# Doped $\mathrm{MoS}_{2}$ Polymorph for an Improved Hydrogen Evolution Reaction 

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#### Abstract

Green hydrogen produced from solar energy could be one of the solutions to the growing energy shortage as non-renewable energy sources are phased out. However, the current catalyst materials used for photocatalytic water splitting (PWS) cannot compete with other renewable technologies when it comes to efficiency and production cost. Transition-metal dichalcogenides, such as molybdenum disulfides $\left(\mathrm{MoS}_{2}\right)$, have previously proven to have electronic and optical properties that could tackle these challenges. In this work, optical properties, the d-band center, and Gibbs free energy are calculated for seven $\mathrm{MoS}_{2}$ polymorphs using first-principles calculations and density functional theory (DFT) to show that they could be suitable as photocatalysts for PWS. Out of the seven, the two polymorphs $3 \mathrm{H}_{\mathrm{a}}$ and $2 \mathrm{R}_{1}$ were shown to have d-band center values closest to the optimal value, while the Gibbs free energy for all seven polymorphs was within $5 \%$  of each other. In a previous study, we found that $3 \mathrm{H}_{\mathrm{b}}$ had the highest electron  mobility among all seven polymorphs and an optimal bandgap for photocatalytic reactions. The $3 \mathrm{H}_{\mathrm{b}}$ polymorphs were therefore selected for further study. An in-depth analysis of the enhancement of the electronic properties and the Gibbs free energy through substitutional doping with $\mathrm{Al}, \mathrm{Co}, \mathrm{N}$, and Ni was carried out. For the very first time, substitutional doping of $\mathrm{MoS}_{2}$ was attempted. We found that replacing one Mo atom with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, and Ni lowered the Gibbs free energy by a factor of 10 , which would increase the hydrogen evolution reaction of the catalyst. Our study further shows that $3 \mathrm{H}_{\mathrm{b}}$ with one S atom replaced with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, or Ni is dynamically and mechanically stable, while for $3 \mathrm{H}_{\mathrm{b}}$, replacing one Mo atom with Al and Ni makes the structure stable. Based on the low Gibbs free energy, stability, and electronic bandgap $3 \mathrm{H}_{\mathrm{b}}, \mathrm{MoS}_{2}$ doped with Al for one Mo atom emerges as a promising candidate for photocatalytic water splitting.


## 1. INTRODUCTION

Hydrogen produced from water and powered by solar energy is considered one of the most profitable and sustainable alternatives to fossil fuels due to the abundance of water and sunlight. Several technologies are being developed to obtain hydrogen from water, such as electrolysis, photolysis, biological treatment techniques, and thermolysis. ${ }^{1}$ Of these, photocatalytic water splitting is seen as one of the most promising technologies due to its low production cost, good solar-to-hydrogen (STH) efficiency, ease of hydrogen and oxygen separation, and the possibility of both large- and small-scale facilities. ${ }^{2-4}$ However, the current achieved STH, lifetime, and production costs are not good enough to warrant large-scale production facilities. ${ }^{5,6}$
Transition-metal dichalcogenides have been seen as a possible solution to tackle these challenges due to their intriguing electronic and optical properties. Especially $\mathrm{MoS}_{2}$, with its catalytic properties, ${ }^{8,9}$ could be a promising photocatalytic candidate. Atomically $\mathrm{MoS}_{2}$ is a thin material, which means that the catalytic active sites are exposed to the reactants. This feature increases the efficiency of the chemical process. $\mathrm{MoS}_{2}$ is also already used in industry as a catalyst, with established and optimized production methods making it a low-cost material. ${ }^{10}$ Recent work has looked into how $\mathrm{MoS}_{2}$ can be used for
photocatalytic degradation of dyes and found that it is very effective toward organic dyes and organic and toxic pollutants. ${ }^{11-13}$ This is because $\mathrm{MoS}_{2}$ photocatalyzes in the visible region, it has a high specific surface area and an efficient charge separation, and it exists in several different polymorphs. ${ }^{14}$ Through small modifications of the compound composition, crystal structure, or production method, the performance of $\mathrm{MoS}_{2}$ can be significantly altered. ${ }^{15-20}$ The inclusion of dopants or changes in crystal structure has a large impact on the catalytic activity by changing the electronic and optical structure of $\mathrm{MoS}_{2}$. This makes it possible to alter the $\mathrm{MoS}_{2}$ structure in such a way that the STH efficiency is increased. Over the past years, research into $\mathrm{MoS}_{2}$ as a catalyst and cocatalyst for photocatalytic water splitting has increased. ${ }^{18,21-24}$ In a previous article, we investigated 14 different $\mathrm{MoS}_{2}$ polymorphs using density

[^0]
functional theory calculations. Here, we found that seven of these polymorphs were both dynamically and mechanically stable and that they have bandgaps in the range of 1.87 to 2.12 $\mathrm{eV} .{ }^{25}$ This makes them interesting and viable candidates for photocatalytic water splitting.
In this article, we will therefore investigate these seven $\mathrm{MoS}_{2}$ polymorphs further and look into properties influencing their photocatalytic water-splitting potential. The d-band center and Gibbs free energy are used as descriptors for the potential hydrogen production, and the obtained values are compared to the existing literature. In addition, we explore if there are any temperature-dependent phase transitions in the structures and we obtain their absorption spectra. This provides us with a clear overview of the different polymorph's potential as a catalyst for photocatalytic water splitting.

