \to \mathbb{R}$ such that $\eta > 3 - \frac{1}{2\delta}$, the averaged field $T^{\omega}g \in C_T^{\gamma}C^3$ for some $\gamma > \frac{1}{2}$. Similarly, it is stated that almost all continuous paths ω are infinitely regularizing in the sense that for any $g \in C^{\eta}$ with $\eta \in \mathbb{R}$, $T^{\omega}g \in C^{\gamma}C^{\kappa}$ for any $\kappa \in \mathbb{R}$. The "almost surely" statement here is given through the concept of prevalence. We refrain from writing proposition 19 in the most general way in order to avoid going into details regarding the concept of prevalence here. We therefore refer the reader to [17] for more details on this result and this concept.

Remark 21. The statement in Proposition 19 can also be generalized to measurable paths $\omega : [0, T] \rightarrow \mathbb{R}$. Indeed, in [24] the authors show that there exists a class of measurable Volterra-Lévy processes, which provides a regularizing effect, similar to that of Gaussian processes, and a statement similar to that of Proposition 19 can be found there.

In all the following we will use notions of (weighted)-Besov space. For a recap on weighted Lebesgue and Besov spaces, as long as a recap on standard (fractional)-Schauder estimates see Appendix Sections A.1 and A.2.

We continue with some properties which will be useful in later analysis.

Proposition 22. Let w be an admissible weight, $\kappa \in \mathbb{R}$ and $1 \le p, q \le +\infty$. Let $g \in B_{p,q}^{\kappa}(w)$. For all $j \ge -1$ and all $0 \le s \le t \le T$,

$$\Delta_j T^{\omega}_{s,t} g = T^{\omega}_{s,t}(\Delta_j g),$$

where Δ_j denotes the standard Paley-Littlewood block (see Appendix A.2). In particular for all $\varepsilon, \delta > 0$, and for all $g \in B_{p,q}^{\kappa}(w)$, and for $S_k = \sum_{j=-1^k} \Delta_j$,

$$T_{s,t}^{\omega}(\mathcal{S}_k g) \xrightarrow[k \to \infty]{} T_{s,t}^{\omega} g \quad in \ \mathcal{S}' \quad and \quad in \ B_{p,q}^{\kappa-\varepsilon} \left(\langle \cdot \rangle^{\delta} w \right).$$

Suppose that g is a measurable locally bounded function, then $T_{s,t}^{\omega}g$ is also a measurable function and

$$T^{\omega}_{s,t}g(x) = \int_{s}^{t} g(x+\omega_{r}) \,\mathrm{d}r.$$

Proof. Let us remind that the topology on S is the one generated by the family of semi-norm

$$\mathcal{N}_{n}(\phi) = \sum_{|k|,|l| \le n} \sup_{x \in \mathbb{R}^{d}} \left| x^{k} \partial^{l} \phi(x) \right|,$$

where k and l are multi-indices. Furthermore, for any distribution $f \in S'$, there exists $n \ge 0$ and C > 0 such that

$$|\langle \phi, g \rangle| \leq C \mathcal{N}_n(\phi).$$

Let ω be a measurable path from [0, T] to \mathbb{R}^d , then since ϕ is bounded, there is a constant $C_1 > 0$ such that

$$\mathcal{N}_n(\phi(\cdot - w_r)) \leq C\mathcal{N}_n(\phi).$$

In particular

$$\Delta_j T_t^{\omega} g(x) = \left\langle K_j(x-\cdot), T_{s,t}^{\omega} f \right\rangle = \int_0^t \left\langle K_j(x+\omega_r) - \cdot \right\rangle, g \right\rangle \mathrm{d}r = \int_0^t \Delta_j g(x+\omega_r) \,\mathrm{d}r.$$

Thanks to the previous remark, we are allowed to perform any Fubini arguments, and for all $\phi \in S$,

$$\langle \phi, \Delta_j T_t^{\omega} g \rangle = \int_{\mathbb{R}^d} \phi(x) \int_0^t \Delta_j g(x + \omega_r) \, \mathrm{d}r \, \mathrm{d}x = \int_0^t \int_{\mathbb{R}^d} \phi(x - \omega_r) \Delta_j g(x) \, \mathrm{d}x \, \mathrm{d}r,$$

which ends the proof by using properties of weighted-Besov spaces.

For the purpose of this section, we will fix a distribution $g \in \mathcal{S}'(\mathbb{R}^d)$ and a measurable path $\omega : [0, T] \to \mathbb{R}^d$ with the property that there exists an averaged field $T^{\omega}g : [0, T] \times \mathbb{R}^d \to \mathbb{R}^d$ defined as in Definition 18 which is Hölder continuous in time and three times locally differentiable in space. More specifically, we will assume that $T^{\omega}g \in \mathcal{C}_T^{\gamma}\mathcal{C}^{\kappa}(w)$ for some admissible weight function $w : \mathbb{R}^d \to \mathbb{R}_+ \setminus \{0\}$.

Our first goal is to show how the averaged field $T^w g$ can be seen as a function from $[0, T] \times C^\beta \to C^\beta$ for some $\beta \ge 0$.

Proposition 23. Let w be an admissible weight (see Section A.1). For a measurable path $\omega : [0, T] \to \mathbb{R}^d$ and $g \in S'(\mathbb{R}^d)$, suppose $T^{\omega}g \in C_T^{\gamma}C^{\kappa}(w)$ for some $\gamma > \frac{1}{2}$ and $\kappa \ge 3$. Then for all $x, y \in C^{\beta}$ with $\beta \in [0, 1)$, we have that

$$\begin{split} \left\| T_{s,t}^{\omega} g(x) \right\|_{\mathcal{C}^{\beta}} & \leq \| \nabla T_{s,t}^{\omega} g(x) \|_{\mathcal{C}^{\beta}} \leq \sup_{|z| \leq \|x\|_{\mathcal{C}^{\beta}}} w^{-1}(z) \| T^{\omega} g \|_{\mathcal{C}^{\gamma}_{T} \mathcal{C}^{\kappa}(w)} |t-s|^{\gamma} \\ & \left\| T_{s,t}^{\omega} g(x) - T_{s,t}^{\omega} g(y) \right\|_{\mathcal{C}^{\beta}} \leq \sup_{|z| \leq \|x\|_{\mathcal{C}^{\beta}} \vee \|y\|_{\mathcal{C}^{\beta}}} w^{-1}(z) \| T^{\omega} g \|_{\mathcal{C}^{\gamma}_{T} \mathcal{C}^{\kappa}(w)} \| x-y\|_{\mathcal{C}^{\beta}} |t-s|^{\gamma} \\ & \left\| \nabla T_{s,t}^{\omega} g(x) - \nabla T_{s,t}^{\omega} g(y) \right\|_{\mathcal{C}^{\beta}} \leq \sup_{|z| \leq \|x\|_{\mathcal{C}^{\beta}} \vee \|y\|_{\mathcal{C}^{\beta}}} w^{-1}(z) \| T^{\omega} g \|_{\mathcal{C}^{\gamma}_{T} \mathcal{C}^{\kappa}(w)} \| x-y\|_{\mathcal{C}^{\beta}} |t-s|^{\gamma} . \end{split}$$

Proof. For $x \in C^{\beta}$ note that for any $\xi \in \mathbb{R}^d$

$$\left|T_{s,t}^{\omega}g(x(\xi))\right| \le w^{-1}(x(\xi))\left|w(x(\xi))T_{s,t}^{\omega}g(x(\xi))\right| \le \sup_{|z|\le \|x\|_{\mathcal{C}^{\beta}}} w^{-1}(z)\left\|T^{\omega}g\right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}}.$$

By similar computations using that $T^{\omega}g$ is (weighted) differentiable in space, it is straightforward to verify that

$$\left\|T_{s,t}^{\omega}g(x)\right\|_{\mathcal{C}^{\beta}} \vee \left\|\nabla T_{s,t}^{\omega}g(x)\right\|_{\mathcal{C}^{\beta}} \leq \sup_{|z| \leq \|x\|_{\mathcal{C}^{\beta}}} w^{-1}(z) \left\|T^{\omega}g\right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)} |t-s|^{\gamma}.$$

Now consider $x, y \in C^{\beta}$. Again using the differentiability of $T^{\omega}g$, and elementary rules of calculus we observe that for any $\xi \in \mathbb{R}^d$

$$\left|T_{s,t}^{\omega}g(x(\xi)) - T_{s,t}^{\omega}g(y(\xi))\right| \leq \sup_{|z| \leq ||x||_{\mathcal{C}^{\beta}} \vee ||y||_{\mathcal{C}^{\beta}}} w^{-1}(z) \left\|T^{\omega}g\right\|_{\mathcal{C}_{T}^{\gamma}\mathcal{C}^{\kappa}(w)} \|x - y\|_{\mathcal{C}^{\beta}} |t - s|^{\gamma},$$

where we have used that $|x(\xi) - y(\xi)| \le ||x - y||_{\mathcal{C}^{\beta}}$ for all $\xi \in \mathbb{R}^d$. At last, for $\xi, \xi' \in \mathbb{R}^d$, observe that

$$T_{s,t}^{\omega}g(x(\xi)) - T_{s,t}^{\omega}g(y(\xi)) - T_{s,t}^{\omega}g(x(\xi')) + T_{s,t}^{\omega}g(y(\xi'))$$

$$= \int_{0}^{1} \nabla T_{s,t}^{\omega}g(l(x(\xi) - y(\xi)) + y(\xi))(x(\xi) - x(\xi') - y(\xi) + y(\xi')) dl$$

$$+ \int_{0}^{1} \int_{0}^{1} D^{2}T_{s,t}^{\omega}g(\Lambda(l,l'))(l(x(\xi) - x(\xi')) + (1 - l)(y(\xi) - y(\xi'))) \otimes (x(\xi') - y(\xi')) dl' dl,$$

with

$$\Lambda(l,l') = l'(l(x(z) - x(z')) + (1 - l)(y(z) - y(z'))) + lx(z') + (1 - l)y(z').$$

Since $T^{\omega}g$ is twice (weighted) differentiable, it follows that

$$\left\|T_{s,t}^{\omega}g(x)-T_{s,t}^{\omega}g(y)\right\|_{\mathcal{C}^{\beta}}\leq \sup_{|z|\leq \|x\|_{\mathcal{C}^{\beta}}\vee\|y\|_{\mathcal{C}^{\beta}}}w^{-1}(z)\left\|T^{\omega}g\right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)}\|x-y\|_{\mathcal{C}^{\beta}}|t-s|^{\gamma}.$$

A similar argument for $\nabla T^{\omega}g$, using that $\kappa \geq 3$, reveals that

$$\left\|\nabla T^{\omega}_{s,t}g(x) - \nabla T^{\omega}_{s,t}g(y)\right\|_{\mathcal{C}^{\beta}} \leq \sup_{|z| \leq \|x\|_{\mathcal{C}^{\beta}} \vee \|y\|_{\mathcal{C}^{\beta}}} w^{-1}(z) \left\|T^{\omega}g\right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)} \|x-y\|_{\mathcal{C}^{\beta}} |t-s|^{\gamma},$$

which concludes this proof.

Remark 24. In order to apply Proposition 15 to $X_{s,t}(x) := T^{\omega}g(x)$ for $x \in C^{\beta}$ we need to relate the spatial derivative considered in Proposition 23 with the Fréchet derivative. A simple point-wise Taylor expansion reveals that for two functions $x, y \in C^{\beta}$ and any $\xi \in \mathbb{R}^d$

(3.1)
$$T^{\omega}g(x(\xi) + y(\xi)) - T^{\omega}g(x(\xi)) = \nabla \left[T^{\omega}g(x(\xi))\right]y(\xi) + O(y(\xi)),$$

where the spatial derivative is in the sense of Proposition 15 and $\nabla[T^{\omega}g(x(\xi))] = \nabla T^{\omega}g(x(\xi)) \otimes \nabla x(\xi)$. Thus, the Fréchet derivative of $T^{\omega}g$ is given by the mapping $x \mapsto \nabla T^{\omega}g(x) \otimes \nabla x$.

Thus, given that $T^{\omega}g$ satisfies the assumptions of Proposition 23, it is also Fréchet differentiable, and by linearity, satisfies the assumptions of 15.

4. Existence and uniqueness of the mSHE

We are now ready to formulate the multiplicative stochastic heat equation using the abstract non-linear Young integral and the averaged field $T^{\omega}g$ introduced in the previous section.

4.1. Standard multiplicative stochastic heat equation with additive noise

Recall that the mSHE with additive (time)-noise is given in its mild form by

(4.1)
$$u_t = P_t u_0 + \int_0^t P_{t-s} \xi g(u_s) \,\mathrm{d}s + \omega_t, \quad u_0 \in \mathcal{C}^\beta$$

for some $\beta > 0$, where $\xi \in C^{-\vartheta}$ for some appropriately chosen $\vartheta \in (0, 1)$ satisfying $0 \le \vartheta < \beta < 2 - \vartheta$. Here *P* denotes the standard heat semi-group acting on functions *u* through convolution, and $\omega : [0, T] \to \mathbb{R}$ is a measurable path. We will see in this section that *g* can be chosen to be distributional given that ω is sufficiently irregular.

Similarly as what is done for pathwise regularization by noise for SDEs (e.g. [8]), we begin to consider (4.1) with a smooth function g, and we set $\theta = u - \omega$, and then study the integral equation

(4.2)
$$\theta_t = P_t u_0 + \int_0^t P_{t-s} \xi g(\theta_s + \omega_s) \,\mathrm{d}s.$$

The integral term can be written in the form of a non-linear Young–Volterra integral. To motivate this, we begin with the following observation: Define the linear operator $S_t = P_t \xi$, whose action is defined by

(4.3)
$$S_t f := \int_{\mathbb{R}} P_t(x - y)\xi(y)f(y) \, \mathrm{d}y.$$

For $\beta > \vartheta$ the classical Schauder estimates for the heat equation tells us that

$$\|S_t f\|_{\mathcal{C}^{\beta}} = \|S_t f\|_{\mathcal{C}^{-\vartheta+2\frac{\beta+\vartheta}{2}}} \lesssim t^{-\frac{\beta+\vartheta}{2}} \|\xi f\|_{\mathcal{C}^{-\vartheta}} \lesssim t^{-\frac{\beta+\vartheta}{2}} \|\xi\|_{\mathcal{C}^{-\vartheta}} \|f\|_{\beta},$$

where in the last estimate we have used that the product between the distribution $\xi \in C^{-\vartheta}$ and the function $f \in C^{\beta}$ since $\beta > \vartheta$. Thus we can view *S* as a bounded linear operator from C^{β} to itself for any $\beta > \vartheta$. This motivates the next proposition:

Proposition 25. If $\vartheta \ge 0$ and $\beta > \vartheta$, then the operator *S* defined in (4.3) is a linear operator on C^{β} and satisfies *Hypothesis* 13 with singularity $\rho = \frac{\beta + \vartheta}{2}$.