Based on these results and the results from our previous paper, ${ }^{25}$ one of the polymorphs $\left(3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}\right)$ was selected for further experimentation as it was deemed to have the highest potential for PWS based on its absorption spectra, bandgap, and high electron mobility. It was decided to dope $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, and Ni to see how that would affect the Gibbs free energy and the band structure. We chose these elements based on existing work and as they represent a wide array of metals and non-metals. Al-doped $\mathrm{MoS}_{2}$ has been proven to be a stable configuration with tunable carrier density, ${ }^{26}$ and it has shown promise as a catalyst for CO oxidation. ${ }^{27}$ In combination with g $\mathrm{C}_{3} \mathrm{~N}_{4}$, Co-doped $\mathrm{MoS}_{2}$ shows promise as an effective photocatalyst due to the distortions in the $\mathrm{MoS}_{2}$ induced by Co atoms. ${ }^{28}$ Co-doped $\mathrm{MoS}_{2}$ has also been shown to give good results as an HER catalyst for electrochemical water splitting. ${ }^{29}$ N and Ni as dopants have improved the catalytic performance of $\mathrm{MoS}_{2}$ for electrochemical water splitting. ${ }^{30}$ It is reported that the catalytic current density of N -doped $\mathrm{MoS}_{2}$ can reach 15 times that of pristine $\mathrm{MoS}_{2} \cdot{ }^{31} \mathrm{Ni}$ on the other hand was found to add more electrochemical sites, improve the conductivity of the catalyst, and increase the turnover frequency. ${ }^{32}$ Although other halogens such as Cl and F have been used with success to activate both the basal plane sites ${ }^{30}$ and the edge sites, ${ }^{33}$ no work has been done with I. Based on this did we decide to use I as a dopant here to investigate how it would affect $\mathrm{MoS}_{2}$ and its photocatalytic properties. Through DFT calculations, the effects of substitutional doping (both Mo and $S$ were replaced with the dopants) on Gibbs free energy, bandgap, optical spectra, stability, and charge density were investigated. Showcasing that through doping, it is possible to create $\mathrm{MoS}_{2}$ catalysts that have the potential of a high STH efficiency.

## 2. COMPUTATIONAL METHOD

Every calculation was performed using the periodic density functional theory framework through the VASP code. ${ }^{34-38} \mathrm{We}$ used the projector-augmented wave (PAW) method to describe the interaction between the core $(\mathrm{Mo}:[\mathrm{Kr}]$ and $\mathrm{S}:[\mathrm{Ne}])$ and the valence electrons. ${ }^{37,39}$ The structures were optimized using the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional and further optimized with the DFT/vdW-DF2 method. ${ }^{40-42}$ For the phonon calculations, the supercell method was used. The real space force constants of the supercell were calculated with the VASP code, and then the PHONOPY code was used to calculate the phonon frequencies from the force constants in the supercell consisting of at least 32 atoms in all systems. Every atom was displaced by a finite displacement of $0.01 \AA$ in the $x, y$, and $z$ directions to get the force-constant matrices for each binary system. $4 \times 4 \times 4 \mathbf{k}$ points and strict
energy convergence criteria $\left(10^{-8} \mathrm{eV}\right)$ were used for the force constant calculations. We then built a dynamical matrix for the different $\mathbf{q}$ vectors in the Brillouin zone. By solving this matrix, the phonon frequency eigenvalues and phonon mode eigenvector were found. The PHONOPY ${ }^{43}$ code was then used to calculate the thermodynamic properties through a summation of the phonon eigenvectors. Based on this the energy versus temperature curves were generated. Our previous calculations suggested ${ }^{44}$ that structural parameters in oxides could be reliably predicted only by using a large energy cutoff to guarantee basis-set completeness. Hence, we have used a cutoff of 500 eV . The electronic properties were computed by using the screened hybrid functional as proposed by Heyd, Scuseria, and Ernzerhof (HSE06) and generalized gradient approximation $(G G A+U)$ for the polymorphs optimized at the PBE level. ${ }^{45}$ If not specified differently, we used a Monkhorst-Pack $9 \times 9 \times 9 \mathrm{k}$ mesh for the structural optimization and the electronic polymorph studies. Band polymorphs were computed by solving the periodic Kohn-Sham equation on $10 \mathbf{k}$ points along each direction of high symmetry of the irreducible part of the first Brillouin zone. Gibbs free energy and optical properties were calculated using a Gamma $8 \times 8 \times 1 \mathrm{k}$ mesh. VASPKIT ${ }^{46}$ and SUMO ${ }^{47}$ have been used for postprocessing and plot production. For the doped structures, GGA + U was used for the bandgap calculations; see Table 1 for the chosen $U$ values.

Table 1. U Values Chosen for the GGA + U Calculations

| atom | $U$ |
| :--- | :--- |
| Mo | $3.28^{48}$ |
| S | $4.0^{48}$ |
| Al | $4.0^{49}$ |
| Co | $3.0^{50}$ |
| I | $8.0^{51}$ |
| N | $7.0^{52}$ |
| Ni | $6.0^{53}$ |

## 3. RESULTS AND DISCUSSION

### 3.1. Energy vs Volume and Temperature Dependency.

 $\mathrm{MoS}_{2}$ can exist in several different crystal polymorphs or polytypes in three dimensions due to the large variation in stacking sequences and the number of successive $\mathrm{S}-\mathrm{Mo}-\mathrm{S}$ sandwiches along the hexagonal $c$ axis. The different variants are called $1 T, 2 \mathrm{H}, 3 \mathrm{R}, 4 \mathrm{H}_{2}$, and 6 R . Here, the integer stands for the number of $\mathrm{S}-\mathrm{Mo}-\mathrm{S}$ sandwiches per unit cell along the hexagonal $c$ axis, while the $T, H$, and $R$ indicate the type of symmetry: tetragonal, hexagonal, and rhombohedral, respectively. In our previous work, we studied 14 novel polymorphs to understand how the stacking sequence and layer-layer distance influence the properties of the compound. ${ }^{25}$ Seven of these were found to be both mechanically and dynamically stable, making them interesting for further research. ${ }^{25}$ To increase our understanding of the properties of the polymorph, we have now investigated how temperature would influence the relative stability of the polymorphs and if it would induce any phase transitions. Figure 1 shows energy versus temperature curves for the seven stable polymorphs. The starting point for these calculations was the preferred volume found in our previous work. As seen below, increasing the temperature does not change the relative stability between the polymorphs. In fact, only between 400 and 600 K , we can see a small deviation and $3 \mathrm{H}_{\mathrm{a}}$ transitions to be the most stable polymorph at 500 K .

Figure 1. Energy vs temperature curve for the seven studied polymorphs. Zoomed-in insert between 0 and 200 K is added to show the small difference between the structures.

However, the change is only 0.9 eV as seen in Figure 1 and we believe that this phase transition is a numerical artifact caused by the calculation. This means that temperature is not a large factor for these polymorphs and one can choose the most appropriate one for a specific application without having to consider phase transitions.
3.2. Optical Properties. In 1996, Norskov and co-workers proposed the d-band center theory, which correlates the energy of the d-band center of gravity of a metal catalyst with the adsorption energy, activation energy, and dissociation energy of small molecules. ${ }^{54,55}$ They state that the electron density of states (DOS) close to the Fermi level is correlated to the adsorbate substrate adsorption energy due to interactions between electrons occupying d-type orbitals of the metal (dband) and those of the adsorbate. ${ }^{54,55}$ In general, if the d-band center is downshifted compared to the Fermi level, it is an indication of good catalytic activity.

Although there have been significant efforts to verify the dband theory, it has not been proven to be valid for all systems. In general, more complex systems have other factors, not included in the d-band center theory, influencing the catalytic activity. However, our systems are deemed to be of a simple enough character that the d-band center will indicate the catalytic activity. Pt is considered to have an optimal position for the d band center for the most favorable hydrogen binding energy, -1.929 eV with respect to the Fermi level. ${ }^{56}$ Our results are presented in Table 2 below. 1H and 2T have positively shifted dband centers compared to the Fermi level; in addition, they are far away from the optimal value of -1.929 eV . This means that they have higher energy in the antibonding states and the bond between the structure surface and adsorbed material will be

Table 2. d-Band Center for the Stable Polymorphs

| polymorph | d-band center $[\mathrm{eV}]$ | Fermi level $[\mathrm{eV}]$ |
| :---: | :---: | :---: |
| 1 H | 0.334 | -2.552 |
| 2 H | 0.276 | 3.511 |
| $2 \mathrm{R}_{1}$ | -1.520 | 3.285 |
| 2 T | 0.328 | -0.882 |
| $3 \mathrm{H}_{\mathrm{a}}$ | -1.541 | 0.003 |
| $3 \mathrm{H}_{\mathrm{b}}$ | 0.369 | 2.865 |
| 4 T | -0.544 | 0.555 |

higher, ${ }^{57,58}$ leading to a decrease in the chance of photocatalytic reactions. However, the closer the d-band center is to the Fermi level, the better charge carrier transfer. ${ }^{59} 2 \mathrm{H}, 2 \mathrm{R}_{1}, 3 \mathrm{H}_{a}, 3 \mathrm{H}_{\mathrm{b}}$, and 4 T have downshifted d-band centers, indicating that they could be effective photocatalysts, especially $2 \mathrm{R}_{1}$ and $3 \mathrm{H}_{2},-1.520$ and -1.541 eV , respectively, which have d-band center values not far from the optimal value of -1.929 eV . The binding strength of H will be less compared to that of the other structures, thus demanding less energy to be removed from the surface and create $\mathrm{H}_{2}$ atoms.

An important parameter for a catalyst used for photocatalytic water splitting is the absorption coefficient as it determines how far into the material the light will travel before it is absorbed. For water-splitting applications, the catalytic materials must have a high absorption coefficient in the visible region, to ensure that photons are absorbed. In addition to the absorption coefficient, we have also calculated the refractive index of the material to better understand the optical properties of the polymorphs.

All the polymorphs showcase the same rise in absorption coefficient, see Figure 2a, at $2 \mathrm{eV}(620 \mathrm{~nm})$, which corresponds to orange light, followed by a larger and wider upswing at 2.7 to 3.5 eV ( 460 to 355 nm ), which is in the blue and violet range of visible light. This indicates that polymorphs can absorb photons at both ends of the visible light spectra, but there is still a considerable number of photons not absorbed. This could be remedied by introducing dopants that would influence the optical properties of the polymorphs, making the polymorphs tuned for the specific needs of photocatalytic water splitting.

Our calculations show that $2 \mathrm{R}_{1}\left(2.67 \times 10^{5} \mathrm{~cm}^{-1}\right)$ has the highest absorption coefficient within the visible spectra. It is closely followed by 2 H and $3 \mathrm{H}_{\mathrm{b}}$ at $2.33 \times 10^{5}$ and $2.28 \times 10^{5}$ $\mathrm{cm}^{-1}$, while 1 H has the lowest absorption coefficient $\left(1.62 \times 10^{5}\right.$ $\mathrm{cm}^{-1}$ ) in the visible region. $2 \mathrm{R}_{1}, 2 \mathrm{H}$, and $3 \mathrm{H}_{\mathrm{b}}$ are, based on the absorption coefficient, the best candidates for photocatalytic water splitting.

The refractive index (Figure 2b) of the material determines how light propagates inside a material, and the higher the refractive index, the slower light will travel through the material. This results in higher irradiance in the material, which enhances the photocatalytic activity of the material. We see that $2 \mathrm{R}_{1}, 2 \mathrm{H}$, and $3 \mathrm{H}_{\mathrm{b}}$ have the largest peaks around 5 in the visible region. This means that they are classified as high refractive index materials ( $n>3.9$ ), ${ }^{60}$ making them ideal candidates for photonic devices such as solar cells, photonic crystals, and PWS. This is because they have a larger scattering cross section ${ }^{61}$ and smaller mode volumes ${ }^{62}$ than normal refractive index materials. Our results agree with those found by Islam et al. ${ }^{63}$ and we see that $\mathrm{MoS}_{2}$ has a larger refractive index than $\mathrm{TiO}_{2}(2.48-2.62) .{ }^{64}$
3.3. Gibbs Free Energy. Photocatalytic water splitting can be split into two half-reactions: hydrogen evolution reaction (HER) and oxygen evolution reaction (OER), with the overall HER being able to follow two different reaction mechanisms: Volmer-Tafel and Volmer-Heyrovsky. In the Volmer step, protons are reduced to produce adsorbed hydrogen on the electrode surface. After this step, the reaction can either follow the Tafel route or the Heyrovsky route. The Volmer-Tafel reaction involves two adsorbed hydrogen atoms on the electrode surface adjacent to each other combining into a hydrogen molecule. The Volmer-Heyrovsky reaction on the other hand utilizes a proton from the surrounding water to react with an adsorbed hydrogen atom to produce a hydrogen molecule. ${ }^{65}$

$$
2 \mathrm{H}^{+}+2 \mathrm{e}^{-} \rightarrow 2 \mathrm{H}_{\mathrm{ads}} \text { (Volmer) }
$$



Figure 2. Calculated absorption coefficient (a) and refractive index (b).