Proof. This follows from the properties of the heat kernel proven in Corollary 50 in Section A.3, and the fact that the para-product ξf satisfies the bound $\|\xi f\|_{C^{-\vartheta}} \le \|\xi\|_{C^{-\vartheta}} \|f\|_{C^{\beta}}$ due to the assumption that $\beta > \vartheta$. See [2] for more details on the para-product in Besov spaces.

The integral in (4.2) can therefore be written as

$$\int_0^t P_{t-s}\xi g(\theta_s+\omega_s)\,\mathrm{d}s = \int_0^t S_{t-s}g(\theta_s+\omega_s)\,\mathrm{d}s.$$

Furthermore, the classical Volterra integral on the right hand side can be written in terms of the averaged field $T^{\omega}g$, in the sense that

$$\int_0^t S_{t-s}g(\theta_s+\omega_s)\,\mathrm{d}s = \int_0^t S_{t-s}T_{\mathrm{d}s}^\omega g(\theta_s),$$

where for a partition \mathcal{P} of [0, t] with infinitesimal mesh, we define

$$\int_0^t S_{t-s} T^{\omega}_{\mathrm{ds}} g(\theta_s) := \lim_{|\mathcal{P}| \to 0} \sum_{[u,v] \in \mathcal{P}} S_{t-u} T^{\omega}_{u,v} g(\theta_u).$$

It is not difficult to see (and we will rigorously prove it later) that when g is smooth, the above definition agrees with the classical Riemann definition of the integral. However, the advantage with this formulation of the integral is that $T^{\omega}g$ might still make sense as a function when g is only a distribution, and thus we will see that this formulation allows for an extension of the concept of integration to distributional coefficients g.

Proposition 26. Consider a measurable path $\omega : [0, T] \to \mathbb{R}$ and $g \in \mathcal{S}'(\mathbb{R})$, and suppose $T^{\omega}g \in \mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)$ for some $\gamma > \frac{1}{2}$ and $\kappa \ge 3$. For some $0 < \beta < 1$, suppose $S : [0, T] \to \mathcal{L}(\mathcal{C}^{\beta})$ satisfies Hypothesis 13 for some $0 \le \rho < 1$ such that $\gamma - \rho > 1 - \gamma$. Take $\gamma - \rho > \varsigma > 1 - \gamma$ Then for any $\gamma \in \mathcal{C}^{\varsigma}_{T}\mathcal{C}^{\beta}$, the integral

(4.4)
$$\Theta(y)_t := \int_0^t S_{t-s} T^{\omega}_{\mathrm{ds}} g(y_s) := \lim_{|\mathcal{P}| \to 0} \sum_{[u,v] \in \mathcal{P}} S_{t-u} T^{\omega}_{u,v} g(y_u)$$

exists as a non-linear Young-Volterra integral according to Lemma 14.

Proof. Let $E := C^{\beta}$. Since $T^{\omega}g \in C_T^{\gamma}C^{\kappa}(w)$ for some $\gamma > \frac{1}{2}$ and $\kappa \ge 3$, it follows from Proposition 23 that $T^{\omega}g$ can be seen as a function from $[0, T] \times E$ to E that satisfies (i)-(ii) in Lemma 14 with $H(x) := \sup_{|z| \le x} w^{-1}(z) ||T^{\omega}g||_{C_T^{\gamma}C^{\kappa}(w)}$. Lemma 14 then implies that (4.4) exists, and satisfies the bound in (2.3).

As an immediate consequence of the above proposition and Proposition 15, we obtain the following Proposition:

Proposition 27. Under the same assumptions as in Proposition 26, for two paths $x, y \in \mathscr{C}^{5}_{T} \mathcal{C}^{\beta}$ we have that

$$\left[\Theta(y) - \Theta(x)\right]_{\varsigma} \lesssim_{P,\xi} \left\{ \sup_{|z| \le \|x\|_{\infty} \lor \|y\|_{\infty}} w^{-1}(z) \left\| T^{\omega}g \right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}} \right\} \left(\|x_{0} - y_{0}\|_{\mathcal{C}^{\beta}} + [x - y]_{\varsigma} \right).$$

Moreover, let $\{g_n\}_{n\in\mathbb{N}}$ be a sequence of smooth functions converging to $g \in \mathcal{S}'(\mathbb{R})$ such that $T^{\omega}g_n \to T^{\omega}g \in \mathcal{C}^{\gamma}_T \mathcal{C}^{\kappa}(w)$. Then the sequence of integral operators $\{\Theta_n\}_{n\in\mathbb{N}}$ built from $T^{\omega}g_n$ converge to Θ , in the sense that for any $y \in \mathscr{C}^{\varsigma}_T \mathcal{C}^{\beta}$

$$\lim_{n \to \infty} \left\| \Theta_n(y) - \Theta(y) \right\|_{\mathcal{C}_T^{\varsigma} \mathcal{C}^{\beta}} = 0,$$

for any $1 - \gamma < \varsigma < \gamma - \rho$.

Proof. A combination of Proposition 26 and Proposition 15 gives the first claim. For the second, consider the field $X_{s,t}^n(x) = (T_{s,t}^{\omega}g_n - T_{s,t}^{\omega}g)(x)$. Then X satisfies condition (i)-(ii) of Lemma 14 with $H(x) = \sup_{|z| \le x} w^{-1}(z) ||T^{\omega}g_n - T^{\omega}g||_{\mathcal{C}_T^{\gamma}\mathcal{C}^{\kappa}(w)}$. Thus by Proposition 15 it follows that the integral operator $\Theta^{X^n} : \mathscr{C}^{\varsigma}\mathcal{C}^{\beta} \to \mathscr{C}^{\varsigma}\mathcal{C}^{\beta}$ built from X^n satisfies for any $y \in \mathscr{C}_T^{\varsigma}\mathcal{C}^{\beta}$

$$\left\|\Theta^{X^{n}}(y)\right\|_{\mathcal{C}^{5}_{T}\mathcal{C}^{\beta}}\lesssim_{S,\xi}\sup_{|z|\leq \|y\|_{\infty}}w^{-1}(z)\left\|T^{\omega}g_{n}-T^{\omega}g\right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)}\left(1+[y]_{\varsigma}\right)$$

Thus taking the limits when $n \to \infty$, it follows that $\|\Theta^{X^n}(y)\|_{\mathcal{C}^{\varsigma}_T\mathcal{C}^{\beta}} \to 0$ due to the assumption that $\|T^{\omega}g_n - T^{\omega}g\|_{\mathcal{C}^{\gamma}_T\mathcal{C}^{\kappa}(w)} \to 0$ when $n \to \infty$.

We are now ready to prove existence and uniqueness of solutions to (4.2) in the non-linear Volterra-Young formulation, through an application of the abstract existence and uniqueness results of equations on the form

$$\theta_t = p_t + \Theta(\theta)_t, \quad t \in [0, T],$$

proven in Theorem 16. In combination with Proposition 26 and 27 we consider $p_t = P_t \psi$ for some $\psi \in C^{\beta}$, Θ is the integral operator constructed in (4.4) with $S_t = P_t \xi$ for some $\xi \in C^{-\vartheta}$ with $\beta > \vartheta$.

Theorem 28. Consider a measurable path $\omega : [0, T] \to \mathbb{R}$ and $g \in S'(\mathbb{R})$, and suppose $T^{\omega}g \in C_T^{\gamma}C^{\kappa}(w)$ for some $\gamma > \frac{1}{2}$ and $\kappa \geq 3$.

Let $0 < \vartheta < 1$ and $\vartheta < \beta < 2 - \vartheta$. Let $\xi \in C^{-\vartheta}$ and $\rho = \frac{\beta + \vartheta}{2}$ and suppose that $1 - \gamma < \gamma - \rho$. Suppose $\psi \in C^{\beta}$. Then there exists a time $\tau \in (0, T]$ such that there exists a unique $\theta \in \mathcal{B}_{\tau}(P\psi) \subset \mathscr{C}_{\tau}^{\varsigma} C^{\beta}$ which solves the non-linear Young

equation

(4.5)
$$\theta_t = P_t \psi + \int_0^t P_{t-s} \xi T_{\mathrm{ds}}^\omega g(\theta_s), \quad t \in [0, \tau],$$

for any $\gamma - \rho > \varsigma > 1 - \gamma$ where $S_t := P_t \xi$ as defined Proposition 25, and the integral is understood in the sense of proposition 26. There exists a $C = C(P, \xi) > 0$ such that the solution θ satisfies the following bound

(4.6)
$$[\theta]_{\varsigma;\tau} \le C \Big([P\psi]_{\varsigma,\tau} + 2 \sup_{|z| \le 1 + \|\psi\|_{\mathcal{C}^{\beta}} + [P\psi]_{\varsigma;\tau}} w^{-1}(z) \|T^{\omega}g\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}} \Big)$$

Moreover, if $w \simeq 1$, then there exists a unique $\theta \in \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$ which is a global solution, in the sense that $\tau = T$.

Proof. By Proposition 25, it follows that $S_t := P_t \xi$ is a linear operator on C^{β} since $\beta > \vartheta$. It then follows from Proposition 26 that the non-linear Young integral

$$\Theta(\mathbf{y})_t := \int_0^t S_{t-s} T_{\mathrm{d}s}^\omega g(\mathbf{y}_s),$$

exists as a map $\Theta : [0, T] \times \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta} \to \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$, and for any $y \in \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$ with $\gamma - \rho > \varsigma > 1 - \gamma$ there exists a constant C > 0 such that

(4.7)
$$\left[\Theta(y)\right]_{\varsigma} \leq C \sup_{|z| \leq \|y\|_{\infty}} w^{-1}(z) \left\| T^{\omega}g \right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)} \left(1 + [y]_{\varsigma}\right) T^{\gamma - \rho - \varsigma}.$$

Moreover by Proposition 27 we have that for two paths $x, y \in \mathscr{C}^{\varsigma}_{T} \mathcal{C}^{\beta}$ there exists a constant $C = C(S, \xi) > 0$ such that

(4.8)
$$\left[\Theta(x) - \Theta(y) \right]_{\varsigma} \leq C \sup_{|z| \leq \|x\|_{\infty} \vee \|y\|_{\infty}} w^{-1}(z) \|T^{\omega}g\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)} (\|x_{0} - y_{0}\|_{\mathcal{C}^{\beta}} + [x - y]_{\varsigma}) T^{\gamma - \rho - \varsigma}.$$

Thus by theorem 16, for any $\tau > 0$ such that

$$\tau \leq \left[4(1+[P\psi]_{\varsigma})C\sup_{|z|\leq 1+\|\psi\|_{\mathcal{C}^{\beta}}+[P\psi]_{\varsigma}}w^{-1}(z)\right]^{-\frac{1}{\gamma-\rho-\varsigma}}$$

there exists a unique $\theta \in \mathscr{C}^{\varsigma}([0, \tau]; \mathcal{C}^{\beta}(\mathbb{R}))$ for $1 - \gamma < \varsigma < \gamma - \rho$ which satisfies (4.5). Thus local existence and uniqueness holds for (4.5).

We now move on to prove the bound in (4.6) First observe that

$$[\theta]_{\varsigma;\tau} \leq [P\psi]_{\varsigma,\tau} + \left[\Theta(\theta)\right]_{\varsigma;\tau}.$$

From Lemma 14, we know that there exists a $C = C(P, \xi) > 0$ such that

$$\left[\Theta(\theta)\right]_{\varsigma;\tau} \le C \sup_{|z| \le \|\theta\|_{\infty}} w^{-1}(z) \left\| T^{\omega} g \right\|_{\mathcal{C}^{\gamma}_{T} \mathcal{C}^{\kappa}} \left(1 + [\theta]_{\varsigma;\tau} \right) \tau^{\gamma - (\rho + \varsigma)}$$

Recall from Theorem 16 that $\tau > 0$ is chosen such that θ is contained in the unit ball centred at $S\psi$, i.e. $\theta \in \mathcal{B}_{\tau}(P\psi)$, and thus $[\theta]_{\varsigma;\tau} \leq [P\psi]_{\varsigma;\tau} + 1$, which yields

$$[\theta]_{\varsigma;\tau} \leq [P\psi]_{\varsigma,\tau} + 2 \sup_{|z| \leq 1 + \|\psi\|_{\mathcal{C}^{\beta}} + [P\psi]_{\varsigma;\tau}} w^{-1}(z) \|T^{\omega}g\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}}.$$

At last, if $w \simeq 1$, then for any $x, y \in \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$, there exists an M > 0 such that

$$C \sup_{|z| \le \|x\|_{\infty} \lor \|y\|_{\infty}} w^{-1}(z) \le M.$$

where *C* is the largest of the constants in (4.7) and (4.8). By Theorem 17 it follows that a global solution exists, i.e. there exists a unique $\Theta \in \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$ satisfying (4.5).

With the above theorem at hand, we are ready to prove Theorem 2 and Corollary 4.

Proof of Theorem 2. From Theorem 28, we know that there exists a $\tau > 0$ such that there exists a unique $\theta \in \mathscr{C}^{\varsigma}_{\tau} C^{\beta}$ which solves (4.5). Thus we say that there exists a unique local solution u in the sense of Definition 1. As proven in Theorem 28, there exists a global solution if $T^{\omega}g \in C^{\gamma}_{T}C^{\kappa}(w)$ with $w \simeq 1$. This concludes the proof of theorem 2.