Figure 3. Location of the adsorbed hydrogen atom. (a) Top, (b) interface, and (c) bridge.

$$
\begin{aligned}
& \mathrm{H}_{\mathrm{ads}}+\mathrm{H}_{\mathrm{ads}} \rightarrow \mathrm{H}_{2(\mathrm{~g})}(\text { Tafel }) \\
& \mathrm{H}^{+}+\mathrm{H}_{\mathrm{ads}}+\mathrm{e}^{-} \rightarrow \mathrm{H}_{2(\mathrm{~g})}(\text { Heyrovsky })
\end{aligned}
$$

where $\mathrm{H}^{+}$is a proton and $\mathrm{H}_{\text {ads }}$ is adsorbed hydrogen atoms.
What is clear from this is that regardless of the path taken by the HER, the adsorption energy of the hydrogen atom is crucial when looking at the catalytic activity of a material. The adsorption energy of hydrogen is linked to the Gibbs free energy ( $\Delta G_{H}$ ) of hydrogen on the compound surface and is widely accepted as an indicator of catalytic activity. ${ }^{66,67}$
The Gibbs free energy can be calculated by the following equation:

$$
\Delta G_{\mathrm{H}}=E_{\mathrm{ads}}+\Delta E_{\mathrm{ZPE}}-T \Delta S_{H}
$$

Here, $E_{\text {ads }}$ is the adsorption energy, $\Delta E_{\text {ZPE }}$ is the zero-point energy difference of $\mathrm{H}_{2}$ in the adsorbed and gas phase state with its values ranging from 0.01 to 0.04 eV , and $T \Delta S_{\mathrm{H}}$ is entropy changes for finite variations at constant temperature $T$. For the materials we study here, $\Delta E_{\text {ZPE }}$ is set to 0.04 eV based on calculations done by Nørskov et al. ${ }^{68}$ The value for the term $\left(T \Delta S_{\mathrm{H}}\right)$ can be approximated to -0.2 eV at $298 \mathrm{~K},{ }^{69}$ which lets us rewrite the equation to $\Delta G_{\mathrm{H}}=E_{\text {ads }}+0.24 \mathrm{eV} .{ }^{68,70,71}$ In general, there are three possibilities for hydrogen adsorption:

1) Exothermic hydrogen adsorption: when this happens, the coverage of hydrogen atoms will be too high and the
reaction will slow down. ${ }^{69}$ This occurs for $\Delta G_{\mathrm{H}}<0$. However, the transformation can occur naturally. ${ }^{72}$ A large and negative $\Delta G_{H}$ also means that $H_{\text {ads }}$ has a weak interaction with the electrode surface and this will slow down the Volmer step, which consequently slows down the overall reaction. ${ }^{67}$
2) Endothermic hydrogen adsorption: here the energy barrier for proton $\left(\mathrm{H}^{+}\right)$formation is too high. ${ }^{69} \Delta G_{\mathrm{H}}>$ 0 . In addition, the natural direction of the reaction is opposite to what is desired (a nonspontaneous reaction). ${ }^{72}$ Here, the $\mathrm{H}_{\mathrm{ads}}$ atoms are bound strongly with the electrode surface, which makes the initial Volmer step easy. However, the following Tafel or Heyrovsky steps difficult. ${ }^{67}$
3) The system is at equilibrium and we have enough of all the reactants. ${ }^{69,72} \Delta G_{H}=0$.
It is well known that an optimal photocatalytic material should have $\Delta G_{\mathrm{H}}=0$ and thus the adsorption energy $E_{\text {ads }}$ should be of the order of -0.24 eV . 67,69

We have calculated the adsorption energies of H at three different locations as it is position-dependent. ${ }^{73}$ Our chosen locations are $S$ top (Figure 3a), $S$ interface (Figure 3b), and $S$ bridge (Figure 3c). The adsorption energy was calculated as

$$
E_{\mathrm{H}}^{\text {ads }}=E_{\mathrm{T}}[\text { system }+\mathrm{H}]-E_{\mathrm{T}}[\text { system }]-\frac{1}{2} E_{\mathrm{H}_{2}}
$$

where $E_{\mathrm{T}}[$ system +H$]$ represents the total energy of $\mathrm{MoS}_{2}$ with an adsorbed H atom, $E_{\mathrm{T}}$ [system] is the energy of $\mathrm{MoS}_{2}$ without an adsorbed molecule, and $E_{\mathrm{H}_{2}}$ is the energy of a hydrogen molecule in the gas phase. The Gibbs free energies are presented in Table 3. From the calculated values, we see that there are not

Table 3. Calculated Gibbs Energy Values for the Seven Stable Polymorphs Calculated at the Three Different Locations

| polymorph | bridge Gibbs $[\mathrm{eV}]$ | inter Gibbs $[\mathrm{eV}]$ | top Gibbs $[\mathrm{eV}]$ |
| :---: | :---: | :---: | :---: |
| 1 H | 1.77 | 1.77 | 1.77 |
| 2 H | 1.75 | 1.75 | 1.85 |
| $2 \mathrm{R}_{1}$ | 1.75 | 1.75 | 1.75 |
| 2 T | 1.77 | 1.77 | 1.77 |
| $3 \mathrm{H}_{\mathrm{a}}$ | 1.77 | 1.77 | 1.76 |
| $3 \mathrm{H}_{\mathrm{b}}$ | 1.75 | 1.75 | 1.75 |
| 4 T | 1.76 | 1.76 | 1.75 |

so many differences in potential catalytic activity among different polymorphs, with Gibbs values ranging from 1.75 eV to 1.85 eV . The positive values of $\Delta G_{\mathrm{H}}$ indicate that the adsorbed hydrogens are bound strongly to the $\mathrm{MoS}_{2}$ surface. This makes the initial Volmer step easy to undergo, while the following Tafel or Heyrovsky steps are more difficult, ${ }^{67}$ which reduces the overall STH efficiency.
It is generally believed that the basal plane of $\mathrm{MoS}_{2}$ is close to inert and it is the edges that contribute to the catalytic activity. ${ }^{74}$ If we compare our results for the basal plane with that of Seo et al. ${ }^{75}$ who got Gibbs energy values ranging from -0.36 to 0.66 eV for $1 \mathrm{~L}, 2 \mathrm{~L}$, and 3 L monolayer $\mathrm{MoS}_{2}$ structure using the edge as active sites, this assumption appears to be true.
3.4. Doped $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ Polymorph. $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ was chosen over the other structures for its better electron mobility, suitable band gap, and absorption coefficient. We performed substitutional doping of both Mo atoms and $S$ atoms. For $S$ substitution, a doping percentage of $6.250 \%$ was chosen, while for Mo $12.5 \%$, that is, we replaced one $S$ atom and one Mo atom.
3.4.1. Gibbs Free Energy for Doped Polymorph. Gibbs free energies calculated for Al, Co, I, N, or Ni doped $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ are shown in Tables 4 and 5 . We found that replacing one $S$ atom

Table 4. Gibbs Energy Values for Substitutional Doped $3 \mathbf{H}_{\mathrm{b}}{ }^{\boldsymbol{a}}$

| Mo replaced on $3 \mathrm{H}_{\mathrm{b}}$ | bridge Gibbs $[\mathrm{eV}]$ | inter Gibbs $[\mathrm{eV}]$ | top Gibbs $[\mathrm{eV}]$ |
| :---: | :---: | :---: | :---: |
| Al | 0.77 | -0.17 | 0.77 |
| Co | 1.17 | -0.02 | 1.18 |
| I | 0.73 | 0.29 | 0.73 |
| N | -0.28 | 0.27 | -0.28 |
| Ni | 0.96 | -0.20 | 0.90 |

${ }^{a}$ One Mo is replaced with either one Al, Co, I, N, or Ni atom.
Table 5. Gibbs Energy Values for Substitutional Doped $3 \mathrm{H}_{\mathrm{b}}{ }^{\boldsymbol{a}}$

| S replaced on $3 \mathrm{H}_{\mathrm{b}}$ | bridge Gibbs $[\mathrm{eV}]$ | inter Gibbs $[\mathrm{eV}]$ | top Gibbs $[\mathrm{eV}]$ |
| :--- | :---: | :---: | :---: |
| Al | -6.99 | -6.45 | -5.55 |
| Co | -6.38 | -4.83 | -5.06 |
| I | -11.18 | -11.16 | -11.16 |
| N | -12.44 | -12.44 | -12.44 |
| Ni | -7.46 | -6.13 | -6.42 |

${ }^{a}$ One S atom is replaced with either one $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, or Ni atom.
with either Al, Co, I, N, or Ni decreased the Gibbs free energy; however, the values are much further away from the optimal
value (zero). In general, negative values indicate that the adsorption process is exothermic and that the surface will be flooded with hydrogen atoms, hindering the other reactants needed for a complete reaction. The large negative values also mean that the $\mathrm{H}_{\text {ads }}$ atom has a weak interaction with the material surface, which results in a slower Volmer step and a reduced overall turnover rate. ${ }^{67}$ Due to the size of the supercell, a smaller doping percentage was not possible, and a larger percentage was not attempted. Nonetheless, the addition of a dopant influenced the $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ structure Gibbs free energy and electronic structure, showcasing the influence of dopants as a powerful tool in developing new photocatalysts.

When the Mo atom is substituted with a dopant atom, we see a clear improvement in the Gibbs free energy as all dopants bring it closer to zero. There is also a greater disparity between the different adsorption sites, where the best results are seen for hydrogen adsorbed at the interface (Figure 3b). This is especially significant for Co that sports a rather high Gibbs free energy for the bridge and top location ( $>1 \mathrm{eV}$ ), while for the interface site, the Gibbs free energy is -0.02 eV , making it an incredible photocatalyst when the hydrogens are adsorbed at this location. Unfortunately, it is not possible to force hydrogen to be adsorbed at specific locations. However, the hydrogen production rate must be tested experimentally before a final verdict can be cast on this compound. N -doped $\mathrm{MoS}_{2}$ has on average the lowest values and they are comparable with that of the edge sites, ${ }^{75}$ which proves that doping the basal layer of $\mathrm{MoS}_{2}$ can be a valid photocatalyst.

Through doping, the electronic structure of the investigated polymorphs is changed, leading to changes in the band structures and bandgaps. Thus, GGA band structure calculations were performed to investigate how the different dopants influence the band gap. It is also important to investigate the structures' stability, band structure, and optical spectra before one can conclude whether a doped structure can perform as a photocatalyst.
3.4.2. GGA Band gaps. The electronic structure is one of the key properties when determining if a material is suitable for photocatalytic processes or photovoltaic cells as it determines the amount of absorbed sunlight.

Our GGA band gap calculations for Mo substituted with Al or I and S substituted with Co or Ni are presented in Figure 4, while the rest are seen in the Supporting Information (Figures S1 and S2). Our calculations show that only four out of the 10 doped structures are semiconductors and those are Mo substituted with Al or I and S substituted with Co or Ni. Of these four, Mo substituted with Al appears to have an intermediate bandgap.