Proof of Corollary 4. Let $\{g_n\}$ be a sequence of smooth functions converging to g such that $T^{\omega}g_n$ converges to $T^{\omega}g$ in $\mathcal{C}_T^{\gamma}\mathcal{C}^{\kappa}$, and define the field $Y_{s,t}^n(x) = (t-s)g_n(x+\omega_s)$. It is readily checked that Y^n satisfies (i)-(ii) of Lemma 14, and thus Θ^{Y^n} is an integral operator from $\mathcal{C}_T^{\varsigma}\mathcal{C}^{\beta} \to \mathcal{C}_T^{\varsigma}\mathcal{C}^{\beta}$. Furthermore, it follows from Theorem 16 together with the same reasoning as in the proof of Theorem 28 that there exists a unique solution θ^n to the equation

(4.9)
$$\theta_t^n = P_t \psi + \Theta^{Y^n}(\theta)_t.$$

In particular, $\|\theta^n\|_{\mathscr{C}^{\varsigma}_{\tau}C^{\beta}} < \infty$ for all $n \in \mathbb{N}$. Note that then Θ^{Y^n} corresponds to the classical Riemann integral on \mathcal{C}^{β} . Indeed, since g_n is smooth, the integral Θ^{Y^n} is given by

$$\Theta^{Y^n}(\mathbf{y})_t := \int_0^t S_{t-s}g(\mathbf{y}_s)\,\mathrm{d}s,$$

where we recall that $S_t = P_t \xi$. Furthermore, it is readily checked that Θ^{Y^n} agrees with the integral operator $\Theta^{T^{\omega}g_n}$, since g_n is smooth, $T_{s,t}^{\omega}g_n(x) = \int_s^t g(x + \omega_r) dr$ is differentiable in time. Thus note that the difference between the two approximating integrals

$$\Theta_{\mathcal{P}}^{Y^n}(y)_t = \sum_{[u,v]\in\mathcal{P}[0,t]} S_{t-u}g_n(y_u + \omega_u)(v-u)$$
$$\Theta_{\mathcal{P}}^{T^{\omega}g_n}(y)_t = \sum_{[u,v]\in\mathcal{P}[0,t]} S_{t-u}\int_u^v g_n(y_u + \omega_r) \, \mathrm{d}r$$

can be estimated by

$$\left\|\sum_{[u,v]\in\mathcal{P}}S_{t-u}\int_{u}^{v}g(y_{u}+\omega_{u})-g(y_{u}+\omega_{r})\,\mathrm{d}r\,\right\|_{\mathcal{C}^{\beta}}\to0$$

when $|\mathcal{P}| \to 0$. Therefore $\Theta^{Y^n} \equiv \Theta^{T^{\omega}g_n}$. Let $\theta \in \mathscr{C}^{\varsigma}_{\tau} \mathcal{C}^{\beta}$ be the solution to (4.5) constructed from $T^{\omega}g$ with integral operator Θ . We now investigate the difference between θ and θ^n as found in (4.9). It is readily checked that

$$\left[\theta^{n}-\theta\right]_{\varsigma} \leq \left[\Theta(\theta)-\Theta^{T^{\omega}g_{n}}(\theta_{n})\right]_{\varsigma}+\left[\Theta^{T^{\omega}g_{n}}(\theta)-\Theta^{Y^{n}}\right]_{\varsigma}.$$

As already argued, the last term on the right hand side is equal to zero, and we are therefore left to show that the first term on the right hand side converge to zero. Invoking Proposition 15 (in particular (2.17)) and following along the lines of the proof of Proposition 27, it is readily checked that there exists a $C = C(P, \xi) > 0$ such that

$$\left[\theta^{n}-\theta\right]_{\varsigma;\tau} \leq C \sup_{|z|\leq \|\theta\|_{\infty}\vee\|\theta^{n}\|_{\infty}} w^{-1}(z) \left(\left\|T^{\omega}g_{n}-T^{\omega}g\right\|_{\mathcal{C}_{T}^{\gamma}\mathcal{C}^{\kappa}(w)}\left(1+[\theta]_{\varsigma;\tau}\vee\left[\theta^{n}\right]_{\varsigma;\tau}\right)+\left[\theta-\theta^{n}\right]_{\varsigma;\tau}\right)\tau^{\gamma-\rho-\varsigma}$$

We can now choose a parameter $\bar{\tau} > 0$ sufficiently small, such that

$$\left[\theta^{n}-\theta\right]_{\varsigma;\bar{\tau}} \leq 2 \sup_{|z|\leq \|\theta\|_{\infty}\vee\|\theta^{n}\|_{\infty}} w^{-1}(z) \left(\left\|T^{\omega}g_{n}-T^{\omega}g\right\|_{\mathcal{C}^{\gamma}_{T}\mathcal{C}^{\kappa}(w)}\left(1+[\theta]_{\varsigma;\tau}\vee\left[\theta^{n}\right]_{\varsigma;\tau}\right)\right)$$

From (4.6) we know that there exists a constant M > 0 such that

$$\left\|\theta^{n}\right\|_{\infty} \vee \|\theta\|_{\infty} \vee \left[\theta^{n}\right]_{\varsigma;\tau} \vee [\theta]_{\varsigma;\tau} \leq M,$$

and it follows that

(4.10)
$$\left[\theta^n - \theta\right]_{\varsigma;\bar{\tau}} \leq 2 \sup_{|z| \leq M} w^{-1}(z) \left\| T^{\omega} g_n - T^{\omega} g \right\|_{\mathcal{C}^{\gamma}_{T} \mathcal{C}^{\kappa}(w)} (1+M).$$

The above inequality can, by similar arguments, be proven to hold for any sub-interval $[a, a + \overline{\tau}] \subset [0, \tau]$ and thus we conclude by [13, Exc. 4.24] that there exists a constant $C = C(P, \xi, M, \overline{\tau}, w) > 0$ such that

$$\left[\theta^n - \theta\right]_{\varsigma;\tau} \le C \left\| T^{\omega} g_n - T^{\omega} g \right\|_{\mathcal{C}_T^{\gamma} \mathcal{C}^{\kappa}(w)}$$

Taking the limit when $n \to \infty$ it follows that $\theta^n \to \theta$ due to the assumption that $T^{\omega}g_n \to T^{\omega}g$. If the solution is global, i.e. $w \simeq 1$, then the same arguments as above holds, and the inequality in (4.10) can be proven to hold on any sub-interval $[a, a + \overline{\tau}] \subset [0, T]$ and then by the same arguments as above $\theta^n \to \theta$ in $\mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$.

4.2. The drifted multiplicative SHE

So far we have proven existence and uniqueness of solutions to equations on the form

$$u_t = P_t \psi + \int_0^t P_{t-s} \xi g(u_s) \, \mathrm{d}s + \omega_t, \quad s \in [0, T], u_0 = \psi \in \mathcal{C}^\beta(\mathbb{R}).$$

That is, we have throughout the text assumed that $\xi \in C^{-\vartheta}$ with $\vartheta < \beta$ is a multiplicative spatial noise. More generally, it is frequently considered a drift term in the above equation as well, such that

$$u_{t} = S_{t}\psi + \int_{0}^{t} P_{t-s}b(u_{s}) \,\mathrm{d}s + \int_{0}^{t} P_{t-s}\xi g(u_{s}) \,\mathrm{d}s + \omega_{t}, \quad s \in [0, T], u_{0} = \psi \in \mathcal{C}^{\beta}(\mathbb{R}).$$

Our results extends easily to this type of drifted equation as well, as long as $T^{\omega}b \in C_T^{\gamma}C^3(w)$. Indeed, again setting $\theta = u - \omega$, we consider the non-linear Young–Volterra equation

(4.11)
$$\theta_t = P_t \psi + \int_0^t P_{t-s} T^{\omega}_{ds} b(\theta_s) + \int_0^t P_{t-s} \xi T^{\omega}_{ds} g(\theta_s).$$

It is straightforward to construct the non-linear Young–Volterra integral of the form $\int_0^t P_{t-s} T_{ds}^{\omega} b(\theta_s)$. Since there is no multiplicative noise ξ , the linear operator P_t on C^{β} is not singular in its time argument. An application of Lemma 14 would then allow to construct $y \mapsto \Theta^b(y) := \int_0^t P_{t-s} T_{ds}^{\omega} b(y_s)$ as an operator from $\mathscr{C}_T^{\varsigma} C^{\beta}$ into itself. The second integral is constructed as before as an operator from $\mathscr{C}_T^{\varsigma} C^{\beta}$ into itself; let us denote this by Θ^g . Then (4.11) can be written as the abstract equation

$$\theta_t = P_t \psi + \Theta(\theta)_t,$$

where $\Theta(y) := \Theta^b(y) + \Theta^g(y)$ for $y \in \mathscr{C}_T^{\varsigma} \mathcal{C}^{\beta}$. An application of Theorem 16 then provides local existence and uniqueness. If both $T^{\omega}b$ and $T^{\omega}g$ are contained in some unweighted Besov spaces, global existence and uniqueness also holds, as proven in Theorem 17.

While there is no spatial noise in the drift, we do not get any extra regularizing effect from the heat kernel P_t on the drift coefficient b, as what is also observed for the simple non-linear heat equation of the form

$$\partial_t u = \Delta u + b(u),$$

where Lipschitz is still required for b. Considering Schauder type stability estimates for the difference

$$\int_0^t P_{t-s}b(u_s)\,\mathrm{d}s - \int_0^t P_{t-s}b(u_s')\,\mathrm{d}s$$

where $u, u' \in C^{\beta}$, we see indeed that we need to assume Lipschitz regularity on b in order to get a Lipschitz estimate for the above difference. Formulating similar estimates in terms of non-linear Young theory will behave in a similar way. As a consequence, when including time homogeneous drifts b as in (4.11) we expect the same regularity assumptions on b as we have proven for g.

5. Averaged fields with Lévy noise

In this section we will construct the averaged field associated to a fractional Levy process, and prove it's regularizing properties. In contrast to the works of [24], we will not consider the regularity of the local time associated with a process in order to obtain the regularizing effect, but estimate the regularity of the averaged field $T^{\omega}g$ directly based on probabilistic techniques developed in [8]. This has the benefit that it improves the regularity of the averaged field, as discussed in the beginning of Section 3.

5.1. Linear fractional stable motion

In this section we introduce a class of measurable processes which have the regularization property, and allows us to deal with our equations.

Definition 29. Let $\alpha \in (0, 2]$. We say that L is a d-dimensional symmetric α -stable Lévy process if

(i) $L_0 = 0$,

- (ii) L has independent and stationary increments,
- (iii) L is cádlàg and
- (iv) there exists a constant such that for all $\lambda \in \mathbb{R}^d$

$$\mathbb{E}\left[\exp(iL_t\cdot\xi)\right] = e^{-c_{\alpha}|\xi|^{\alpha}t}.$$

Following [41], we define a fractional process with respect to a stable Lévy process as followed:

Definition 30. Let $\alpha \in (0, 2]$ and L, \tilde{L} be two independent *d*-dimensional α -stable symmetric Lévy processes. Let $H \in (0, 1) \setminus \{\alpha^{-1}\}$. For all $t \in \mathbb{R}_+$, we define the α -linear fractional stable motion (α -LFSM) of Hurst parameter H by

$$L_t^H = \int_0^t (t - v)^{H - \frac{1}{\alpha}} \, \mathrm{d}L_t + \int_0^{+\infty} (t + v)^{H - \frac{1}{\alpha}} - v^{H - \frac{1}{\alpha}} \, \mathrm{d}\tilde{L}_v$$

We extend this definition to $H = \frac{1}{\alpha}$ by setting $L^{\frac{1}{\alpha}} = L$.

The existence of the LFSM is proved in [41], Chapter 3. Note that when $\alpha = 2$, we recover the standard fractional Brownian motion (up to a multiplicative constant), and its different representations. Here we have chosen the moving average representation, in the fractional Brownian motion case, one could also use the harmonizable representation or something else. Note that when $\alpha < 2$, those different representations are no longer equivalent.

We gather here some properties about the LFSM which will become useful in the subsequent analysis of the regularity of averaged fields associated to LFSM. One can again consult [41], but also [31] for the first three points. The last point is a direct consequence of the definition of L^{H} .

Proposition 31. Let $\alpha \in (0, 2]$ and let $H \in (0, 1) \cup \{\alpha^{-1}\}$, and let L^H be an α -LFSM.

- (i) L^H is almost surely measurable.
- (ii) L^H is continuous if and only if $\alpha = 2$ or $H > \frac{1}{\alpha}$.
- (iii) L^H is H self-similar and its increments are stationary.
- (iv) For all $t \ge 0$ define $\mathcal{F}_t = \sigma(\{\tilde{L}_{(\tilde{t})}\}_{\tilde{t}\in\mathbb{R}_+}) \cup \{L_s: s \le t\})$. Then L^H is $(\mathcal{F}_t)_{t\ge 0}$ adapted.
- (v) For all $0 \le s \le r$, we have

$$L_r^H = L_{s,r}^{1,H} + L_{s,r}^{2,H},$$

with

$$L_{s,r}^{1,H} = \int_s^r (r-v)^{H-\frac{1}{\alpha}} \,\mathrm{d}L_v$$

and

$$L_{s,r}^{2,H} = \int_0^s (r-v)^{H-\frac{1}{\alpha}} \, \mathrm{d}L_v + \int_0^{+\infty} (r+v)^{H-\frac{1}{\alpha}} - v^{H-\frac{1}{\alpha}} \, \mathrm{d}\tilde{L}_v.$$

Hence, $L_{s,r}^{1,H}$ is independent of \mathcal{F}_r while $L_{s,r}^{2,H}$ is measurable with respect to \mathcal{F}_r , and for all $x \in \mathbb{R}^d$

$$\mathbb{E}\left[e^{i\xi\cdot L_{s,r}^{1,H}}\right] = e^{-c_{\alpha}|\xi|^{\alpha}(s-r)^{\alpha H}}$$

Note that we can (and we will) reformulate the last point by saying that for all $g \in S$,

$$\mathbb{E}\left[g\left(x+L_{s,r}^{1,H}\right)\right] = P_{c_{\alpha}(r-s)^{\alpha H}}^{\frac{\alpha}{2}}g(x)$$

Finally, let us recall a useful result about martingales in Lebesgue spaces due to Pinelis [39] and adapted in weighted spaces (see Appendix A.1).

Proposition 32. Let $p \ge 2$, let w be an admissible weight, and let $(M_n)_n$ be a $(\Omega, \mathcal{F}, (\mathcal{F}_n)_n, \mathbb{P})$ martingale with value in $L^p(\mathbb{R}^d, w; \mathbb{R})$. Suppose that there exists a constant c > 0 such that for all $n \ge 0$, $||M_{n+1} - M_n||_{L^p(w)} \le c$. Then there exists two constants $C_1 > 0$ and $C_2 > 0$ such that for all $x \ge 0$,

$$\mathbb{P}\left(\left\|M_N - \mathbb{E}[M_N]\right\|_{L^p(w)} \ge x\right) \le C_1 e^{-\frac{C_2 x^2}{Nc^2}}.$$

5.2. Averaging operator

In this section we give another proof of the regularizing effect of the fractional Brownian motion. This proof is similar to the one in [8], but provides bounds in general weighted Besov spaces directly. For other proofs of the same kind of results, one can consult [9,17] and [30]. Other proofs may rely on Burkolder-Davis-Gundy inquality in UMD spaces. For information about UMD spaces, one can consult [28].

Lemma 33. Let w be an admissible weight. Let $\alpha \in (0, 2]$ and let $H \in (0, 1) \cup \{\alpha^{-1}\}$. Let $v \in (0, 1)$. Then there exists three constants $K, C_1, C_2 > 0$ such that for all $j \ge -1$, all $g \in S'(\mathbb{R}^d; \mathbb{R})$ all $s \le t$ and all x > K

$$\mathbb{P}\left(\frac{\|\int_{s}^{t} \Delta_{j}g(\cdot+L_{r}^{H})\,\mathrm{d}r\|_{L^{p}(w)}}{|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|\Delta_{j}g\|_{L^{p}(w)}} \ge x\right) \le C_{1}e^{-C_{2}x^{2}}.$$

Proof. Let $0 \le s \le t$ and let $N \in \mathbb{N}$ to be fixed from now. Let us define $t_n = \frac{n}{N}(t-s) + s$, and $\mathcal{G}_n = \mathcal{F}_{t_n}$. Let us define

$$M_n(x) = \mathbb{E}\left[\int_s^t \Delta_j g\left(x + L_r^H\right) \mathrm{d}r |\mathcal{F}_{t_n}\right].$$

Hence, $(M_n)_n$ is a martingale with respect to $(\mathcal{G}_n)_n$ and with value in $L^p(\mathbb{R}^d, w; \mathbb{R})$ for all $p \ge 2$, and furthermore. Note also that

$$M_N = \int_s^t \Delta_j g(\cdot + L_r^H) \,\mathrm{d}r$$

Furthermore, we have for all $0 \le n \le N - 1$

$$M_{n+1} - M_n = \int_{t_n}^{t_{n+1}} \Delta_j g(\cdot + L_r^H) \, \mathrm{d}r + A_{n+1} - A_n$$

with

$$A_n = \mathbb{E}\left[\int_{t_n}^t \Delta_j g\left(\cdot + L_r^H\right) \mathrm{d}r |\mathcal{F}_{t_n}\right].$$

It is also readily checked that for all n = 0, ..., N we have

$$\left\|\int_{t_n}^{t_{n+1}} \Delta_j g\left(\cdot + L_r^H\right) \mathrm{d} r\right\|_{L^p(w)} \leq \frac{|t-s|}{N} \|\Delta_j g\|_{L^p(w)}$$

Furthermore, using the decomposition of Proposition 31, one get for all $n \in \{0, ..., N\}$,

$$A_{n}(x) = \int_{t_{n}}^{t} \mathbb{E} \Big[\Delta_{j} g \big(x + L_{t_{n},r}^{1,H} + L_{t_{n},r}^{2,H} \big) |\mathcal{F}_{t_{n}} \Big] dr$$
$$= \int_{t_{n}}^{t} P_{(r-t_{n})^{\alpha H}}^{\frac{\alpha}{2}} \Delta_{j} g \big(x + L_{t_{n},r}^{2,H} \big) dr,$$

Hence, thanks to Proposition 48, the following bound holds:

$$\|A_n\|_{L^p(w)} \lesssim \int_{t_n}^t e^{-c2^{\alpha j}(r-t_n)^{\alpha H}} \|\Delta_j g\|_{L^p(w)} \lesssim 2^{-\frac{j}{H}} \|\Delta_j g\|_{L^p(w)} \int_0^{+\infty} e^{-cr^{\alpha H}} \, \mathrm{d}r.$$

Finally, we have the bound

(5.1)
$$\|M_{n+1} - M_n\|_{L^p(w)} \lesssim \left(\frac{t-s}{N} + 2^{-\frac{j}{H}}\right) \|\Delta_j g\|_{L^p(w)}$$

Note also that we have the straightforward bound $||M_{n+1} - M_n||_{L^p(w)} \leq (t-s) ||\Delta_j g||_{L^p(w)}$. Therefore, by interpolation, this gives

(5.2)
$$\|M_{n+1} - M_n\|_{L^p(w)} \lesssim (t-s)^{1-\nu} \left(\frac{t-s}{N} + 2^{-\frac{j}{H}}\right)^{\nu} \|\Delta_j g\|_{L^p(w)}.$$

It is readily checked that

$$\mathbb{P}\left(\|M_{N}\|_{L^{p}(w)} \geq x|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|\Delta_{j}g\|_{L^{p}(w)}\right)$$

$$\leq \mathbb{P}\left(\|M_{N}-\mathbb{E}[M_{N}]\|_{L^{p}(w)} \geq \frac{x}{2}|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|\Delta_{j}g\|_{L^{p}(w)}\right)$$

$$+\mathbb{P}\left(\|\mathbb{E}[M_{N}]\|_{L^{p}(w)} \geq \frac{x}{2}|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|g\|_{L^{p}(w)}\right).$$

Furthermore, note that $\mathbb{E}[M_N] = \mathbb{E}[A_0]$, and $||A_0||_{L^p(w)} \le (t-s)||\Delta_j g||_{L^p(w)}$, where $||A_0||_{L^p(w)} \le 2^{-\frac{j}{H}} ||\Delta_j g||_{L^p(w)}$. Hence, there exists a constant K > 0 such that

$$\left\|\mathbb{E}[M_N]\right\|_{L^p(w)} \leq \frac{K}{2} |t-s|^{1-\frac{\nu}{2}} 2^{-j\frac{\nu}{2H}} \|\Delta_j g\|_{L^p(w)}.$$

Therefore, for x > K,

$$\mathbb{P}\bigg(\big\|\mathbb{E}[M_N]\big\|_{L^p(w)} \ge \frac{x}{2}|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|\Delta_j g\|_{L^p(w)}\bigg) = 0.$$

Finally, this gives us, thanks to Proposition 32, there exists a constant \tilde{C}_2 such that

$$\mathbb{P}\left(\left\|\int_{s}^{t} \Delta_{j}g\left(\cdot+L_{r}^{H}\right) \mathrm{d}r\right\|_{L^{p}(w)} \geq x|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|\Delta_{j}g\|_{L^{p}(w)}\right)$$
$$\lesssim \exp\left(-\tilde{C}_{2}\frac{x^{2}|t-s|^{2-\nu}2^{-j\frac{\nu}{H}}\|\Delta_{j}g\|_{L^{p}(w)}^{2}}{N((t-s)^{1-\nu}(\frac{t-s}{N}+2^{-j\frac{j}{H}})^{\nu})^{2}\|\Delta_{j}g\|_{L^{p}(w)}^{2}}\right).$$

When optimizing in N, one obtains the desired result.

Remark 34. Thanks to the bound (5.2) and to UMD spaces standard arguments (see for example [28], one could rather use the Burkoldher-Davis-Gundy inequality, and get for all $1 , and all <math>m \ge 2$

$$\mathbb{E}\left[\left\|\int_{s}^{t}\Delta_{j}g\left(\cdot+L_{r}^{H}\right)\mathrm{d}r\right\|_{L^{p}(w)}^{m}\right]\lesssim_{m}|t-s|^{\frac{m}{2}}2^{-\frac{m}{2H}j}\|\Delta_{j}g\|_{L^{p}(w)}^{m}.$$

This would allow for a more standard proof of regularizing effect of LFSM (using Kolmogorov continuity theorem), but prevents to have Gaussian tails.

Lemma 35. Let X be a non-negative random variable such that there exists $K, C_1, C_2 > 0$ such that for all x > K,

$$\mathbb{P}(X \ge x) \le C_1 e^{-C_2 x^2}$$

Then for all $m \ge 1$

$$\mathbb{E}\left[X^{2m}\right] \le K^{2m} + \frac{C_1}{2} \frac{m!}{C_2^m}.$$

Proof. A simple computation reveals that

$$\mathbb{E}[X^{2m}] = \int_0^K 2m X^{2m-1} \mathbb{P}(X \ge x) \, \mathrm{d}x + \int_K^{+\infty} 2m X^{2m-1} \mathbb{P}(X \ge x) \, \mathrm{d}x$$
$$\leq K^{2m} + C_1 \int_0^{+\infty} e^{-C_2 x^2} \, \mathrm{d}x$$

$$= K^{2m} + C_1 C_2^{-m} \int_0^{+\infty} e^{-x^2} dx$$

= $K^{2m} + \frac{C_1}{2} \frac{m!}{C_2^m}$.

Finally, we have all the tools to prove the following theorem of regularity of the averaging operator with respect the LFSM. We will use a version of the Garsia-Rodemich-Rumsey inequality from Friz and Victoir ([12] -Theorem 1.1 p. 571).

Theorem 36. Let w be an admissible weight. Let $\alpha \in (0, 2]$. Let $H \in (0, 1) \cup \{\alpha^{-1}\}$. Let $\kappa \in \mathbb{R}$ and $\nu \in [0, 1]$. Let $2 \le p \le \infty$ and let $1 \le q \le +\infty$. There exists a constant C > 0 such that for any $g \in B_{p,q}^{\kappa}(w)$ a positive random variable F exists with

$$\mathbb{E}\big[e^{CF^2}\big] < +\infty,$$

and such that for any $0 \le s \le t \le T$

$$\|T_{s,t}^{L^{H}}g\|_{B_{p,q}^{\kappa+\frac{\nu}{2H}}(w)} \le F|t-s|^{1-\frac{\nu}{2}-\frac{\varepsilon}{2}}\|g\|_{B_{p,q}^{\kappa}(w)} \quad if q < +\infty$$

and

$$\|T_{s,t}^{L^{H}}g\|_{B_{p,\infty}^{\kappa+\frac{\nu}{2H}-\delta}(w)} \leq F|t-s|^{1-\frac{\nu}{2}-\frac{\varepsilon}{2}}\|g\|_{B_{p,\infty}^{\kappa}(w)}.$$

Proof. First consider $1 \le q < \infty$. Thanks to Lemma 33 and Lemma 35, we know that there exists a constant $c_q > 0$ such that for all $0 \le s < t \le T$

$$\mathbb{E}\left[\left(\frac{\|\Delta_{j}T_{s,t}^{L^{H}}g\|_{L^{p}(w)}}{|t-s|^{1-\frac{\nu}{2}}2^{-j\frac{\nu}{2H}}\|\Delta_{j}g\|_{L^{p}(w)}}\right)^{q}\right] \leq c_{q}$$

and that for all $m \ge 1$,

$$\mathbb{E}\left[\left\|\Delta_{j}T_{s,t}^{L^{H}}g\right\|_{L^{p}(w)}^{2m}\right] \lesssim \left(K^{2m} + \frac{m!}{C_{2}^{m}}\right)|t-s|^{(2-\nu)m}2^{-j\frac{\nu m}{H}}\|\Delta_{j}g\|_{L^{p}(w)}^{2m}.$$

Let us take $C < C_2$. We have

$$\mathbb{E}\left[\exp\left(C\left(\frac{\|T_{s,t}^{L^{H}}g\|_{B_{p,q}^{\kappa+\frac{\nu}{2H}}(w)}}{|t-s|^{1-\frac{\nu}{2}}\|g\|_{B_{p,q}^{\kappa}(w)}}\right)^{2}\right)\right] = A + B,$$

with

$$A = \sum_{2m \le q} \frac{C^m}{m! |t - s|^{(2-\nu)m} \|g\|_{B^{\kappa}_{p,q}(w)}^{2m}} \mathbb{E}\bigg[\bigg(\sum_{j \ge 1} 2^{j(\kappa + \frac{\nu}{2H})q} \|\Delta_j T_{s,t}^{L^H} g\|_{L^p(w)}^q \bigg)^{\frac{2m}{q}} \bigg]$$

and

$$B = \sum_{2m>q} \frac{C^m}{m!} \mathbb{E}\bigg[\bigg(\sum_{j\geq -1} \frac{2^{jq\kappa} \|\Delta_j g\|_{L^p(w)}^q}{\sum_{j\geq -1} 2^{jq\kappa} \|\Delta_j g\|_{L^p(w)}^q} \bigg(\frac{\|\Delta_j T_{s,t}^{LH} g\|_{L^p(w)}}{|t-s|^{1-\frac{\nu}{2}} 2^{-j\frac{\nu}{2}} \|\Delta_j g\|_{L^p(w)}} \bigg)^q \bigg)^{\frac{2m}{q}} \bigg].$$

For A, we use Jensen inequality in the concave case, and we have thanks to Lemma 35

$$A \leq \sum_{2m \leq q} \frac{C^{m}}{m! |t - s|^{(2-\nu)m} \|g\|_{B^{\kappa}_{p,q}(w)}^{2m}} \left(\sum_{j \geq 1} 2^{j(\kappa + \frac{\nu}{2H})q} \mathbb{E} \left[\left\| \Delta_{j} T^{L^{H}}_{s,t} g \right\|_{L^{p}(w)}^{q} \right\| \right] \right)^{\frac{2m}{q}}$$
$$\lesssim \sum_{2m < q} \frac{C^{m} (c_{q}^{\frac{2}{q}})^{m}}{m!}$$

.

$$\lesssim e^{Cc_q^{\frac{2}{q}}}.$$

For B, we use Jensen inequality in the convex case, and we have again thanks to Lemma 35

$$\begin{split} B &= \sum_{2m>q} \frac{C^m}{m!} \mathbb{E} \bigg[\bigg(\sum_{j\geq -1} \frac{2^{jq\kappa} \|\Delta_j g\|_{L^p(w)}^q}{\sum_{j\geq -1} 2^{jq\kappa} \|\Delta_j g\|_{L^p(w)}^q} \bigg(\frac{\|\Delta_j T_{s,t}^{L^H} g\|_{L^p(w)}}{|t-s|^{1-\frac{\nu}{2}} 2^{-j\frac{\nu}{2}} \|\Delta_j g\|_{L^p(w)}} \bigg)^q \bigg)^{\frac{2m}{q}} \bigg] \\ &\leq \sum_{2m>q} \frac{C^m}{m!} \sum_{j\geq -1} \frac{2^{jq\kappa} \|\Delta_j g\|_{L^p(w)}^q}{\sum_{j\geq -1} 2^{jq\kappa} \|\Delta_j g\|_{L^p(w)}^q} \mathbb{E} \bigg[\bigg(\frac{\|\Delta_j T_{s,t}^{L^H} g\|_{L^p(w)}}{|t-s|^{1-\frac{\nu}{2}} 2^{-j\frac{\nu}{2}} \|\Delta_j g\|_{L^p(w)}} \bigg)^{2m} \bigg] \\ &\lesssim \sum_{2m>q} \frac{C^m}{m!} \bigg(K^{2m} + \frac{m!}{C_2^m} \bigg) \\ &\lesssim e^{CK^2} + \frac{1}{1 - \frac{C}{C_2}}. \end{split}$$

Hence, if we define

$$F = \int_0^T \int_0^T \exp\left(C\left(\frac{\|T_{s,t}^{L^n}g\|_{B_{p,q}^{\kappa+\frac{\nu}{2H}}(w)}}{|t-s|^{1-\frac{\nu}{2}}\|g\|_{B_{p,q}^{\kappa}(w)}}\right)^2\right) \mathrm{d}s \,\mathrm{d}t,$$

τī

we are exactly in the scope of the Garsia-Rodemich-Rumsey inequality with $\psi(x) = e^{Cx^2}$ and $p(u) = u^{1-\frac{\nu}{2}}$, which leads to the wanted result when $q < +\infty$. When $q = \infty$ one needs to deal with the supremum in q, and thus needs to lose a bit in q. We leave it to the reader since its a direct adaptation of the previous proof.