Mo substituted with $\mathrm{Co}, \mathrm{N}$, and Ni and S substituted with Al are metals, while S substituted with N is a semi/anisotropic metal. This means that the material becomes a metal or semiconductor depending on the direction taken through the Brillouin zone, which makes them non-appropriate for photocatalytic applications. When S is substituted with I, the material becomes an insulator making it unsuitable as a catalyst.
3.4.3. Formation Energy. We calculated the formation energy to compare the stability of the doped structures with each other. This was done with the following expression: ${ }^{76}$

$$
\Delta H=E_{\text {doped }}-\left[E_{\text {structure }}+E_{\text {dopant }}-E_{\text {atom }}\right]
$$

where $E_{\text {doped }}$ and $E_{\text {structure }}$ are the total energies of the supercell with and without dopants, respectively, $E_{\text {dopant }}$ is the total energy of the dopant, and $E_{\text {atom }}$ is the total energy of the replaced atom. The results are shown below in Table 6. From these calculations,


Figure 4. GGA band structures for doped $3 \mathrm{H}_{\mathrm{b}}$. Here, Mo is substituted with $\mathrm{Al}(\mathrm{a})$, Mo is substituted with $\mathrm{I}(\mathrm{b})$, S is substituted with Co (c), and S is substituted with $\mathrm{Ni}(\mathrm{d})$.

Table 6. Calculated Formation Energy for the Doped $3 \mathrm{H}_{\mathrm{b}}$ $\mathrm{MoS}_{2}$ Structure

| structure | formation energy $(\Delta H)$ |
| :--- | :---: |
| $\mathrm{Mo}_{7} \mathrm{~S}_{16} \mathrm{Al}$ | 3.19 |
| $\mathrm{Mo}_{7} \mathrm{~S}_{16} \mathrm{Co}$ | 3.51 |
| $\mathrm{Mo}_{7} \mathrm{~S}_{16} \mathrm{I}$ | 7.68 |
| $\mathrm{Mo}_{7} \mathrm{~S}_{16} \mathrm{~N}$ | 3.85 |
| $\mathrm{Mo}_{7} \mathrm{~S}_{16} \mathrm{Ni}$ | 4.27 |
| $\mathrm{Mo}_{8} \mathrm{~S}_{15} \mathrm{Al}$ | 11.13 |
| $\mathrm{Mo}_{8} \mathrm{~S}_{15} \mathrm{Co}$ | 11.37 |
| $\mathrm{Mo}_{8} \mathrm{~S}_{15} \mathrm{I}$ | 15.62 |
| $\mathrm{Mo}_{8} \mathrm{~S}_{15} \mathrm{~N}$ | 11.79 |
| $\mathrm{Mo}_{8} \mathrm{~S}_{15} \mathrm{Ni}$ | 12.22 |

it is clear that replacing Mo with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, or Ni results in a more thermodynamically stable structure compared to replacing
one S atom in $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$. Mo substituted with Al appears to be the most stable of the doped structures, closely followed by Mo substituted with Co and then N. The calculated formation energies for all doped materials results are positive, which means that these compounds will not be created spontaneously and external factors (e.g., pressure, temperature, or energy) are needed to create them as the reactants are more stable than the product.
3.4.4. Mechanical Stability. In addition to influencing the electronic structure, the dopants will also affect the mechanical and dynamical stability of the polymorph. The elastic constants of a material can be used as a descriptor for how the material will react to an applied force, as either applied strain or stress. Thus, it is possible to use them to understand the mechanical stability of a material.

The stress and strain have three shear and three tensile components. Therefore, it is possible to describe the elastic

Table 7. Calculated Single-Crystal Elastic Constants $C_{i j}$ (in GPa), Bulk Modulus B (in GPa), Shear Modulus G (in GPa), Poisson's Ratio $\nu$, and Young's Modulus E (in GPa) ${ }^{a}$

| polymorph | $3 \mathrm{H}_{\mathrm{B}}{ }^{25}$ | Al for Mo | Co for Mo | I for Mo | N for Mo | Ni for Mo | Al for S | Co for S | I for S | N for S | Ni for S |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| crystal system | hexagonal | hexagonal | hexagonal | hexagonal | oblique | hexagonal | hexagonal | hexagonal | hexagonal | hexagonal | hexagonal |
| $\mathrm{C}_{11}$ | 176 | 223 | 161 | 212 | 254 | 196 | 301 | 300 | 296 | 326 | 302 |
| $\mathrm{C}_{12}$ | 45 | 68 | 72 | 56 | -10 | 90 | 99 | 103 | 93 | 107 | 100 |
| $\mathrm{C}_{13}$ | 0.6 | 0 | 0 | 0 |  | 0 | 0 | 0 | 0 | 0 | 0 |
| $\mathrm{C}_{16}$ |  |  |  |  | 23 |  |  |  |  |  |  |
| $\mathrm{C}_{26}$ |  |  |  |  | -53 |  |  |  |  |  |  |
| $\mathrm{C}_{66}$ | 46 | 77 | 44 | 78 | 116 | 53 | 101 | 98 | 101 | 109 | 101 |
| Born | yes | yes | yes | yes | yes | yes | yes | yes | yes | yes | yes |
| $\mathrm{B}_{\mathrm{V}}$ | 50 | 146 | 116 | 134 | 119 | 143 | 200 | 202 | 195 | 216 | 201 |
| $\mathrm{B}_{\mathrm{R}}$ | 2 | 146 | 116 | 134 | 117 | 143 | 200 | 202 | 195 | 216 | 201 |
| $\mathrm{B}_{\mathrm{H}}$ | 26 | 146 | 116 | 134 | 118 | 143 | 200 | 202 | 195 | 216 | 201 |
| $\mathrm{G}_{\mathrm{V}}$ | 34 | 77 | 44 | 78 | 123 | 53 | 101 | 98 | 101 | 109 | 101 |
| $\mathrm{G}_{\mathrm{R}}$ | 0.8 | 77 | 44 | 78 | 109 | 53 | 101 | 98 | 101 | 109 | 101 |
| $\mathrm{G}_{\mathrm{H}}$ | 17 | 77 | 44 | 78 | 116 | 53 | 101 | 98 | 101 | 109 | 101 |
| E | 43 | 202 | 128 | 197 | $\begin{gathered} 171-292 \\ (232) \end{gathered}$ | 154 | 268 | 264 | 267 | 291 | 269 |
| G | 17 | 77 | 44 | 78 | $\begin{gathered} 103-116 \\ (110) \end{gathered}$ | 53 | 101 | 98 | 101 | 109 | 101 |
| $\nu$ | 0.23 | 0.306 | 0.450 | 0.263 | $\begin{gathered} -0.167- \\ 0.197 \\ (0.015) \end{gathered}$ | 0.462 | 0.329 | 0.345 | 0.315 | 0.328 | 0.330 |
| $\mathrm{G}_{\mathrm{H}} / \mathrm{B}_{\mathrm{H}}$ | 0.653 | 0.527 | 0.379 | 0.582 | 0.932 | 0.370 | 0.505 | 0.485 | 0.518 | 0.505 | 0.502 |