With the above construction of the averaged field associated with the fractional Lévy process, we are ready to prove Theorem 5.

Proof of Theorem 5. As in Section 4, let us first consider a distribution $\xi \in C^{-\vartheta}$ and an initial condition $\psi \in C^{\beta}$ with $0 < \vartheta < 1$ and $\beta = \vartheta + \varepsilon_1$. The singularity $\rho = \frac{\vartheta + \beta}{2}$ is one of the limiting thresholds in all the previous computations. Hence, the best choice of $2 - \vartheta > \beta > \vartheta$ for the space of the initial condition is for a small $\varepsilon_1 > 0$. Note also that in Theorem 36 one wants to take ν as close as possible to 1, since $\frac{\vartheta}{2H}$ is the index of regularization of the LFSM. Note also that for some $\varepsilon_2 > 0$ small enough, one has

$$\gamma = 1 - \frac{\nu}{2} - \varepsilon_2.$$

The condition which allows the nonlinear Young-Volterra calculus to work is

$$\frac{1}{2} < \gamma - \rho$$

here it gives us

$$\vartheta < 1 - \nu - \frac{\varepsilon_1}{2} - \varepsilon_2$$

which gives for some $\varepsilon_3 > 0$

$$\nu = 1 - \vartheta - \frac{\varepsilon_1}{2} - \varepsilon_2 - \varepsilon_3.$$

Hence, for an admissible weight w and a function $g \in C^{\kappa}$ with $\kappa > 3 - \frac{1-\vartheta}{2H}$, there exists $\varepsilon_1 > 0$, $\varepsilon_2 > 0$ and $\varepsilon_3 > 0$ such that all the previous condition are satisfied and such that almost surely

$$T^{L^{H}} \in \mathcal{C}^{\gamma}_{T}\mathcal{C}^{3}(w).$$

Applying Theorem 2 and 4 we have the desired result.

Proof of Corollary 6. The proof is straightforward when one recall that the space white noise ξ is almost surely in $C^{-\vartheta}$ for any $\vartheta > \frac{d}{2}$. One can then apply Theorem 5.

6. Conclusion

We have proven that a certain class of measurable paths $\omega : [0, T] \to \mathbb{R}$ provides strong regularizing effects on the multiplicative stochastic heat equation of the form of (1.8). In particular, we prove that there exists measurable paths ω such that local existence and uniqueness of such equations holds even when the non-linear function g is only a Schwartz distribution and $\xi \in C^{-\vartheta}$ for some $\vartheta < 1$, thus allowing for very rough spatial noise in the one dimensional setting. To this end, we apply the concept of non-linear Young integration, and extends this to the infinite dimensional setting with Volterra operators. This sheds new light on the application of the "pathwise regularization by noise" techniques developed in [8] to the context of SPDEs, and we believe that this program can be taken further in several directions in the future, and we provide some thoughts on such developments here.

In the current article we restricted our analysis to the spatial space white noise on \mathbb{T} (Corollary 6). Our techniques could be extended to \mathbb{T}^d , but with the same techniques, one could not allow for white spatial noise ξ . Indeed, recall that if $\xi : \mathbb{T}^d \to \mathbb{R}$ then $\xi \in C^{-\frac{d}{2}-\varepsilon}$ for any $\varepsilon > 0$ and thus even in d = 2 the noise is too rough, since the best we can deal with in the Young setting developed above is when $\xi \in C^{-\vartheta}$ with $\vartheta < 1$. However, there is a possibility that techniques from the theory of paracontrolled calculus as developed in [19] could be applied here to make sense of the product, and thus a generalization could then be possible. In that connection, one would possibly need "second order correction terms" associated to the averaged field $T^{\omega}g$ which at this point is unclear (at least to us) how one should construct.

Another possible direction would be to allow multiplicative space-time noise, i.e. consider ξ as a distribution on $[0, T] \times \mathbb{R}^d$ (or \mathbb{T}^d). When ξ is depending on time, one can no longer use the non-linear Young integral in the same way as developed here. The situation looks like the one encountered when considering regularization by noise for ODEs with multiplicative (time dependent) noise of the form

(6.1)
$$y_t = y_0 + \int_0^t b(y_s + \omega_s) \,\mathrm{d}\beta_s + \omega_t, \quad y_0 \in \mathbb{R}^d.$$

Existence and uniqueness of the above equation when β is a fractional Brownian motion with $H > \frac{1}{2}$ was recently established in [18], even when b is a distribution. The key idea in this result was to consider the average operator

$$\Gamma^{\omega}_{s,t}b(x) := \int_{s}^{t} b(x+\omega_{r}) \,\mathrm{d}\beta_{r}$$

and use a recently developed probabilistic lemma by Hairer and Li [21] to show that the regularity of $\Gamma^{\omega}b$ is linked to the regularity of $T^{\omega}b$. By considering an averaging operator given on the form on a Banach space

(6.2)
$$\Pi^{\omega}_{s,t}b(x) = \int_{s}^{t} b(x+\omega_{s}) \,\mathrm{d}\xi_{s}$$

where $x \in E$ and $\xi_t \in E$ is a time-colored spatially-white noise, for some Banach space *E*. If one can extend the lemmas of Hairer and Li to the infinite dimensional setting, showing the connection between the regularity of $\Pi^{\omega}b$ and $T^{\omega}b$ (when considering $T^{\omega}b$ as an infinite dimensional averaged field, as in Proposition 23) there is a possibility that one could prove regularization by noise for stochastic heat equations with space-time noise on the form

$$x_t = P_t \psi + \int_0^t P_{t-s} g(x_s) \,\mathrm{d}\xi_s + \omega_t,$$

by using similar techniques as developed in the current article. We leave a deeper investigation into these possibilities open for future work.

Appendix A: Basic concepts of Besov spaces and properties of the heat kernel

We gather here some material about Besov spaces, heat kernel estimates and embedding in those spaces. This section is strongly inspired by [3] and [42]. See also [33] for the weighted besov norms. For the sake of the comprehensions, we give elementary proofs of the main points for developing the theory. This is the purpose of Section A.1 and A.2. For the sake of the Volterra sewing lemma (see Section 2.1), we need a few non-standard estimates on the heat kernel action on Besov spaces. This is the purpose of Section A.3. In Section B we prove the elementary Cauchy-Lipschitz theorem for multiplicative SHE without additive perturbation.

A.1. Weighted Lebesgue spaces

In order to work in a general setting, we will define the weighted Besov spaces. To do so, let us define the class of admissible weight, following Triebel Chapter 6 [42]. In [33] one can find a more general definition for weights, which somehow allows the same kind of estimates.

Definition 37. We say that $w \in C^{\infty}(\mathbb{R}^d; \mathbb{R}_+ \setminus \{0\})$ is an admissible weight function if

(i) For all $\kappa \in \mathbb{N}^d$, there exists a positive constant c_{κ} such that

$$\forall x \in \mathbb{R}^d, \quad \partial^{\kappa} w(x) \le c_{\kappa} w(x).$$

(ii) There exists $\lambda' \ge 0$ and c > 0 such that

(A.1)
$$w(x) \le cw(y) \left(1 + |x - y|^2\right)^{\frac{\lambda}{2}}$$

Furthermore for any admissible weight we define the weighted L^p spaces as the following:

$$L^{p}(\mathbb{R}^{d};\mathbb{R}|w) = \left\{ f:\mathbb{R}^{d} \to \mathbb{R}: \|f\|_{L^{p}(\mathbb{R}^{d},w)} = \|wf\|_{L^{p}(\mathbb{R}^{d};\mathbb{R})} < +\infty \right\}.$$

A direct consequence of the definition is that the product of two admissible weights is also an admissible weight. Furthermore, standard polynomials weights are of course admissible, as proved in the following Proposition. Finally, Hölder inequality in weighted Lebesgue spaces is straightforward.

Proposition 38. Let $\lambda \in \mathbb{R}$, and let us define for all $x \in \mathbb{R}^d$, $\langle x \rangle = (1 + |x|^2)^{\frac{1}{2}}$. Then $\langle \cdot \rangle^{\lambda}$ is an admissible weight with $\lambda' = |\lambda|$.

Proof. Let us remark that for all $x, y \in \mathbb{R}^d$

$$(1 + y \cdot (x - y))^2 \le 2(1 + |y|^2|x - y|^2) \le 2(1 + |y|^2|x - y|^2)^2.$$

Hence,

$$1 + y \cdot (x - y) \le \sqrt{2} (1 + |y|^2 |x - y|^2)$$

then

$$\begin{aligned} \langle x \rangle^2 &= 1 + |x - y|^2 + |y|^2 + 2y \cdot (x - y) \\ &\leq |x - y|^2 + |y|^2 + 2\sqrt{2} (1 + |y|^2 |x - y|^2) \\ &\leq 2\sqrt{2} (1 + |x - y|^2 + |y|^2 + |y|^2 |x - y|^2) \\ &= 2\sqrt{2} \langle y \rangle^2 \langle x - y \rangle^2. \end{aligned}$$

and it therefore follows that

$$\langle x \rangle \le 2^{\frac{3}{4}} \langle y \rangle \langle x - y \rangle.$$

If $\lambda \geq 0$,

$$\langle x \rangle^{\lambda} \le 2^{\frac{\lambda 3}{4}} \langle y \rangle^{\lambda} \langle x - y \rangle^{\lambda},$$

and

$$\langle x \rangle^{-\lambda} = \langle x - y \rangle^{\lambda} (\langle x \rangle \langle y - x \rangle)^{-\lambda} \le 2^{\frac{3\lambda}{4}} \langle y \rangle^{-\lambda} \langle x - y \rangle^{\lambda},$$

which concludes the proof.

Lemma 39. Let w be an admissible weight and λ' defined as in Equation (A.1). Then for $1 \le p, q, r \le +\infty$ with $\frac{1}{p} + \frac{1}{q} = 1 + \frac{1}{r}$ and any measurable functions f and g,

$$\|f * g\|_{L^{p}(w)} \le \|f\|_{L^{p}(\langle \cdot \rangle^{\lambda'})} \|g\|_{L^{p}(w)}$$

Proof. The result is a direct consequence of the definition of admissible weights and of standard Young inequality. Indeed,

$$\|f * g\|_{L^{r}(w)} = \left(\int_{\mathbb{R}^{d}} \left| \int_{\mathbb{R}^{d}} f(y)g(x - y) \, \mathrm{d}y \right|^{r} w(x)^{r} \, \mathrm{d}x \right)^{\frac{1}{r}}$$

$$\lesssim \left(\int_{\mathbb{R}^{d}} \left| \int_{\mathbb{R}^{d}} \langle x - y \rangle^{\lambda'} f(x - y)w(y)g(y) \, \mathrm{d}y \right|^{r} \, \mathrm{d}x \right)^{\frac{1}{r}}$$

$$= \left\| \left(\langle \cdot \rangle^{\lambda'} f \right) * (wg) \right\|_{L^{r}}$$

$$\leq \left\| \langle \cdot \rangle^{\lambda'} f \right\|_{L^{p}} \|wg\|_{L^{q}}$$

$$= \|f\|_{L^{p}(\langle \cdot \rangle^{\lambda'})} \|g\|_{L^{q}(w)}$$

which is the desired result.

Furthermore, as usual in order to define Besov spaces we will work with functions with compactly supported Fourier transform. In order to deal with such functions, let us prove a Bernstein-type lemma in weighted L^p space:

Lemma 40. Let w be an admissible weight. Let $C = \{\xi \in \mathbb{R}^d : c_1 \le |\xi| \le c_2\}$ be an annulus and \mathcal{B} be a ball. There exists a constant C > 0 such that for all $1 \le p \le p' \le +\infty$, $n \ge 0$, $a \ge 1$, and for any function $f \in L^p(w)$, we have (i) If supp $\hat{f} \subset a\mathcal{B}$ then for any $k \in \mathbb{N}^d$

$$\|D^n f\|_{L^{p'}(w)} = \sup_{|k|=n} \|\partial^k f\|_{L^{p'}(w)} \le C^{n+1} a^{n+d(\frac{1}{p}-\frac{1}{p'})} \|f\|_{L^p(w)}.$$

(ii) If supp $\hat{f} \subset a\mathcal{A}$ then

$$\frac{1}{C^{n+1}}a^n \|f\|_{L^p(w)} \le \|D^n f\|_{L^p(w)} \le C^{n+1}a^n \|f\|_{L^p(w)}.$$

Here \hat{f} is the Fourier transform of f.

Proof. Let \hat{K} be a function such that $\hat{K} \equiv 1$ on \mathcal{B} and $\operatorname{supp}(\hat{K})$ is compactly supported and let us define $K_a = a^d K(a \cdot)$, and recall that \hat{f} denotes the Fourier transform of f. If $\operatorname{supp} \hat{f} \subset a\mathcal{B}$, then $f = K_a * f$, and $\partial^k f = (\partial^k K_a) * f = a^{|k|}((\partial^k K)_a) * f$, where $(\partial^k K)_a = a^d \partial^k K(a \cdot)$. Hence, by the previous weighted Young inequality, for $1 \le p \le p' \le +\infty$ and $\frac{1}{r} = 1 - (\frac{1}{p} - \frac{1}{p'})$,

$$\|\partial^k f\|_{L^{p'}(w)} = a^{|k|} \|(\partial^k K)_a * f\|_{L^{p'}(w)} \le a^{|k|} \|(\partial^k K)_a\|_{L^{r}(\langle \cdot \rangle^{\lambda'})} \|f\|_{L^{p}(w)}.$$

Furthermore, since $a \ge 1$ and $\lambda' \ge 0$, one has $\langle a^{-1}x \rangle^{\lambda'} \le \langle x \rangle^{\lambda'}$, and

$$\begin{split} \| \left(\partial^k K \right)_a \|_{L^r(\langle \cdot \rangle^{\lambda'})}^r &= a^{rd} \int_{\mathbb{R}^d} \left| \partial^k K(ax) \right|^r \langle x \rangle^{r\lambda'} \, \mathrm{d}x \\ &= a^{(r-1)d} \int_{\mathbb{R}^d} \left| \partial^k K(x) \right|^r \langle a^{-1}x \rangle^{r\lambda'} \, \mathrm{d}x \\ &\leq a^{(r-1)d} \int_{\mathbb{R}^d} \left| \partial^k K(x) \right|^r \langle x \rangle^{r\lambda'} \, \mathrm{d}x. \end{split}$$

It follows that

$$\left\|\left(\partial^k K\right)_a\right\|_{L^r(\langle\cdot\rangle^{\lambda'})} \lesssim a^{(1-\frac{1}{r})} \lesssim a^{d(\frac{1}{p}-\frac{1}{p'})}.$$

Gathering the above considerations proves (i).

The second inequality of (ii) is just a sub-case of (i). For the first inequality of the (ii), consider a smooth function L such that supp \hat{L} is included in an annulus and such that $\hat{L} \equiv 1$ on \mathcal{A} . Following [3] Lemma 2.1 and (1.23) page 25, there exists some real numbers (A_k) such that

$$|\xi|^{2n} = \sum_{|k|=n} A_k (-i\xi)^k (i\xi)^k,$$

with $(\xi_1, ..., \xi_d)^{(k_1, ..., k_d)} = \xi_1^{k_1} \cdots \xi_d^{k_d}$. Hence, we have

$$\sum_{|k|=n} A_k \frac{(-i\xi)^k}{|\xi|^{2n}} \hat{L}(a^{-1}\xi) \widehat{\partial^k f}(\xi) = L(a^{-1}\xi) \hat{f}(\xi) \sum_{|k|=n} A_k \frac{(-i\xi)^k (i\xi)^k}{|\xi|^{2n}} = \hat{f}(\xi).$$

For $k \in \mathbb{N}^d$ with |k| = n let us define

$$L^{k}(x) = A_{k} \int_{\mathbb{R}^{d}} (-i\xi)^{k} |\xi|^{-2n} \hat{L}(\xi) e^{i\xi \cdot x} dx$$

and we have $f = a^{-n} \sum_{|k|=n} L_a^k * \partial^k f$. One can use the weighted Young inequality to obtain that

$$\|f\|_{L^{p}(w)} \lesssim a^{-|n|} \sum_{|k|=n} \|L_{a}^{k}\|_{L^{1}(\langle\cdot\rangle^{\lambda'})} \|\partial^{k} f\|_{L^{p}(w)}$$

Furthermore, since $\langle a^{-1}x \rangle^{\lambda'} \leq \langle x \rangle^{\lambda'}$, it follows that $\|L_a^k\|_{L^1(\langle \cdot \rangle^{\lambda'})} \leq \|L^k\|_{L^1(\langle \cdot \rangle^{\lambda'})}$. Finally, we get

$$||f||_{L^{p}(w)} \lesssim a^{-|n|} ||D^{n}f||_{L^{p}(w)}$$

which proves our claim.

A.2. Weighted Besov spaces and standards estimates

Definition 41. Let $\mathcal{A} = \{\lambda \in \mathbb{R}^d : \frac{3}{4} \le |\lambda| \le \frac{8}{3}\}$. There exists two radial functions χ and φ such that $\operatorname{supp}(\chi) = B(0, \frac{3}{4})$, $\operatorname{supp}(\varphi) \subset \mathcal{A}$,

$$\forall \lambda \in \mathbb{R}^d \chi(\lambda) + \sum_{j \ge 0} \varphi(2^{-j}\lambda) = 1.$$

and for $j \ge 1$,

$$\operatorname{supp}(\chi) \cap \operatorname{supp}(\varphi(2^{-j}\cdot)) = \emptyset$$

and for $|j - j'| \ge 2$,

$$\operatorname{supp}(\varphi(2^{-j}\cdot)) \cap \operatorname{supp}(\varphi(2^{-j'}\cdot)) = \varnothing$$

For all $f \in S'$ and $j \ge 0$, we define the in-homogeneous Paley-Littlewood blocks by

$$\Delta_{-1}f = \mathcal{F}^{-1}(\chi \hat{f}), \text{ and } \Delta_j f = \mathcal{F}^{-1}(\varphi(2^{-j}\cdot)\hat{f}),$$

where \hat{f} denotes the Fourier transform of f and \mathcal{F}^{-1} the inverse Fourier transform.

Note that the Paley-Litllewood blocks define a nice appproximation of the unity. We refer to [3] proposition 2.12 for a proof.

Proposition 42. For all $f \in S'$, let us define for all $j \ge -1$ $S_j f = \sum_{j' \le j-1} \Delta_j f$. Then $f = \lim_{j \to \infty} S_j f \quad in \quad S'.$

Definition 43. Let $1 \le p, q \le +\infty$, and let $\kappa \in \mathbb{R}$ and w be an admissible weight. For a distribution $f \in \mathcal{S}'(\mathbb{R}^d; \mathbb{R})$ we define the (in-homogeneous) weighted Besov norm by

$$\|f\|_{B_{p,q}^{\kappa}(w)} = \left\| \left(2^{\kappa_{j}} \|\Delta_{j} f\|_{L^{p}(\mathbb{R}^{d},w)} \right)_{j \geq -1} \right\|_{\ell^{q}(\mathbb{N} \cup \{-1\})}$$

where $||f||_{L^p(\mathbb{R}^d,w)} = (\int_{\mathbb{R}^d} f(x)^p w(x)^p dx)^{\frac{1}{p}}$. When $w \equiv 1$, we only write $B_{p,q}^{\kappa}$.

We gather here some basic properties of (weighted) Besov spaces.

Proposition 44. Let w be an admissible weight.

(i) The space

$$B_{p,q}^{s}(w) = \left\{ f \in \mathcal{S}'(\mathbb{R}^{d}; \mathbb{R}) : \|f\|_{B_{p,q}^{\kappa}(w)} < +\infty \right\}$$

does not depend on the choice of φ and χ .

(ii) The two following quantities $||f||_{B_{p,q}^{\kappa}(w)}$ and $||wf||_{B_{p,q}^{\kappa}}$ are equivalent norms on $B_{p,q}^{\kappa}(w)$.

- (iii) For all $n \ge 0$, $||D^n f||_{B^{\kappa}_{p,q}(w)} \lesssim ||f||_{B^{\kappa+n}_{p,q}(w)}^{r''}$.
- (iv) Let $1 \le p \le p' \le +\infty$ and $1 \le q \le q' \le +\infty$, then for all $\varepsilon > 0$,

$$\|f\|_{B^{\kappa-d(\frac{1}{p}-\frac{1}{p'})}_{p',q}(w)} \lesssim \|f\|_{B^{\kappa}_{p,q}(w)} \lesssim \|f\|_{B^{\kappa}_{p',q}(\langle\cdot\rangle^{d(\frac{1}{p}-\frac{1}{p'})+\varepsilon}w)}$$

and

$$\|f\|_{B_{p,\infty}^{\kappa}(w)} \lesssim \|f\|_{B_{p,q}^{\kappa}(w)} \lesssim \|f\|_{B_{p,q'}^{\kappa-\varepsilon}(w)}.$$

(v) For all $\varepsilon, \delta > 0$ and all $\kappa \in \mathbb{R}$ and all $1 \le p, q \le +\infty$,

$$B_{p,q}^{\kappa}(w)$$
 is compactly embedded in $B_{p,q}^{\kappa-\varepsilon}(\langle\cdot\rangle^{-\delta}w)$

(vi) Suppose that $\kappa > 0$ and $\kappa \notin \mathbb{N}$ and for $f : \mathbb{R}^d \mapsto \mathbb{R}$ let us define

$$\|f\|_{\mathcal{C}^{\kappa}(w)} = \sum_{|k| \le [\kappa]} \sup_{x \in \mathbb{R}^d} |w(x)\partial^k f(x)| + \sum_{|k|=[\kappa]} \sup_{0 < |h| \le 1} \sup_{x \in \mathbb{R}^d} \frac{w(x)|\partial^k f(x+h) - \partial^k f(x)|}{|h|^{\kappa - [\kappa]}}.$$

Then

$$\mathcal{C}^{\kappa}(w) = \left\{ f : \|f\|_{\mathcal{C}^{\kappa}(w)} < +\infty \right\} = B^{\kappa}_{\infty,\infty}(w)$$

and furthermore $\|\cdot\|_{\mathcal{C}^{\kappa}(w)}$ and $\|\cdot\|_{B^{\kappa}_{\infty,\infty}(w)}$ are equivalent norms on this space.

Proof. We only prove the weighted inequality in the fourth point, and we refer [42] and the references therein for the other ones. The first and the third inequalities are direct consequences of Lemma 40. For the second one, let us take $1 \le p < p' < +\infty$. We have, thanks to Jensen inequality, for any $\varepsilon > 0$

$$\begin{split} \|\Delta_j f\|_{L^p(w)}^{p'} \lesssim_{d,\varepsilon} \int_{\mathbb{R}^d} |\Delta_j f(x)|^{p'} w(x)^{p'} \langle x \rangle^{\frac{(d+\varepsilon)p'}{p}} \langle x \rangle^{-\frac{(d+\varepsilon)p'}{p'}} \, \mathrm{d}x \\ &= \int_{\mathbb{R}^d} |\Delta_j f(x)|^{p'} \big(w(x) \langle x \rangle^{(d+\varepsilon)(\frac{1}{p} - \frac{1}{p'})} \big)^{p'} \, \mathrm{d}x. \end{split}$$

The constant in the previous inequality does not depends on p'. This gives

$$\|\Delta_j f\|_{L^p(w)} \lesssim_{d,\varepsilon} \|\Delta_j f\|_{L^{p'}(\langle \cdot \rangle^{(d+\varepsilon)(\frac{1}{p}-\frac{1}{p'})}w)}.$$

Finally, in order to deal with product of elements in Besov space, we give the following result, which can be proved thanks to standard techniques (see [3] Lemma 2.69 and 2.84 and [33] Theorem 3.17 and Corollary 3.19 and 3.21). Note that it mostly relies on Hölder and Young inequality, and therefore is available in the context of weighted spaces.

Proposition 45 (Corollary 2.86 in [3] and Corollary 3.19 in [33]). Let w be an admissible weight, $\kappa_2 \leq \kappa_1$ with $\kappa_1 \geq 0$ and suppose that $\kappa_1 + \kappa_2 > 0$, and let $1 \leq p$, p_1 , p_2 , $q \leq +\infty$ such that $\frac{1}{p} = \frac{1}{p_1} + \frac{1}{p_2}$. Let ε , $\delta > 0$. Then for all $f \in \mathcal{B}_{p_1,q}^{\kappa_1}(w)$ and all $g \in \mathcal{B}_{p_2,q}^{\kappa_2}(w)$,

- (i) $(S_j f S_j g)_{j \ge 0}$ converges in $B_{p,q}^{\kappa_2 \varepsilon}(\langle \cdot \rangle^{-\delta} w)$ to a limit in $B_{p,q}^{\kappa_2}(w)$.
- (ii) We have

$$\left\|\lim_{j} \mathcal{S}_{j} f \mathcal{S}_{j} g\right\|_{B^{\kappa_{2}}_{p,q}(w)} \lesssim \|f\|_{B^{\kappa_{1}}_{p,q}(w)} \|g\|_{B^{\kappa_{2}}_{p,q}(w)}$$

(iii) When $\kappa_2 \ge 0$, $\lim_j S_j f S_j g = fg$ the standard point wise product.

Remark 46. The previous proposition shows that the limit does not depend on the choice of the blocks, and therefore it extends canonically the notion of product of functions to a product of distributions, as soon as $\kappa_1 + \kappa_2 > 0$, with $\kappa_1 \ge 0$. For this reason, we will denote by $fg = \lim_j S_j f S_j g$, and this is a bi-linear functional from $B_{p_1,q}^{\kappa_1}(w) \times B_{p_2,q}^{\kappa_2}(w)$ to $B_{p,q}^{\kappa_2}(w)$.

A.3. Heat kernel estimates

In order to deal with heat kernel estimates on weighted Besov spaces, we first need some heat kernel estimates for function whose Fourier transform has a support in an annulus. In order to deal with non-Gaussian noise, we need to consider heat semi-group for fractional Laplacian. For a full study of the fractional Laplacian, we refer to [29]. Here we just define the fractional Laplacian for smooth functions.

Definition 47. Let $\alpha \in (0, 2]$. For any function $f \in S(\mathbb{R}^d; \mathbb{R})$ we define the fractional Laplace operator $\Delta^{\frac{\alpha}{2}} = -(-\Delta)^{\frac{\alpha}{2}}$ by

$$\Delta^{\frac{\alpha}{2}} f = \mathcal{F}^{-1} \left(-|\cdot|^{\alpha} \hat{f}(\cdot) \right)$$

Furthermore, we define the semi-group associated to $\Delta^{\frac{\alpha}{2}}$

$$P_t^{\frac{\alpha}{2}}f = \mathcal{F}^{-1}\left(e^{-|\cdot|^{\alpha}t}\hat{f}(\cdot)\right).$$

Finally we extend this definition to the whole space S' by the standard procedure.

Note that when $\alpha = 2$, the previous definition gives the standard Laplace operator. When this is the case we simply write *P* instead of *P*¹.