${ }^{a}$ Subscript V indicates the Voigt bound, R indicates the Reuss bound, and H indicates the Hill bound.
constants of a crystal using a $6 \times 6$ symmetric matrix with 27 components. 21 of those 27 are independent of the others, and in addition, we can use any existing symmetry in the crystal to reduce the number of components. The stiffness matrix, $C_{i j}$, which the $6 \times 6$ matrix also is known as, can be used to calculate the bulk modulus, Poisson coefficient, and Lame constants. In general, it has been found that the DFT elastic constants are within $10 \%$ of the experimental values. ${ }^{77}$
For hexagonal polymorphs, the stability criteria are ${ }^{78}$

$$
\begin{aligned}
& \mathrm{B}_{\mathrm{H} 1}=C_{11}>\left|C_{12}\right| \\
& \mathrm{B}_{\mathrm{H} 2}=C_{11}>0
\end{aligned}
$$

and for oblique polymorphs the stability criteria are ${ }^{78}$

$$
\begin{aligned}
& \mathrm{B}_{\mathrm{O} 1}=C_{11}>0 \\
& \mathrm{~B}_{\mathrm{O} 2}=C_{11} C_{22}>C_{12} C_{12} \\
& \mathrm{~B}_{=3}=\left|C_{i j}\right|>0
\end{aligned}
$$

The $3 \mathrm{H}_{\mathrm{B}}$ structure was found to be mechanically stable with all five dopants and for both Mo and $S$ substitutional doping as they all fulfill the Born criteria. This means that the doped structure can keep its pore size and structure under a mechanical load. To further investigate how the doped structure would react to applied mechanical forces, we used the stiffness matrix to calculate the Voigt(V), Reuss (R), and Hill (H) moduli. With these, we calculated the bulk modulus $B$, shear modulus $G$, Young's modulus $E$, and Poisson's ratio $\nu$ for all doped $3 \mathrm{H}_{\mathrm{b}}$ $\mathrm{MoS}_{2}$ structures. The values are shown in Table 7.

Young's modulus characterizes how a structure reacts to a force applied lengthwise, both under compression and extension. All of our doped structures have positive Young's modulus ranging from 128 GPa (higher than for bronze, titanium, and annealed copper ${ }^{79}$ ) up to 291 GPa (beryllium ${ }^{79}$ ), indicating that they are quite stiff materials and would react slowly to applied pressure. To determine whether the doped
structures are ductile or brittle, we calculated the shear bulk modulus ratio, where 0.5 is the cutoff below which materials are categorized as ductile. ${ }^{48}$ Our results show that seven of the doped structures are brittle, although three of them are just above the threshold. As they are so close to the threshold values, experimental tests would be needed to determine if $\mathrm{MoS}_{2}$ structures with one S atom substituted for one $\mathrm{Al}, \mathrm{N}$, or Ni atom are brittle. Our simulations show that only when Co and Ni are substituted for Mo or Co substituted for S in $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ do the doped materials become ductile. ${ }^{80,81}$

Poisson's ratio is a measure of the Poisson effect, a material's deformation perpendicular to the direction of the applied force. A value close to zero indicates that the material's structure (e.g., foam) would collapse in the compression, while materials with a much larger bulk modulus than shear modulus (rubber) will have a Poisson ratio around $0.5 .{ }^{82} \mathrm{Co}$ and Ni substitutions for Mo have quite high values ( 0.450 and 0.462 ), putting them in the upper region together with gold and saturated clay. I substitution for Mo on the other hand is in the mid-range where similar materials would be magnesium and variants of cast iron. The rest are in the range of 0.306 to 0.345 , which is similar to the Poisson ratio of stainless steel, aluminum alloy and copper. Compared to the values obtained for $3 \mathrm{H}_{\mathrm{b}}$ in our previous paper, we see that the shear modulus has increased for the doped structures, indicating that they will withstand higher values of shear stress. Adding dopants has also made the structures ductile in the case of Co and Ni substitution for Mo and Co substitution for $S$. This showcases the power of doping when it comes to altering and improving mechanical properties of materials.
3.4.5. Dynamical Stability. It is not enough for a material to be mechanically stable; it must also be dynamically stable. To verify this, we performed phonon and phonon density of states calculations. In addition, we have calculated the phonon dispersion curves, at the equilibrium volume, along the high symmetry direction of the Brillouin zone and these results are seen in Figures 5 and 6. The $3 \mathrm{H}_{\mathrm{b}}$ structures where one Mo atom


Figure 5. Phonon density of states for $3 \mathrm{H}_{\mathrm{b}}$ doped with $\mathrm{Al}(\mathrm{a}), \mathrm{Co}(\mathrm{b}), \mathrm{I}(\mathrm{c}), \mathrm{N}(\mathrm{d})$, and $\mathrm{Ni}(\mathrm{e})$ for Mo .
was replaced with one $\mathrm{Co}, \mathrm{I}$, or one N atom have clear and distinctive negative frequencies, as seen in Figure 5b-d, respectively, which makes the dynamically unstable structures.
In general, if a material is mechanically stable and dynamically unstable, it could have a metastable phase and by using dopants, nanoparticles and nanoobjects the structures can be stabilized.
For Mo replaced with either one Al or one Ni atom, small and similar negative modes are seen for both structures between A-L, М-Г, Г-К, and Н-А.

The imaginary modes seen in Figure 5a,e are believed to be caused by lattice distortions stemming from the creation of the supercell. The effect of these distortions can potentially be dealt with by changing the supercell size. Phonon calculations also require extremely $\left(10^{-8}\right)$ accurate forces, ${ }^{83}$ while the volume relaxation calculations were performed using $10^{-6}$ as the accuracy limit. This difference in accuracy could have caused these small imaginary modes we see for Al or Ni . Another possible explanation is a slightly wrong magnetic moment in the
input files for the phonon calculations. Wolloch et al. found that quite small differences in the magnetic moment can help stabilize a structure. ${ }^{83}$ Based on this, we conclude that these two structures are dynamically stable.

For the S-substituted structures presented in Figure 6, we see the same type of negative modes, as seen in Figure 5a,e, appear for $S$ replaced with Mo (Figure 6a) and I (Figure 6c). While for $S$ replaced with either $\mathrm{Co}, \mathrm{N}$, or Ni, Figure $6 \mathrm{~b}, \mathrm{~d}, \mathrm{e}$ respectively, the negative frequency is near zero and an artifact from the calculations. Thus, we can conclude that all the S-substituted structures are dynamically stable.
3.4.6. Charge Density Calculations. The charge density map shown in Figure 7a for $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ clearly shows that the electrons reside around the individual atoms. In addition, the spherically shaped charge distribution indicates that the bonding between $\mathrm{Mo}-\mathrm{S}$ and $\mathrm{S}-\mathrm{S}$ is primarily ionic. The charge transfer plots in Figure 7b reveal a depletion of charge around the Mo atoms, while we get a charge accumulation around the $S$ atoms

(a)

(c)

(e)

(b)

(d)

Figure 6. Phonon density of states for $3 \mathrm{H}_{\mathrm{b}}$ with S substituted for $\mathrm{Al}(\mathrm{a})$, $\mathrm{Co}(\mathrm{b}), \mathrm{I}(\mathrm{c}), \mathrm{N}(\mathrm{d})$, and $\mathrm{Ni}(\mathrm{e})$.


Figure 7. Calculated charge density map (a), charge transfer plot (b), and electron localization function (c) for $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$.
and between the Mo and $S$ atoms. As the charge transfer plot shows, the charge gathers around the $S$ atoms and indicates a slight polar character for the compound The electron localization function (ELF), Figure 7c, is a measure of the spatial localization of electrons that helps to map electron pair probability in multielectron systems. Our ELF findings are consistent with that of the charge transfer plot, and we can thus expect the electron to be found outside of the $S$ atoms.

## 4. CONCLUSIONS

Seven $\mathrm{MoS}_{2}$ polymorphs were studied with regard to phase transition and their photocatalytic capabilities were analyzed using catalytic indicators such as d-band center, Gibbs free energy, and absorption coefficient. The most promising candidate, $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$, was chosen for further enquiries and was substitutionally doped with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, and Ni atoms at two different locations. This study shows the following:

- No phase transitions caused by temperature changes were found for $1 \mathrm{H}, 2 \mathrm{H}, 2 \mathrm{R}_{1}, 2 \mathrm{~T}, 3 \mathrm{H}_{\mathrm{a}}, 3 \mathrm{H}_{\mathrm{b}}$, and 4 T polymorphs.
- All seven polymorphs have their absorption peak at 620 nm in the visible spectra.
- $2 \mathrm{R}_{1}$ has the highest maximum $\left(2.67 \times 10^{5} \mathrm{~cm}^{-1}\right)$ followed by 2 H and $3 \mathrm{H}_{\mathrm{b}}$ at $2.33 \times 10^{5}$ and $2.28 \times$ $10^{5} \mathrm{~cm}^{-1}$, respectively.
- $2 \mathrm{R}_{1}, 2 \mathrm{H}, 3 \mathrm{H}_{\mathrm{b}}, 4 \mathrm{~T}$, and $3 \mathrm{H}_{\mathrm{a}}$ are the high refractive index materials ( $n>3.9$ ).
- Different dopants of $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ resulted in varying band gaps.
- Only four were seen to be semiconductors: Mo substituted with Al and I and S substituted with Co and Ni.
- S substituted with I is an insulator.
- The remaining five structures, Mo substituted with $\mathrm{Co}, \mathrm{N}$, or Ni and S substituted with Al or N , are metallic.
- Mo substitutional doped $\mathrm{MoS}_{2}$ with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, and Ni reduces the Gibbs free energy by a factor of 10 .
- $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ doped with N for $\mathrm{Mo}(-0.28-0.27 \mathrm{eV})$ has a lower Gibbs free energy than the edge sites of $\mathrm{MoS}_{2}$.
- Only Co, I, and N for Mo-doped $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ are dynamically unstable.
- The other seven structures are both mechanically and dynamically stable.
- Structures with Mo replaced are more thermodynamically stable than structures with S substituted.
Our calculations show that $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ doped with Al and Ni for one Mo atom and $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ doped with $\mathrm{Al}, \mathrm{Co}, \mathrm{I}, \mathrm{N}$, and Ni for one S atom are dynamically and mechanically stable. This makes it possible to synthesize and perform experimental tests on these structures. Through doping, we lowered the Gibbs free energy of $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ and made the basal plane competitive with the edge sites, especially for N for Mo. Our study indicates that $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ substitutionally doped with Al for Mo atoms is a promising candidate for photocatalytic water splitting. However, the best doping percentage and optimal dopant need to be found through further theoretical and experimental studies. Future work should also investigate if the dopants introduce new photocatalytic sites and through that increases the overall STH efficiency. In addition, it should be looked into if the dopant promotes carrier separation and carrier mobility.


## ASSOCIATED CONTENT

## (s) Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsomega.3c02623.

GGA band structure for $3 \mathrm{H}_{\mathrm{b}}-\mathrm{MoS}_{2}$ substitutional doped with Al, Co, I, N, and Ni for either one Mo atom or one S atom (PDF)

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## Author Contributions

H.E.: conceptualization, methodology, software, formal analysis, writing - original draft, writing - review \& editing, and visualization; D.V.: supervision, resources, validation, writing review and editing; P.V.: validation, supervision, writing review and editing.

## Notes

The authors declare no competing financial interest.

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