Proposition 48. Let w be an admissible weight and λ' be defined as in (A.1). Let $\alpha \in (0, 2]$. Let $\mathcal{A} = \{\xi \in \mathbb{R}^d : c_1 \leq |\xi| \leq c_2\}$ be an annulus, let $a \geq 1$ and let $f : \mathbb{R}^d \to \mathbb{R}^d$ be a function such that $supp \hat{f} \subset a\mathcal{A}$. Let $0 < s \leq t$. There exists a constant c > 0 such that for all $p \geq 1$,

$$\left\|P_t^{\frac{\alpha}{2}}f\right\|_{L^p(w)} \lesssim e^{-cta^{\alpha}} \|f\|_{L^p(w)},$$

and for all $\rho \geq 0$,

$$\|(P_t - P_s)f\|_{L^p(w)} \lesssim a^{-\alpha\rho} \left|\frac{1}{s^{\rho}} - \frac{1}{t^{\rho}}\right| \|f\|_{L^p(w)}$$

Finally, for all $s \le u \le \tau' \le \tau$, and for all $\rho > 0$,

$$\begin{split} \left\| \left((P_{\tau-s} - P_{\tau-u}) - (P_{\tau'-s} - P_{\tau'-u}) \right) f \right\|_{L^{p}(w)} \\ \lesssim a^{-\alpha\rho} \left(\frac{1}{(\tau'-u)^{\rho}} - \frac{1}{(\tau'-s)^{\rho}} - \frac{1}{(\tau-u)^{\rho}} + \frac{1}{(\tau-s)^{\rho}} \right) \| f \|_{L^{p}(w)} \end{split}$$

Proof. We follow the proof of [3], Lemma 2.4. The first affirmation is the result of this lemma but in the context of weighted spaces and of fractional operator. We also refer to [33] Lemma 2.10 to a proof. Since the proof of the second and third points are similar to the proof of the first one, we will not detail the first point. For the second and third one, let us define

$$E(s,t;\xi) := e^{-|\xi|^{\alpha}t} - e^{-|\xi|^{\alpha}s}$$

Note that in that case, we have

$$(P_t - P_s)f(x) = \int_{\mathbb{R}^d} E(s, t; \xi)\hat{f}(\xi)e^{i\xi x} \,\mathrm{d}x.$$

Thanks to the hypothesis, there also exists a smooth function $\varphi : \mathbb{R}^d \to \mathbb{R}^d$ such that $\varphi \equiv 1$ on \mathcal{A} and $\operatorname{supp} \varphi \subset \{\xi \in \mathbb{R}^d : \frac{c_1}{2} \le |\xi| \le 2c_2\}$. Let us define for all $s \le t$

$$K(s,t;x) = \int E(s,t;\xi)\varphi(\xi)e^{i\xi x}d\xi$$

and $K_a(s, t, x) = a^d K(s, t, ax)$. Hence we have

$$(P_t - P_s)f(x) = \int_{\mathbb{R}^d} E(s, t; \xi)\hat{f}(\xi)e^{i\xi \cdot x}d\xi$$
$$= \int_{\mathbb{R}^d} E(s, t; \xi)\varphi(a^{-1}\xi)\hat{f}(\xi)e^{i\xi \cdot x}d\xi$$
$$= a^d \int_{\mathbb{R}^d} E(s, t; a\xi)\varphi(\xi)\hat{f}(\xi)e^{i\xi \cdot ax}d\xi.$$

And since $E(s, t, a\xi) = E(a^{\alpha}s, a^{\alpha}t; \xi)$, one has

$$(P_t - P_s) * f(x) = K_a(a^{\alpha}s, a^{\alpha}t; \cdot) * f(x)$$

Finally, thanks to Young inequality, one has

$$\|(P_t - P_s) * f\|_{L^p(w)} \le \|K_a(a^{\alpha}s, a^{\alpha}t; \cdot)\|_{L^1(\langle \cdot \rangle^{\lambda'})} \|f\|_{L^p(w)}.$$

Note also that since $\lambda' \ge 0$ and $a \ge 1$, we have $\langle a^{-1}x \rangle^{\lambda'} \le \langle x \rangle^{\lambda'}$, and

$$\|K_a(a^{\alpha}s, a^{\alpha}t; \cdot)\|_{L^1(\langle \cdot \rangle^{\lambda'})} = a^d \int_{\mathbb{R}^d} K_{a^{\alpha}s, a^{\alpha}t, ax} \langle x \rangle^{\lambda'} dx$$
$$= \int_{\mathbb{R}^d} K_{a^{\alpha}s, a^{\alpha}t, ax} \langle a^{-1}x \rangle^{\lambda'} dx$$
$$= \leq a^d \int_{\mathbb{R}^d} K_{a^{\alpha}s, a^{\alpha}t, x} \langle x \rangle^{\lambda'} dx$$
$$= \|K(a^{\alpha}s, a^{\alpha}t; \cdot)\|_{L^1(\langle \cdot \rangle^{\lambda'})}$$

Hence, it is enough to prove the proposition for a = 1 and for all $0 < s' \le t'$, and then specify $s' = a^{\alpha}s$ and $t' = a^{\alpha}t$.

Let $M \in \mathbb{N}$ such that $2M > d + \lambda'$. In order to prove the proposition, it is then enough to bound $|(1 + |x|^2)^M K(s, t; x)|$ by the wanted quantity $\frac{1}{x^{\rho}} - \frac{1}{t^{\rho}}$. We have

$$(1+|x|^2)^M K(s,t;x) = \int_{\mathbb{R}^d} ((1-\Delta)^M e^{ix\cdot})(\xi)\varphi(\xi)E(s,t;\xi)d\xi$$
$$= \int_{\mathbb{R}^d} e^{ix\xi}(1-\Delta)^M (\varphi(\cdot)E(s,t;\cdot))(\xi)d\xi$$

And thanks to the Faà di Bruno formula, there exists constants $(c_{\nu,\kappa})$ where ν and κ are multi-indices such that

$$(1+|x|^2)^M K(s,t;x) = \sum_{|\nu|+|\kappa| \le 2M} c_{\nu,\kappa} \int_{\mathbb{R}^d} e^{i\xi x} \partial^{\nu} \varphi(\xi) \partial^{\kappa} E(s,t;\xi) d\xi.$$

Hence, in order to prove the proposition, one only has to bound all the derivatives up to order 2*M* of $E(s, t; \cdot)$ for $\xi \in A$ by the wanted quantity $\frac{1}{s^{\rho}} - \frac{1}{t^{\rho}}$.

Note that the same strategy could be used in the case where we have the rectangular increment of the semi-group if we replace $E(s, t; \xi)$ by

$$F(s, u, \tau', \tau; \xi) := \left(e^{-|\xi|^{\alpha}(\tau-s)} - e^{-|\xi|^{\alpha}(\tau-u)}\right) - \left(e^{-|\xi|^{\alpha}(\tau'-s)} - e^{-|\xi|^{\alpha}(\tau'-u)}\right).$$

The same partial conclusion holds, one only has to control all the derivatives of $F(s, u, \tau', \tau; \cdot)$ up to order 2*M* for $\xi \in A$. Furthermore, observe that

$$E(s,t;\xi) = \int_s^t -|\xi|^{\alpha} e^{-r|\xi|^{\alpha}} \,\mathrm{d}r.$$

Hence, by a direct induction and since $\xi \in A$, for every multiindexe $k = (k_1 \cdots, k_d) \in \mathbb{N}^d$, there exists a polynomial

$$P_{k,\xi}(t) = \sum_{l=0}^{|k|} a_l^k(\xi) t^l,$$

where for all $l \in \{0, ..., n\}$, $\xi \to a_l^k(\xi)$ are non-negative smooth functions on \mathcal{A} , and for all $\xi \in \mathcal{A}$, $a_n^k(\xi) \neq 0$ and such that

$$\partial^k \left(|\cdot|^2 e^{-r|\cdot|^{\alpha}} \right)(\xi) = P_{k,\xi}(r) e^{-r|\xi|^{\alpha}}$$

Furthermore, since $c_1 \le |\xi| \le c_2$, there exists a constant c > 0 (depending on \mathcal{A}) such that

$$\left|\partial^{k}\left(|\cdot|^{\alpha}e^{-r|\cdot|^{\alpha}}\right)(\xi)\right| \lesssim \tilde{P}_{k,\xi}(r)e^{-r|\xi|^{2}} \lesssim_{k,\mathcal{A}} e^{-cr} \lesssim r^{-\rho-1}$$

for any $\rho \geq -1$. Hence

$$\left|\partial^k E(s,t;\cdot)\right| \lesssim \int_s^t r^{-\rho} \,\mathrm{d}r \lesssim \frac{1}{s^{\rho}} - \frac{1}{t^{\rho}}.$$

With the previous discussion, this gives the wanted result, when we replace t by $a^{\alpha}t$ and s by $a^{\alpha}s$. In order to deal with the last estimates, one only has to remember that

$$F(s, u, \tau', \tau; \xi) = -\int_{u}^{s} |\xi|^{\alpha} \left(e^{-|\xi|^{\alpha}(\tau-r)} - e^{-|\xi|^{\alpha}(\tau'-r)}\right) \mathrm{d}r$$
$$= \int_{s}^{u} \int_{\tau'-r}^{\tau-r} |\xi|^{2\alpha} e^{-|\xi|^{\alpha}v} \,\mathrm{d}v \,\mathrm{d}r$$

The same argument as before gives the bound

$$\left|\partial^k F(s, u, \tau', \tau; \cdot)\right| \lesssim \int_s^u \int_{\tau'-r}^{\tau-r} e^{-cv} \,\mathrm{d}v \,\mathrm{d}r$$

Finally, for any $\rho \ge 0$,

$$\begin{aligned} \left|\partial^k F(s, u, \tau', \tau; \cdot)\right| &\lesssim \int_s^u \int_{\tau'-r}^{\tau-r} v^{-(\rho+2)} \,\mathrm{d}v \,\mathrm{d}r \\ &\lesssim \frac{1}{(\tau'-u)^\rho} - \frac{1}{(\tau'-s)^\rho} - \frac{1}{(\tau-u)^\rho} + \frac{1}{(\tau-s)^\rho} \end{aligned}$$

Again, this give the wanted result when recalling that one must replace τ , τ' , u and s by $a^{\alpha}\tau$, $a^{\alpha}\tau'$, $a^{\alpha}u$ and $a^{\alpha}s$.

Let us gives the following useful and straightforward corollary for the action of the fractional heat semi-group on weighted Besov spaces. Let us first remind a rather elementary but useful lemma ([11], Lemma 4.4). For the sake of the reader, we provide a full proof of it.

Lemma 49. Let $\rho \ge 0$ and $\theta \in [0, 1]$. There is a constant c > 0 such that for any $0 < s \le t$,

$$\frac{1}{s^{\rho}} - \frac{1}{t^{\rho}} \le c(t-s)^{\theta} s^{-(\rho+\theta)}.$$

Let $\rho \ge 0$ and $\theta, \theta' \in [0, 1]$. There exists a constant such that

$$\frac{1}{(\tau'-u)^{\rho}} - \frac{1}{(\tau'-s)^{\rho}} - \frac{1}{(\tau-u)^{\rho}} + \frac{1}{(\tau-s)^{\rho}} \le c(\tau-\tau')^{\theta}(u-s)^{\theta'}(\tau'-u)^{-(\rho+\theta+\theta')}$$

Proof. Let us remark that

$$\frac{1}{s^{\rho}} - \frac{1}{t^{\rho}} = (1+\rho) \int_{s}^{t} r^{-(\rho+1)} \,\mathrm{d}r \le (1+\rho)(t-s)s^{-(\rho+1)}$$

The result follows by a standard interpolation between the two inequalities. For the second inequality, set

$$B = \frac{1}{(\tau' - u)^{\rho}} - \frac{1}{(\tau' - s)^{\rho}} - \frac{1}{(\tau - u)^{\rho}} + \frac{1}{(\tau - s)^{\rho}}$$

We have, thanks to the same integral representation,

$$B \lesssim (\tau' - u)^{-\rho},$$

$$B \lesssim (\tau - \tau')(\tau' - u)^{-(\rho+1)},$$

$$B \lesssim (u - s)(\tau' - u)^{-(\rho+1)},$$

and

$$B \lesssim (u-s) \big(\tau - \tau'\big) \big(\tau' - u\big)^{-(\rho+2)}.$$

Let us suppose, without loss of generality that $\theta' \le \theta$ that is $\theta' = \alpha \theta$ for some $0 \le \alpha < 1$. We have, using the first and the last inequalities,

$$B \lesssim (\tau - \tau')^{\theta} (u - s)^{\theta} (\tau' - u)^{-(\rho + 2\theta)}.$$

Using the first and the second one, we have

$$B \lesssim (\tau - \tau')^{ heta} (\tau' - u)^{-(
ho + 2 heta)}.$$

We interpolate those last two inequalities to have

$$B \lesssim (\tau - \tau')^{\theta} (u - s)^{\alpha \theta} (\tau' - u)^{-(\rho + 2\alpha \theta + (1 - \alpha)\theta)} = (\tau - \tau')^{\theta} (u - s)^{\theta'} (\tau' - u)^{-(\rho + \theta + \theta')}.$$

Corollary 50. Let w be an admissible weight and $\alpha \in (0, 2]$. Let $\kappa \in \mathbb{R}$ and let $1 \le p, q \le +\infty$. Let $0 \le t$, then for all $\rho \ge 0$,

$$\left\|P_t^{\frac{\alpha}{2}}f\right\|_{B_{p,q}^{\kappa+\alpha\rho}(w)} \lesssim t^{-\rho} \|f\|_{B_{p,q}^{\kappa}(w)}$$

For all $\theta \in [0, 1]$ and all $\rho > 0$,

$$\| \left(P_t^{\frac{\alpha}{2}} - P_s^{\frac{\alpha}{2}} \right) f \|_{B_{p,q}^{\kappa+\alpha\rho}(w)} \lesssim |(t-s)^{\theta} s^{-(\rho+\theta)} \| f \|_{B_{p,q}^{\kappa}(w)}$$

Furthermore, for any $\theta, \theta' \in [0, 1]$ *and all* $\rho \ge 0, s \le u < \tau' \le \tau$,

$$\left\| \left(\left(P_{\tau-s}^{\frac{\alpha}{2}} - P_{\tau-u}^{\frac{\alpha}{2}} \right) - \left(P_{\tau'-s}^{\frac{\alpha}{2}} - P_{\tau'-u}^{\frac{\alpha}{2}} \right) \right) f \right\|_{B_{p,q}^{\kappa+\alpha\rho}(w)} \lesssim \left(\tau - \tau' \right)^{\theta} (u-s)^{\theta'} \left(\tau' - u \right)^{-(\rho+\theta+\theta')} \| f \|_{B_{p,q}^{\kappa}(w)} \leq \left(\tau - \tau' \right)^{\theta} (u-s)^{\theta'} \left(\tau' - u \right)^{-(\rho+\theta+\theta')} \| f \|_{B_{p,q}^{\kappa}(w)} \leq \left(\tau - \tau' \right)^{\theta} (u-s)^{\theta'} \left(\tau' - u \right)^{-(\rho+\theta+\theta')} \| f \|_{B_{p,q}^{\kappa}(w)} \leq \left(\tau - \tau' \right)^{\theta} (u-s)^{\theta'} \left(\tau' - u \right)^{-(\rho+\theta+\theta')} \| f \|_{B_{p,q}^{\kappa}(w)} \leq \left(\tau - \tau' \right)^{\theta} (u-s)^{\theta'} \left(\tau' - u \right)^{-(\rho+\theta+\theta')} \| f \|_{B_{p,q}^{\kappa}(w)} \leq \left(\tau - \tau' \right)^{\theta} (u-s)^{\theta'} \left(\tau' - u \right)^{-(\rho+\theta+\theta')} \| f \|_{B_{p,q}^{\kappa}(w)} \leq \left(\tau - \tau' \right)^{\theta} \left($$

Proof. The first bound is a direct consequence of the first bound of Proposition 48. Indeed, for all $j \ge 0$,

$$\|\Delta_{j}P_{t}^{\frac{\alpha}{2}}f\|_{L^{p}(w)} = \|P_{t}(\Delta_{j}f)\|_{L^{p}(w)} \lesssim e^{-c2^{\alpha j}t} \|\Delta_{j}f\|_{L^{p}(w)} \lesssim 2^{-\alpha\rho j}t^{-\rho} \|\Delta_{j}f\|_{L^{p}(w)},$$

and the result follows via the definition of weighted Besov spaces. For the second bound, let us remark that we have for all $0 < s \le t$, and thanks to Proposition 48,

$$\|\Delta_{j} \left(P_{t}^{\frac{\alpha}{2}} - P_{s}^{\frac{\alpha}{2}}\right) f\|_{L^{p}(w)} = \lesssim \left(\frac{1}{s^{\rho}} - \frac{1}{s^{\rho}}\right) 2^{-\alpha\rho j} \|\Delta_{j} f\|_{L^{p}(w)} \lesssim 2^{-\alpha\rho j} (t-s)^{\theta} s^{-(\rho+\theta)} \|\Delta_{j} f\|_{L^{p}(w)}.$$

where the last inequality comes from the previous lemma. This allows us to derive the result by using the definition of B spaces. Finally, note that we have thanks to Proposition 48 and Lemma 49,

$$\left|\Delta_{j}\left(\left(P_{\tau-u}^{\frac{\alpha}{2}}-P_{\tau-s}^{\frac{\alpha}{2}}\right)-\left(P_{\tau'-u}^{\frac{\alpha}{2}}-P_{\tau'-s}^{\frac{\alpha}{2}}\right)\right)f\right\|_{L^{p}(w)} \lesssim 2^{-\alpha\rho j}\left(\tau-\tau'\right)^{\theta}(u-s)^{\theta'}\left(\tau'-u\right)^{\rho+\theta+\theta'}\|\Delta_{j}f\|_{L^{p}(w)}.$$

And again one can conclude with the definition of the weighted Besov spaces.

Appendix B: Cauchy–Lipschitz theorem for mSHE in standard case

We give a short proof of local well-posedness of the mSHE in a simple context. For more on this, one can consult [32] and the generalization for more irregular noise with more involve techniques [20] and [22].

Theorem 51. Let $\alpha \in (0, 2]$. Let $\vartheta \in (0, \min(1, \alpha))$, let $\vartheta < \beta < \alpha - \vartheta$. Let $g \in C^2$ and let $u_0 \in C^\beta$ and $\xi \in C^{-\vartheta}$. There exists a unique local solution of the equation

$$\partial_t u = \Delta^{\frac{\alpha}{2}} u + g(u)\xi$$

in the mild form

$$u(t,\cdot) = P_t^{\frac{\alpha}{2}} u_0 + \int_s^t P_{t-s}^{\frac{\alpha}{2}} \xi g\left(u(s,\cdot)\right) \mathrm{d}s$$

Proof. First, let us take $u \in C^{\beta}$. We have

$$\|g(u(t,\cdot)\|_{\mathcal{C}^{\beta}} \le \|g\|_{\mathcal{C}^{2}} \|u(t,\cdot)\|_{\mathcal{C}^{\beta}}$$

For $u \in C([0, T]; C^{\beta})$, let us remark that thanks to standard Bony estimates in Besov-Hölder spaces (Proposition 45)

$$\|\xi g(u(t,\cdot)\|_{\mathcal{C}^{-\vartheta}} \lesssim \|g\|_{\mathcal{C}^2} \|\xi\|_{\mathcal{C}^{-\vartheta}} \|u(t,\cdot)\|_{\mathcal{C}^{\beta}}.$$

Finally for $0 \le s < t \le T$,

$$\left\|P_{t-s}\xi g(u(t,\cdot))\right\|_{\mathcal{C}^{\beta}} \lesssim \frac{1}{(t-s)^{\frac{\beta+\rho}{\alpha}}} \|g\|_{\mathcal{C}^{2}} \|\xi\|_{\mathcal{C}^{-\vartheta}} \|u(t,\cdot)\|_{\mathcal{C}^{\beta}}.$$

Hence, the application Γ

$$\Gamma(u)(t,x) = P_t u_0 + \int_0^t P_{t-s} \xi g(u(s,\cdot)) \,\mathrm{d}s$$

is well-defined from $C^0([0, T]; C^\beta)$ to itself. Furthermore, Let us remark that for $u, v \in C^\beta$,

$$\left\|g(u)-g(v)\right\|_{\mathcal{C}^{\beta}}\lesssim \|g\|_{\mathcal{C}^{2}}\|u-v\|_{\mathcal{C}^{\beta}},$$

hence for $u, v \in C^0([0, T]; \mathcal{C}^\beta)$,

$$\sup_{t\in[0,T]} \left\| \Gamma(u)(t,\cdot) - \Gamma(v)(t,\cdot) \right\|_{\beta} \lesssim \int_{0}^{t} \frac{1}{(t-s)^{\frac{\beta+\rho}{\alpha}}} \|g\|_{\mathcal{C}^{2}} \|\xi\|_{\mathcal{C}^{-\vartheta}} \left\| u(t,\cdot) - v(t,\cdot) \right\|_{\mathcal{C}^{\beta}} \mathrm{d}s$$
$$\lesssim T^{1-\frac{\beta+\rho}{\alpha}} \sup_{t\in[0,T]} \left\| u(t,\cdot) - v(t,\cdot) \right\|_{\mathcal{C}^{\beta}}.$$

Hence, be standard Schauder fixed point, for T small enough, there is a unique $u \in C^0([0, T]; C^\beta)$.

Acknowledgements

FH gratefully acknowledges financial support from the STORM project 274410, funded by the Research Council of Norway. We would like thank Prof. Nicolas Perkowski for several fruitful discussions on this topic, and the two anonymous referees whose comments and suggestions greatly improved this article.

References

- S. Athreya, O. Butkovsky, K. Lê and L. Mytnik. Well-posedness of stochastic heat equation with distributional drift and skew stochastic heat equation, 2020. Available at arXiv:2011.13498 [math].
- [2] H. Bahouri, J.-Y. Chemin and R. Danchin. Fourier Analysis and Nonlinear Partial Differential Equations. Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences] 343. Springer, Heidelberg, 2011. MR2768550 https://doi.org/10.1007/ 978-3-642-16830-7

- H. Bahouri, J.-Y. Chemin and R. Danchin. Fourier Analysis and Nonlinear Partial Differential Equations. Springer, Berlin, 2011. MR2768550 https://doi.org/10.1007/978-3-642-16830-7
- [4] I. Bailleul, A. Debussche and M. Hofmanová. Quasilinear generalized parabolic Anderson model equation. Stoch. Partial Differ. Equ. Anal. Comput. 7 (1) (2019) 40–63. MR3916262 https://doi.org/10.1007/s40072-018-0121-1
- [5] C. Bellingeri, P. K. Friz and M. Gerencsér. Singular Paths Spaces and Applications, 2020.
- [6] J. Benedikt, V. E. Bobkov, P. Girg, L. Kotrla and P. Takac. Nonuniqueness of solutions of initial-value problems for parabolic p-Laplacian. Electron. J. Differential Equations 38 (2015), 2015. MR3335768
- [7] O. Butkovsky and L. Mytnik. Regularization by noise and flows of solutions for a stochastic heat equation. Ann. Probab. 47 (1) (2019) 165–212. MR3909968 https://doi.org/10.1214/18-AOP1259
- [8] R. Catellier and M. Gubinelli. Averaging along irregular curves and regularisation of ODEs. *Stochastic Process. Appl.* 126 (8) (2016) 2323–2366. MR3505229 https://doi.org/10.1016/j.spa.2016.02.002
- [9] L. Coutin, R. Duboscq and A. Réveillac. The Itô-Tanaka Trick: A non-semimartingale approach, 2019.
- [10] A. M. Davie. Differential equations driven by rough paths: An approach via discrete approximation. Appl. Math. Res. Express. AMRX 40 (2007) abm009. MR2387018
- [11] A. Deya and S. T. Rough. Volterra equations. I. The algebraic integration setting. *Stoch. Dyn.* **9** (3) (2009) 437–477. MR2566910 https://doi.org/10. 1142/S0219493709002737
- [12] P. Friz. Mini-Course on Rough Paths, 2009.
- [13] P. K. Friz and M. Hairer. A Course on Rough Paths: With an Introduction to Regularity Structures. Universitext. Springer, Cham, 2014. MR3289027 https://doi.org/10.1007/978-3-319-08332-2
- [14] H. Fujita and S. Watanabe. On the uniqueness and non-uniqueness of solutions of initial value problems for some quasi-linear parabolic equations. *Comm. Pure Appl. Math.* 21 (6) (1968) 631–652. MR0234129 https://doi.org/10.1002/cpa.3160210609
- [15] L. Galeati. Nonlinear young differential equations: A review. J. Dynam. Differential Equations (2021).
- [16] L. Galeati and M. Gubinelli. Prevalence of rho-irregularity and related properties, 2020.
- [17] L. Galeati and M. Gubinelli. Noiseless regularisation by noise. Rev. Mat. Iberoam. (2021). MR4404773 https://doi.org/10.4171/rmi/1280
- [18] L. Galeati and F. A. Harang. Regularization of multiplicative sdes through additive noise, 2020.
- [19] M. Gubinelli, P. Imkeller and N. Perkowski. Paracontrolled distributions and singular PDEs. Forum Math. Pi 3 (2015), e6. MR3406823 https://doi.org/10.1017/fmp.2015.2
- [20] M. Gubinelli, P. Imkeller and N. Perkowski. A Fourier analytic approach to pathwise stochastic integration. *Electron. J. Probab.* 21 (2016). MR3485344 https://doi.org/10.1214/16-EJP3868
- [21] M. Hairer and X.-M. Li. Averaging dynamics driven by fractional Brownian motion. Ann. Probab. 48 (4) (2020) 1826–1860. MR4124526 https://doi.org/10.1214/19-AOP1408
- [22] M. Hairer and N. S. Pillai. Regularity of laws and ergodicity of hypoelliptic SDEs driven by rough paths. Ann. Probab. 41 (4) (2013) 2544–2598. MR3112925 https://doi.org/10.1214/12-AOP777
- [23] F. Harang and A. Mayorcas. Pathwise regularisation of singular interacting particle systems and their mean field limits, 2020.
- [24] F. A. Harang and C. Ling. Regularity of local times associated with Volterra–Lévy processes and path-wise regularization of stochastic differential equations. J. Theor. Probab. (2021).
- [25] F. A. Harang and N. Perkowski. C infinity regularization of odes perturbed by noise. Stoch. Dyn. (2021), 2140010. MR4342752 https://doi.org/10. 1142/S0219493721400104
- [26] F. A. Harang and S. Tindel. Volterra equations driven by rough signals. Stochastic Process. Appl. 142(C) (2021) 34–78. MR4309474 https://doi.org/10.1016/j.spa.2021.08.001
- [27] Y. Hu, D. Nualart and J. Song. A nonlinear stochastic heat equation: Hölder continuity and smoothness of the density of the solution. In Stochastic Processes and their Applications 1083–1103, 123, 2013. MR3005015 https://doi.org/10.1016/j.spa.2012.11.004
- [28] T. Hytönen, J. van Neerven, M. Veraar and L. Weis. Analysis in Banach Spaces: Volume I: Martingales and Littlewood–Paley Theory, 1st edition. Springer, Berlin, 2016. MR3617205
- [29] M. Kwaśnicki. Ten equivalent definitions of the fractional Laplace operator. In Fractional Calculus and Applied Analysis 7–51. Fractional Calculus and Applied Analysis. 20. De Gruyter, Berlin, 2017. MR3613319 https://doi.org/10.1515/fca-2017-0002
- [30] K. Lê. A stochastic sewing lemma and applications. Electron. J. Probab. 25 (2020) 1–55. MR4089788 https://doi.org/10.1214/20-ejp442
- [31] M. Lifshits and T. Simon. Small deviations for fractional stable processes. Ann. Inst. Henri Poincaré Probab. Stat. 41 (4) (2005) 725–752. MR2144231 https://doi.org/10.1016/j.anihpb.2004.05.004
- [32] W. Liu and M. Röckner. Stochastic Partial Differential Equations: An Introduction. Universitext. Springer, Berlin, 2015. MR3410409 https://doi.org/10.1007/978-3-319-22354-4
- [33] J.-C. Mourrat and H. Weber. Global well-posedness of the dynamic ϕ^4 model in the plane. Ann. Probab. **45** (4) (2017) 2398–2476. MR3693966 https://doi.org/10.1214/16-AOP1116
- [34] C. Mueller, L. Mytnik and E. Perkins. Nonuniqueness for a parabolic SPDE with \\frac{3}{4}-\varepsilon \$-Hölder diffusion coefficients. Ann. Probab. 42 (5) (2014) 2032–2112. MR3262498 https://doi.org/10.1214/13-AOP870
- [35] L. Mytnik and E. A. Perkins. Pathwise uniqueness for stochastic heat equations with Hölder continuous coefficients: The white noise case. Probab. Theory Related Fields 149 (2008) 1–96. MR2773025 https://doi.org/10.1007/s00440-009-0241-7
- [36] L. Mytnik, E. A. Perkins and A. Sturm. On pathwise uniqueness for stochastic heat equations with non-Lipschitz coefficients. Ann. Probab. 34 (2005) 1910–1959. MR2271487 https://doi.org/10.1214/009117906000000331
- [37] E. Neuman. Pathwise uniqueness of the stochastic heat equation with spatially inhomogeneous white noise. Ann. Probab. 46 (6) (2018) 3090– 3187. MR3857853 https://doi.org/10.1214/17-AOP1239
- [38] D. Nualart and Y. Ouknine. Regularization of quasilinear heat equations by a fractional noise. Stoch. Dyn. 4 (2) (2004) 201–221. MR2069689 https://doi.org/10.1142/S0219493704001012
- [39] I. Pinelis. Optimum Bounds for the Distributions of Martingales in Banach Spaces. Ann. Probab. 22 (4) (1994) 1679–1706. MR1331198
- [40] T. Rippl and A. Sturm. New results on pathwise uniqueness for the heat equation with colored noise. *Electron. J. Probab.* 18 (2012) 1–46. MR3101643 https://doi.org/10.1214/EJP.v18-2506
- [41] G. Samorodnitsky and M. S. Taqqu. Stable Non-Gaussian Random Processes. Stochastic Modeling: Stochastic Models with Infinite Variance. Chapman & Hall, New York, 1994. MR1280932

- [42] H. Triebel. Theory of Function Spaces. III. Monographs in Mathematics 100. Birkhäuser Verlag, Basel, 2006. MR2250142
- [43] X. Yang and X. Zhou. Pathwise uniqueness for an SPDE with Hölder continuous coefficient driven by α-stable noise. *Electron. J. Probab.* 22 (2017). MR3613697 https://doi.org/10.1214/16-EJP23